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3DDT, A Three-Dimensional Multigroup Diffusion-Burnup Program

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# 3DDT, A Three-Dimensional Multigroup Diffusion-Burnup Program

by

John C. Vigil



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#### 3DDT, A THREE-DIMENSIONAL MULTIGROUP DIFFUSION-BURNUP PROGRAM

#### by

John C. Vigil

#### ABSTRACT

3DDT is a three-dimensional (X-Y-Z or R- $\theta$ -Z) multigroup diffusion theory code. It can be used to compute  $k_{eff}$  or to perform criticality searches on reactor composition, time absorption, or reactor dimensions by either the regular or the adjoint flux equations. Material burnup and fission production buildup can be computed for specified time intervals, and criticality searches can be performed during burnup to compensate for fuel depletion and fission product growth.

All programming is in FORTRAN-IV and variable dimensioning is used to make maximum use of available core storage. Because variable dimensioning is used, no simple restrictions can be placed on individual components of the problem. However, a 16-group problem containing 30 x 30 x 30 mesh points and 80 zones can be accommodated on a 65k computer. Execution times are about 0.0l sec per mesh point per group for a  $k_{eff}$  calculation on a CDC 6600 computer.

#### I. INTRODUCTION AND SUMMARY

The 3DDT (<u>3</u>-<u>D</u>imensional <u>D</u>iffusion <u>Theory</u>) code is an extension to three space dimensions of the 2DB two-dimensional code.<sup>1</sup> Except that the geometry options are X-Y-Z or R- $\theta$ -Z and that programming is for the CDC 6600 computer, 3DDT contains all of the 2DB features. Some of these features are:

- Multigroup calculations of k<sub>eff</sub> or criticality searches on reactor composition, reactor dimensions, or time absorption (alpha) by means of either the regular or the adjoint flux equations.
- Material burnup and fission product buildup can be computed for specified time intervals, and criticality searches can be performed during burnup to compensate for fuel depletion and fission product growth.
- All programming is in FORTRAN-IV and variable dimensioning is used to make maximum use of available storage.
- Convergence is accelerated by group rebalancing, successive overrelaxation, and line inversion.
- Alpha and k<sub>eff</sub> can be used as parametric eigenvalues.
- DTF-IV input formats<sup>2</sup> are used.
- Only downscattering is treated.

A description of the mathematical model is given in the body of the report and follows closely the description given in the 2DB user's manual.<sup> $\perp$ </sup> Input instructions, storage requirements, a simplified logical flow diagram, a description of the variable-dimensioned arrays and nonsubscripted common variables, and an executed sample problem are included, respectively, in Appendices A through E. Although the logical flow of the code is patterned after the 2DB code (which in turn is patterned after the DTF-IV and 2DF codes<sup>3</sup>), a considerable amount of programming was required to extend 2DB to three space dimensions. Most of the programming effort was the result of the storage arrangements and the concomitant transfer of data from one storage medium to another.

Both Extended Core Storage (ECS) and disk storage are utilized in 3DDT. In general, fourdimensional arrays, e.g.,  $\phi(x,y,z,g)$ , are stored on the disk using tape file simulation; threedimensional arrays, e.g.,  $\phi(x,y,z)$  for a particular energy group, are stored in the ECS using random access; and two-dimensional arrays, e.g.,  $\phi(x,y)$ for a particular energy group and axial position,

are stored in the central memory. Thus, central memory storage requirements are insensitive to the number of energy groups and the number of axial mesh points.

Because variable dimensioning is used, no simple restrictions can be placed on individual components of the problem such as number of energy groups or number of mesh points. However, most threedimensional problems of practical interest (from the viewpoint of computing time) can be accommodated on a machine with a 130k central memory and many can be accommodated on a 65k machine. Because of the manner in which arrays are stored in 3DDT, very large two-dimensional (in R-Z or X-Z geometry) problems can be run on a 65k machine.

Eigenvalues and fluxes computed with 3DDT for several test problems agreed very well with results obtained with 2DB for corresponding two-dimensional problems. Execution times are similar to those for 2DB (for equal size problems) and are of the order of 0.01 to 0.02 sec per mesh point per group. Thus, a six-group problem containing 30 x 30 x 30 mesh points would require from 30 to 60 min of computing time on the CDC 6600 computer. The low estimate applies to a  $k_{eff}$  calculation, and the high estimate applies to an implicit eigenvalue search calculation. On the CDC 7600 computer, which is reported to be about four times faster than the CDC 6600, the utility of 3DDT will be greatly enhanced.

Possible areas for future work on 3DDT include (not necessarily in order of priority):

- Incorporation of internal boundary conditions (principally for use in control rod problems).
- 2. Incorporation of a spherical  $(R-\theta-\phi)$  and/or a hexagonal-Z geometry option in addition to the X-Y-Z and R- $\theta-Z$  options currently available.
- Computation of neutron balance by zone (the present version computes only the overall neutron balance).
- Computation of neutron currents for use in a three-dimensional perturbation code.
- 5. Extension of the burnup routines to allow computation of buildup of individual fission products such as xenon and samarium. In the present version, fission products are treated as a single nuclide and are characterized by a single decay constant and a single crosssection set.
- Incorporation of upscattering (the present version is limited to downscattering).

- 7. Incorporation of a source option.
- Use of sequential ECS access (tape file simulation) to make 3DDT compatible with computing systems for which ECS is not available. The present version uses random ECS access, which is much faster than sequential file access.
- II. FORMULATION OF DIFFERENCE EQUATIONS

The neutron balance equations, in the multigroup diffusion approximation with no upscatter, are

$$\nabla \cdot D_{g} \nabla \phi_{g} - \Sigma_{g}^{r} \phi_{g} + S_{g} = 0$$
 (1)

(g≖1,2,...,IGM),

 $s_{g} = \frac{\chi_{g}}{k_{eff}} \sum_{g'=1}^{IGM} \nu \Sigma_{g'}^{f}, \phi_{g'} + \sum_{g'=1}^{g-1} \Sigma_{g' \rightarrow g} \phi_{g'} \quad (2)$ 

and

where

IGM = number of energy groups,

g = energy group index,

Σ°

- $\phi_g$  = neutron flux in group g,
- S = neutron source for group g from fission and downscatter from higher energy groups,

$$D_g = \text{diffusion coefficient for group g}$$
$$D_g = 1/3 \Sigma_g^{\text{tr}},$$

- $\sum_{g}^{tr}$  = macroscopic transport cross section for group g,
- $v\Sigma_g^f$  = average number of neutrons per fission times the macroscopic fission cross section for group g,
- $\sum_{g' \to g} = \text{macroscopic transfer cross section from } g'$ to g,
  - $\Sigma^{r} = macroscopic removal cross section for group g given by$

$$= \Sigma_g^a + \sum_{g'=g+1}^{IGM} \Sigma_{g+g}, \quad ,$$

- $\Sigma^{a} = macroscopic absorption cross section for g group g,$
- $\chi_g$  = fraction of fission neutrons born in group g, and

k<sub>eff</sub> = effective multiplication factor.

The spatial difference equations are obtained by integrating Eq. (1) over the volume associated with each mesh point. Each mesh point is assumed to be at the center of its associated mesh volume. The mesh point arrangement in X-Y-Z geometry is illustrated in Fig. 1. For the (i,j,k) mesh point, labelled 0 in Fig. 1, the radial integration is from  $(x_i - \Delta x_o/2)$  to  $(x_i + \Delta x_o/2)$ , the rotational integration is from  $(y_j - \Delta y_o/2)$  to  $(y_j + \Delta y_o/2)$ ,



Fig. 1. Mesh point arrangement in X-Y-Z geometry.

and the axial integration is from  $(z_k - \Delta z_0/2)$  to  $(z_k + \Delta z_0/2)$ . In this report, the radial direction refers to either the R or the X direction, the rotational direction refers to either the  $\theta$  or the Y direction, and the axial direction refers to the Z direction.

The leakage terms are obtained by first transforming the volume integral to a surface integral using the divergence theorem,

$$\int \nabla \cdot D \nabla \phi \, dV = \int D \nabla \phi \cdot \vec{n} \, dA , \qquad (3)$$

where  $\vec{n}$  is a unit vector normal to the surface. To evaluate the surface integral, the flux gradients at the mesh volume boundaries are approximated by using the two adjacent flux values. The result of volume integration of Eq. (1) for mesh point 0 (see Fig. 1) is thus

$$\sum_{k=1}^{6} \frac{\overline{D}_{k} A_{k} (\phi_{k} - \phi_{o})}{\ell_{k}} - \Sigma_{o}^{r} \phi_{o} V_{o} + S_{o} V_{o} = 0 , \qquad (4)$$

where the group index has been omitted for simplicity and

- $\Sigma^{r}_{o}$  = macroscopic removal cross section associated with mesh point 0,
- S = total neutron source rate associated with mesh point 0,
- $V_{o}$  = volume associated with mesh point 0,
- $\phi_k$  = neutron flux at mesh point k,
- $\ell_k = \text{distance between mesh point } k$  and mesh point 0,
- ${\bf A}_{k}$  = area of the boundary between mesh point  ${\bf k}$  and mesh point 0, and
- $\bar{D}_{k}$  = effective diffusion coefficient between mesh point k and mesh point 0, e.g.,

$$\tilde{D}_{1} = \frac{D_{o} D_{1}(\Delta x_{o} + \Delta x_{1})}{D_{o} \Delta x_{1} + D_{1} \Delta x_{o}},$$

For calculational purposes, Eq. (4) can be written in the simpler form

$$\sum_{k=0}^{6} c_{k} \phi_{k} = s_{0} v_{0} , \qquad (5)$$

where

$$C_{k} = -\frac{\bar{D}_{k} A_{k}}{\ell_{k}}$$
 (k=1,2,...,6) (6)

and

$$C_0 = \Sigma_0^r V_0 - \sum_{k=1}^6 C_k$$
 (7)

The constants  $C_k$  (k=0,1,...,6) are computed for the initial system and stored for use in the flux calculation. They must be recomputed whenever material compositions change (as in concentration and alpha searches or in depletion calculations) or whenever changes are made in the mesh intervals (as in delta calculations).

Three boundary conditions are available: reflective, vacuum, and periodic. The periodic boundary condition is available only for the  $\theta$  (or Y) boundaries. In the flux calculation, the boundary conditions enter only through the values of the constants  $C_k$  at the reactor boundaries. To illustrate for reflective and vacuum boundary conditions, consider the one-dimensional slab reactor shown in Fig. 2. In the figure, the left-hand boundary is reflected and the right-hand boundary is free (vacuum).



Fig. 2. Slab reactor with reflective condition at left boundary and vacuum condition at right boundary.

<u>Reflective Condition.</u> Imagine that a pseudomesh interval, interval 0, has been added on the left-hand side of the left boundary with the same composition and thickness of interval 1. Since  $\nabla \phi = 0$  at the left boundary,  $\phi_0 = \phi_1$  and  $\phi_0 - \phi_1$ vanishes. Therefore the coefficient,  $C_1$ , of  $\phi_0 - \phi_1$ in Eq. (4) is immaterial and may be set equal to zero. The calculation is performed assuming that  $C_1 = 0$ , and thus the imaginary flux  $\phi_0$  does not contribute to Eq. (5).

<u>Vacuum Condition.</u> Now imagine that a pseudomesh interval, interval IM + 1, with the same composition as interval IM and thickness 1.42  $\lambda^{tr}$ , where  $\lambda^{tr}$  is the transport mean free path in interval IM, has been added to the right-hand side of the right boundary. Since  $\phi_{IM} \neq 0$  and  $\phi_{IM+1} = 0$ , the coefficient,  $C_{IM+1}$ , of  $\phi_{IM} - \phi_{IM+1}$  in Eq. (4) cannot be disregarded. From Eq. (6), it is seen that

$$C_{IM+1} = -\frac{D_{IM} A_{IM+1}}{0.5 \Delta x_{IM} + 0.71 \lambda^{tr}}, \qquad (8)$$

where it is assumed that  $\lambda^{tr} = 1/\Sigma^{tr}$ . As in the reflective case, there is no contribution of the imaginary flux to Eq. (5). For the reflective case,  $C_{\rm L} = 0$ ; whereas for the vacuum case,  $\phi_{\rm L} = 0$ .

<u>Periodic Condition.</u> To illustrate the manner in which periodic boundary conditions enter into the computation of the constants  $C_k$  at the boundaries, consider the one-dimensional slab reactor in Fig. 3. On the left-hand side of the left boundary, an imaginary mesh interval, interval 0, of the same composition and thickness as interval IM has been



Fig. 3. Slab reactor with periodic boundary conditions.

added. Similarly, an imaginary mesh interval, interval IM + 1, of the same composition and thickness as interval 1 has been added on the right-hand side of the right boundary. For periodic boundary conditions,  $\phi_o = \phi_{IM}$  and  $\phi_1 = \phi_{IM+1}$ . Since in general  $\phi_1 \neq \phi_{IM}$  (thus  $\phi_o \neq \phi_1$  and  $\phi_{IM} \neq \phi_{IM+1}$ ), the coefficients  $C_1$  and  $C_{IM+1}$  cannot be disregarded. From Eq. (6), it is clear that

$$C_{IM+1} = C_1 = \frac{D_1 A_1}{0.5(\Delta x_1 + \Delta x_{IM})}$$
, (9)

where

$$\overline{D}_{1} = \frac{D_{1} D_{IM} (\Delta x_{1} + \Delta x_{IM})}{D_{1} \Delta x_{IM} + D_{IM} \Delta x_{1}}$$
(10)

The imaginary mesh intervals discussed in the examples above are not a part of the code. They are mentioned only to illustrate the manner in which the boundary conditions enter into the calculation of the constants  $C_k$  at the reactor boundaries.

The adjoint form of the multigroup diffusion equations is the same as Eq. (1) except that the source term, Eq. (2), is changed to

$$s_{g}^{\star} = \frac{\nu \Sigma_{g}^{f}}{k_{eff}} \sum_{g' = 1}^{IGM} \chi_{g'} \phi_{g'}^{\star} + \sum_{g' = g+1}^{IGM} \Sigma_{g \neq g'} \phi_{g'}^{\star} , \qquad (11)$$

where  $\phi_g^*$  is the adjoint flux in group g. Note that the adjoint of a downscatter problem has the characteristics of a more time-consuming upscatter problem in that the flux in groups g<sup>1</sup> > g contribute to the source in group g. The code solves the adjoint form by transposing the scattering matrix and redefining the fission source term as in Eq. (11). In addition, the code inverts the group order of the cross sections, the fission spectrum, and the neutron velocities. By proceeding in this manner, the upscatter problem becomes a downscatter problem and the calculation can proceed through the same calculational loops as for regular problems. III. SOLUTION OF DIFFERENCE EQUATIONS

Standard source-iteration techniques are used to compute eigenvalues and flux profiles. Using an initial flux guess, an initial fission source distribution is calculated. New flux profiles in each group are sequentially computed, beginning in the highest energy group in regular problems and in the lowest energy group in adjoint problems. The flux in each group is normalized immediately before each group calculation by balancing the total source and loss rates. After the new fluxes in all groups have been calculated, a new fission source distribution is computed from the new flux profiles. The multiplication ratio,  $\lambda$ , is obtained by taking the ratio of the new (current iteration) total fission source to the old (previous iteration) total fission source. A fission source cycle, i.e., a complete pass through all the energy groups, is called an outer iteration. The outer iteration process is continued until the fission source  $(\lambda)$  converges or until either the maximum number of outer iterations (input parameter DO5) or the problem time limit (input parameter ITIMOF) is exceeded.

In  $k_{eff}$  calculations, the fission spectrum,  $\chi$ , is divided by  $\lambda$  so that  $\lambda$  approaches unity as the problem converges. The effective multiplication factor,  $k_{eff}$ , is the product of the successive  $\lambda$ 's, which is equivalent to

$$k_{eff} = \frac{\sum_{g} \chi_{g}^{initial}}{\sum_{g} \chi_{g}^{current}}$$
 (12)

Outer iteration convergence is assumed when both  $|1 - \lambda_j| \leq EPS$  and  $|\lambda_j - \lambda_{j-1}| \leq EPS$ , where EPS is the lambda convergence input parameter and j is the outer iteration index. Fission source over-relaxation is employed to accelerate convergence. After the fission source profile,  $F_j$ , for the jth outer iteration is calculated, a second profile is computed by magnifying the difference between  $F_4$ 

and  $F_{i-1}$  according to

$$F_{j}^{\prime} = F_{j} + \beta^{\prime} (F_{j} - F_{j-1})$$
, (13)

where  $\beta'$  is the fission source overrelaxation factor. After  $F'_j$  is computed from Eq. (13), it is normalized to give the same total neutron source as  $F_j$ .

The group fluxes are computed by horizontal  $(R-\theta \text{ or } X-Y)$  planes beginning with the plane at the lowermost axial position. A complete sweep of a horizontal plane for a particular axial position and energy group is called an  $R-\theta$  (or X-Y) iteration. A complete sweep through the entire spatial mesh for a particular energy group is called an inner iteration. Since an inner iteration involves sweeping the horizontal planes at all axial positions, an inner iteration as defined here is sometimes referred to as a Z iteration.

For each energy group, the inner iteration process is continued until either convergence is achieved or the maximum number of inner iterations is exceeded. Two tests must be satisfied for inner iteration convergence. The first test requires that

$$\frac{IGM \int v \Sigma_{g}^{f} \left| \phi_{g}^{i} - \phi_{g}^{i-1} \right| dV}{\sum_{g} \int v \Sigma_{g}^{f} \phi_{g}^{j-1} dV} \leq EPS, \qquad (14)$$

where the integration is over the entire mesh, i is the inner iteration index, and j is the outer iteration index. Equation (14) is basically a convergence test on the total fission source, with the additional feature that the tightest convergence is required of those groups which contribute the most to the fission source. If the first test is satisfied, the code then requires that

$$\left| \left( \phi_{g}^{i} - \phi_{g}^{i-1} \right) / \phi_{g}^{i-1} \right| \leq GO6$$
 (15)

be satisfied at each mesh point. GO6 is the pointwise flux convergence input parameter and is, in general, larger than EPS. Two values for the maximum number of inner iterations per group are specified in the input. The first, SO4, is applied if  $|1 - \lambda| > 10$ EPS and the second, GO7, is applied otherwise. For each outer iteration, the number of inner iterations is at least equal to IGM and is not larger than IGMxIIPG, where IGM is the number of energy groups and IIPG is equal to either SO4 or GO7 as appropriate.

During an inner iteration, the R- $\theta$  plane at each axial position is swept until either the fluxes converge or the maximum number of R- $\theta$  iterations per plane is exceeded. The test for R- $\theta$  convergence is the same as the pointwise test, Eq. (15), for inner iteration convergence. The maximum number of R- $\theta$  iterations per plane is computed internally from IIPG/M05 where M05 is the R- $\theta$  iteration reduction factor specified in the input. For each inner iteration, the number of R- $\theta$  iterations is at least equal to KM and is not larger than IIPGxKM/M05, where KM is the number of axial mesh points.

In each horizontal plane, the fluxes are computed using line inversion. That is, the fluxes on each rotational (or radial) line are assumed to be unknown, and all other fluxes coupled to the unknown fluxes are assumed to be known (using updated values). This procedure leads to a system of coupled linear nonhomogeneous equations which are solved by Gaussian elimination and back-substitution.

The direction (rotational or radial) in which the line inversion is performed depends on the geometry and on the boundary conditions at the rotational boundaries. Line inversion is performed in the rotational direction for all  $R-\theta-Z$  problems and for X-Y-Z problems with periodic boundary conditions. For other X-Y-Z problems, the line inversion is performed in the direction containing the most mesh points.

After each horizontal flux profile,  $\phi_i$ , for the ith R- $\theta$  iteration is calculated, a new profile,  $\phi'_i$ , is computed by magnifying the difference between  $\phi_i$  and  $\phi_{i-1}$  according to

$$\phi'_{i} = \phi_{i} + \beta(\phi_{i} - \phi_{i-1}) , \qquad (16)$$

where  $\beta$  is the overrelaxation factor specified in the input. The code automatically reduces  $\beta - 1$  by a factor of 1.1 when  $|1 - \lambda| < 10$ EPS and also computes the fission source overrelaxation factor,  $\beta$ ', from the expression

$$\beta' = 1.0 + 0.6(\beta - 1) \quad . \tag{17}$$

Tight mesh spacing in the dimension perpendicular to line inversion can cause excessive running time. Thus, if tight mesh spacing is used, it should be in the dimension containing the most mesh intervals. For the same reason, the dummy dimension in two-dimensional problems should contain large mesh intervals (at least three intervals are required). IV. IMPLICIT EIGENVALUE SEARCHES

The 3DDT code permits implicit eigenvalue searches for time absorption (alpha), material composition (concentration), or system dimensions (delta). In contrast to a k<sub>eff</sub> calculation, the fission spectrum is not changed. Instead, the desired parameter is changed to make  $\lambda$  approach unity. In general terms, a converged value of  $\lambda$  is obtained by a sequence of outer iterations for the initial system configuration. Then the system is altered, by changing the desired parameter by an amount specified in the input, and another sequence of outer iterations is performed to obtain a second converged  $\lambda$ . Subsequent parameter changes are determined by either linear interpolation or by parabolic interpolation modified by precautionary safeguards. Eigenvalue search problems are thus basically a sequence of  $k_{eff}$  calculations for several perturbations of the original system.

The outer iteration process for the original system (and corresponding initial eigenvalue guess) is continued until  $|\lambda_j - \lambda_{j-1}| < EPS$ , where j is the outer iteration index and EPS is the lambda convergence criterion specified in the input. The original eigenvalue guess, EV, is then changed to EV-EVM if  $\lambda < 1$  and to EV+EVM if  $\lambda > 1$ . EVM is an input parameter (eigenvalue modifier) for changing the implicit parameter which is the object of the search, and an intelligent use of EVM can materially shorten computation time.

After converging on the first value of  $\lambda$  to a precision of  $|\lambda_j - \lambda_{j-1}| < EPS$  for two successive outer iterations, subsequent values are converged to  $|\lambda_j - \lambda_{j-1}| < EPSA$ , where EPSA is the eigenvalue search convergence criterion specified in the input. In most computations, EPSA can be larger than EPS, saving time by requiring fewer outer iterations without sacrificing final accuracy. When the second converged value of  $\lambda$  is obtained for the second value of EV, a straight-line interpolation is used to modify EV, provided that XLAL  $\leq |1 - \lambda| \leq$  XLAH. XLAL and XLAH are low and high limits on  $|1 - \lambda|$  specified in the input and used to stabilize the search. If  $|1 - \lambda| >$  XLAH, the straight-line interpolation proceeds as though  $|1 - \lambda|$  were equal to XLAH. If  $|1 - \lambda| <$  XLAL, the search continues to completion using a straight-line interpolation with a fixed slope to prevent errors due to subtraction of nearly equal quantities. In all cases where straight-line interpolation is used, the slope of the line can be adjusted to either over- or underpredict by using the parameter oscillation damper, POD, specified in the input.

After a third converged value of  $\lambda$  (corresponding to the third value of EV) is obtained, and if  $|1 - \lambda| \ge XLAL$ , parabolic interpolation is used to obtain the next value of EV. The root of the parabola closest to the previous value of EV is taken as EV. Let  $\lambda^p$  and  $\lambda^{pp}$  be the converged lambdas corresponding to EV<sup>P</sup> and EV<sup>PP</sup>, where EV<sup>P</sup> is the previous eigenvalue and EV<sup>PP</sup> is two eigenvalues back. If the sign of  $\lambda^{pp}$  - 1 is different from the sign of  $\lambda^p$  - 1, the new value of EV must lie in the extreme range of the previous values E<sup>P</sup> and E<sup>PP</sup>. If it does not, EV is set equal to  $0.5(EV^{P} + EV^{PP})$ . In the parabolic search, a slope is computed for the switch to the straight-line method if  $|1 - \lambda| < \lambda$ XLAL. Also, if the two roots of the parabola are imaginary, the straight-line method is used. Final convergence of the search is reached when both  $|1 - \lambda| < EPS$  and  $|\lambda_i - \lambda_{i-1}| < EPS$ .

Alpha Searches. Consider the time-dependent form of Eq. (1),

$$\frac{1}{v_g} \frac{\partial \phi_g(\vec{r}, t)}{\partial t} = \nabla \cdot D_g \nabla \phi_g(\vec{r}, t) - \Sigma_g^r \phi_g(\vec{r}, t) + S_g(\vec{r}, t) .$$
(18)

If we assume that

$$\phi_g(\vec{r},t) = \phi_g(\vec{r}) e^{\alpha t} , \qquad (19)$$

Eq. (16) becomes

$$\nabla \cdot D_{g} \nabla \phi_{g}(\vec{r}) - \left(\sum_{g}^{r} + \frac{\alpha}{v_{g}}\right) \phi_{g}(\vec{r}) + S_{g}(\vec{r}) = 0.$$
(20)

The parameter  $\alpha$  as defined in Eq. (19) is computed as the eigenvalue EV in an alpha calculation. Because the term  $\alpha/\nu_g$  in Eq. (20) appears as an additional absorption term, alpha calculations are referred to as time absorption calculations.

<u>Concentration Searches</u>. Any number of materials can simultaneously be added, depleted, or interchanged in any number of zones during a concentration search. The basic format for specifying concentration searches is illustrated by a simple example. Suppose that mixture 15 is initially composed of materials 1, 2, 3, and 4 with atom densities, respectively, of  $N_1$ ,  $N_2$ ,  $N_3$ , and  $N_4$ . Now suppose that we wish to vary the atom densities of materials 1 and 2 by the same factor in order to make the system critical, while at the same time keeping the atom densities of materials 3 and 4 constant. The IO, II, and I2 tables (see input instructions, Appendix A) for mixture 15 would then be specified as shown in the following tabulation.

Mix	Material	Atom
Number (10)	Number (I1)	Density (I2)
15	0	0
15	1	N <sub>1</sub>
15	2	NZ
15	15	0
15	3	N3
15	4	N <sub>4</sub>

In the first row, the 0 entry in the II table instructs the code to clear a storage area for mixture 15. In the second and third rows, the entries in the II table cause materials 1 and 2 to be added to the current contents of mixture 15 with atom densities, respectively, of  $N_1$  and  $N_2$ . Because the entry in the II table in the fourth row is the same as the mixture number, the current contents of mixture 15 is multiplied by the eigenvalue. Finally, the entries in the II table in the last two rows instruct the code to add materials 3 and 4, with atom densities, respectively, of  $N_3$  and  $N_4$ , to the current contents of mixture 15. These instructions are summarized by the expression.

$$\Sigma_{15} = EV(N_1\sigma_1 + N_2\sigma_2) + N_3\sigma_3 + N_4\sigma_4 , \quad (21)$$

where

 $\Sigma_{15}$  = macroscopic cross section for mixture 15,  $\sigma_i$  = microscopic cross section for material i,  $N_i$  = atom density of material i, and EV = eigenvalue.

From the simple example above as a basis, complex concentration search problems can be constructed; the possibilities are limited only by the ingenuity of the user.

<u>Dimension Searches</u>. In delta calculations. the code searches on reactor dimensions by varying the size of each mesh interval according to

$$\Delta r^{j} = \Delta r_{o}^{j} [1 + (\text{mesh modifier})^{j} EV] . \qquad (22)$$

In Eq. (22),  $\Delta r_o^j$  is the initial mesh spacing for the jth interval and EV is the eigenvalue. Different mesh modifiers can be specified for each radial, rotational, and axial mesh interval. This allows great versatility in the manner in which the system can be changed to achieve criticality.

In the concentration and delta options, either  $k_{eff}$  or alpha can be used as parametric eigenvalues. That is, concentration or dimension searches can be made so that the final system has a specified reactivity or a specified alpha. If no parametric eigenvalue is specified, it is assumed that the final system is to be critical.

V. BURNUP MODEL

The burnup equation for each material in each zone has the form

$$dN_{i}/dt = \lambda_{k}N_{k} + \bar{\phi} \sum_{j=1}^{2} \bar{\sigma}_{j}C_{j}N_{j} + \bar{\phi} \sum_{m=1}^{7} \bar{\sigma}_{m}f_{m}$$
$$- \lambda_{i}N_{i} - \bar{\phi}\bar{\sigma}_{i}N_{i} , \qquad (23)$$

where

- $N_1 =$ atom density of nuclide i in the zone,
- $\lambda_i$  = decay constant for nuclide i,
- $\sigma_1^{\tilde{a}}$  = spectrum-averaged microscopic absorption cross section for nuclide 1,
- $\bar{\sigma}_{j}^{c}$  = spectrum-averaged microscopic capture cross section for nuclide j,
- $\overline{\sigma}_{m}^{f}$  = spectrum-averaged microscopic fission cross section for nuclide m, and
- $\overline{\phi}$  = total zone-averaged flux.

The first term on the right-hand side of Eq. (23) gives the source for nuclide i from decay of nuclide k. The next two terms provide for two capture and seven fission sources. Since fission products are usually included as a single nuclide, the yield of nuclide i from the fission of nuclide m has been assumed to be unity in Eq. (23). Modifications to the code required to allow fission yields for individual fission products are straightforward. The last two terms in Eq. (23) provide for losses of nuclide i by decay and absorption.

The zone-averaged flux and cross sections appearing in Eq. (23) are computed and printed before each burnup interval and, along with the total reactor power and flux profiles, are held constant during the burnup interval. Each burnup interval is arbitrarily subdivided into ten smaller intervals of equal duration. Equation (23) is then solved as a march-out problem using the smaller time intervals.

If Eq. (23) is rewritten in the form

$$d\vec{N}/dt = \vec{f}(\vec{N},t) , \qquad (24)$$

the march-out algorithm can be written as

$$\vec{N}_{J} = \vec{N}_{J-1} + \frac{\delta_{E}}{2} (\vec{f}_{J-1} + \vec{f}_{J})$$
 (25)

where J is the time index and  $\delta t$  is the subdivided time interval. Because  $\vec{N}_J$  must be known in order to compute  $\vec{f}_J$ , Eq. (25) must be solved by an iterative procedure at each time step J. The algorithm used is

$$\vec{N}_{J}^{\nu} = \vec{N}_{J-1} + \frac{\delta t}{2} \left( \vec{f}_{J-1} + \vec{f}_{J}^{\nu-1} \right) , \qquad (26)$$

where v is the iteration index. It is clear from the mathematical model that short burnup intervals should be used if rapid variations in isotopic concentrations or flux profiles are expected.

Although the eigenvalue and material densities for the depleted system are computed and printed after each burnup interval, the zone-averaged flux, cross sections, and reaction rates are not. However, they can be obtained by specifying a final burnup interval of minuscule length. The same stratagem can be used in nonburnup problems to obtain a printout of these quantities.

REFERENCES

- W. W. Little, Jr. and R. W. Hardie, "2DB User's Manual," BNWL-831, Battelle Northwest Laboratory (1968).
- K. D. Lathrop, "DTF-IV, a FORTRAN-IV Program for Solving the Multigroup Transport Equation with Anisotropic Scattering," LA-3373, Los Alamos Scientific Laboratory (1965).
- 2DF is an unpublished, two-dimensional version of the DTF-IV code developed at the Los Alamos Scientific Laboratory.

#### APPENDIX A

#### INPUT INSTRUCTIONS

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In the following pages, input specifications	<u>11</u>	Option
for 3DDT are listed and described in exactly the	0 (or blank)	Code reads data field and no special
order in which they are required by the code. The		action is taken.
input is divided into four categories:	1	Repeat the number in the data field I2 times.
1. a title card,	2	Place I2 linear interpolates between
<ol> <li>input control integers and input control float- ing point numbers intermixed on Cards 2 through 9,</li> </ol>	-	the number in this data field and the number in the following data field (not allowed for integers).
<ol> <li>problem-dependent data (number of cards is variable), and</li> </ol>	3	Terminate reading of the data block. A 3 must follow the last data field
4. burnup data (number of cards is variable).		of each data block.
The formats developed for reading problem-	4	Fill the remainder of the data block with the number in the data field.
dependent data in the DTF-IV code are also used in	5	Reneat the number in the data field
3DDT. With the exception of the cross sections,	5	10 x 12 times.
all problem-dependent data are read by generalized	9	Indicates the end of data to be read
input subroutines in one of two formats:		from the card and code skips to the
6(I1,I2,E9.4) for reading floating point numbers		next card. This option allows inser- tions of data within a block without
and 6(I1,I2,I9) for reading integers. In these for-		repunching the entire block.
mats, the integer in the first column (Il) indi-	Both the	2DB and the DTF-IV cross-section for-
cates the option (described below), the integer in	mats are allow	ved in 3DDT. Algorithms for mixing
the next two columns (I2) indicates the number of	cross sections	and for criticality searches in k eff,
times the option is to be applied, and the number	alpha, concent	ration, and delta calculations are the
in the last nine columns (I9 or E9.4) is the data	same as those	used in the DTF-IV and 2DB codes.
associated with the field. The options for Il are:		

Word	Variable	Format	Comments
Card 1	(required)		
A11	ĺD	12A6	Problem identification card which may contain any alpha- numeric characters desired.
CONTROL	DATA (requir	ed)	
Card 2			
1	A01	112	Cross-section format indicator $(1/2 = BNWL/LASL)$ . See explanatory notes.
2	A02	112	Theory (0/1 = regular/adjoint). Flux guess for adjoint prob- lem must be input in reverse order in terms of groups. Adjoint fluxes are also printed in reverse order. The neutron balance tables, power density, and fission source rate in adjoint cal- culations do not have a direct physical interpretation.
3	IGE	112	Geometry (0/1 = X-Y-Z/R- $\theta$ -Z). Input for $\theta$ is in revolutions, e.g., $\pi$ radians = 0.5 revolution.
4	IZM	112	Number of material zones.
5	104	112	Eigenvalue type $(1/2/3/4 = k/\alpha/C/\delta)$ .
6	M07	112	Initial flux guess (0/1,2,3/4/5 = flat/cards/tape/sine-cosine). See input block NO. If M07=4, the flux tape must be assigned to file TAPE14.

Word	Variable	Format	Comments
Card 3			
1	IM	112	Number of radial (or X) intervals ( $\geq$ 3).
2	JM	112	Number of rotational (or Y) intervals (≥3).
3	КM	112	Number of axial (Z) intervals (≥3).
4	EV	E12.4	Initial eigenvalue guess (typical values are 1.0 for C calcula- tions and 0.0 for all the rest).
5	EVM	E12.4	Eigenvalue modifier used in search calculations. Typical values are $\pm 1.0$ for $\alpha$ , $\pm 0.1$ for C, and $\pm 0.05$ for $\delta$ . Value of EVM should reduce reactivity.
6	EPS	E12.4	Lambda convergence criterion, i.e., convergence criterion on the total fission source. Typical values are 0.00001 $\leq$ EPS $\leq$ 0.0001.
Card 4			
1	B01	112	Left boundary condition (0/1 = vacuum/reflective).
2	B02	112	Right boundary condition (see BOl).
3	B03	112	Front boundary condition $(0/1/2 = vacuum/reflective/periodic)$ .
4	в04	L12	Back boundary condition (see BO3).
5	B05	112	Bottom boundary condition (0/1 = vacuum/reflective).
6	B06	112	Top boundary condition (see B05).
Card 5			
1	MT	112	Total number of materials including mixtures. Mixtures must be assigned identification numbers (see IO block) between MCR+1 and MT.
2	M01	112	Number of mixture specifications (>0 for burnup calculations). This is the length of the blocks (see IO, II, and I2) which specify how mixtures are to be formed from the input cross sections.
3	MCR	112	Number of cross-section sets input from cards ( $\geq$ 1). The sets are assigned identification numbers by the code from 1 to MCR in the order in which they are read.
4	IZ	112	Number of radial (or X) mesh interval modifiers (≥1 for δ option, O otherwise). See R3 block.
5	JZ	112	Number of rotational (or Y) mesh interval modifiers (≥1 for δ option, O otherwise). See TA3 block.
6	КZ	112	Number of axial (Z) mesh interval modifiers (≥1 for δ option, O otherwise). See Z3 block.
Card 6			
1	S02	112	Parametric eigenvalue type (0/1/2 = none/k/ $\alpha$ ). Use with C and $\delta$ options only.
2	S03	E12.4	Parametric eigenvalue (leave blank or zero if SO2 = 0).
3	IGM	112	Number of energy groups $(\geq 1)$ .
4	IHT	112	Position of sigma transport in cross-section table.
5	IHS	112	Position of sigma self-scatter in cross-section table.
6	ITI.	·T12	Cross-section table length.

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Word	Variable	Format	Comments
Card 7			
1	S01	E12.4	Neutron source rate (SO1 > 0) or power in MWT (SO1 < 0). Used to normalize the fluxes.
2	M05	I12	R- $\theta$ iteration reduction factor. The maximum number of R- $\theta$ (or X-Y) mesh sweeps per inner (Z) iteration is set equal to S04/M05 or G07/M05 as appropriate. Suggested value M05 = 2.
3	M06	112	Number of R- $\theta$ (or X-Y) planes with unique zone numbers (1 $\leq$ MO6 $\leq$ KM).
4	S04	112	Inner (Z) iteration maximum per group for $ 1 - \lambda  > 10 \times PS$ (suggested value 5 $\leq$ SO4 $\leq$ 10).
5	D05	112	Maximum number of outer iterations (if running time limita- tion is object, use ITIMOF).
6	G07	112	Maximum number of inner (Z) iterations per group when $ 1 - \lambda  < 10 \times PS$ (suggested value 10 $\leq$ G07 $\leq$ 20).
Card 8			
1	G05	E12.4	Not used at presentleave blank or zero.
2	G06	E12.4	Inner iteration pointwise flux convergence criterion. Sug- gested value EPS $\leq$ GO6 $\leq$ 10xEPS. GO6 cannot be zero.
3	LAL	E12.4	Lower limit on $ 1 - \lambda $ used in search options only. Suggested value: 0.005.
4	LAH	E12.4	Upper limit on $\left 1-\lambda\right $ used in search options only. Suggested value: 0.5.
5	POD	E12.4	Parameter oscillation damper used in search options only. Sug-gested values: 0.5 for $\alpha$ calculation, 1.0 otherwise.
6	EPSA	E12.4	Eigenvalue search convergence criterion. Used only in search calculations. Suggested value EPS $\leq$ EPSA $\leq$ 10xEPS.
Card 9			
1	IPFLX	112	Punch flux dump (0/1 = yes/no). Not recommended for IMxJMxKMxIGM > 6000, instead use IDMTPS. If a dump is called for in a burnup calculation, only the fluxes computed for the undepleted system are dumped. Adjoint fluxes are punched in reverse order in terms of groups.
2	IPCUR	112	Punch atom densities after each burnup interval $(0/1 = no/yes)$ .
3	IDMTPS	112	Write flux dump on magnetic tape (0/1 = yes/no). If yes, must assign a magnetic tape to file TAPE16. Adjoint fluxes are written in reverse order in terms of groups. In burnup calcula- tions, only the fluxes computed for the undepleted system are dumped.
4	ITIMOF	112	Time limit in seconds. For long running problems, this allows a flux dump to be taken before CP time expires.
5	ORF	E12.4	Overrelaxation factor. Suggested value: 1.4.
6	NESFL	112	ECS field length in octal thousands. Must be the same as that requested on the job card.

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#### PROBLEM-DEPENDENT DATA

Block Name	Format	Number of Entries		Comments
CO (required)	6E12.5	MCR	Cross-section se number of sets. the heading card AA(9) (format A6 isotope identifi weight (amu) of laneous addition card is followed (in barns) for e highest energy g for each group s the cross sectio A01=2 and IHT=3, is followed by A the value of v f various interact table is discuss	ts on cards where MCR is the The first card in each set is containing HOLN, ATW, ALAM, 2E6.2, 9A6) where HOLN = cation name, ATW = atomic isotope, and AA(9) = miscel- al identification. The heading by the cross-section tables each group beginning with the roup. For AOl=1, the table tarts on a new card. For AOl=2, ns are a continuous block. If the last set of cross sections NU(I), I=1,IGM where ANU(I) is or group I. The order of the fion cross sections within the sed in the explanatory notes.
NO (optional)	6(I1,I2,E9.4)	Varies	Flux input. The on the option ch as follows:	e length of this block depends nosen in parameter MO7 (card 2)
			M07 Option	Length
			0 1 2 3 4 5 See explanatory flux input.	No input required. IM+JM+KM+IGM IGM(IM+JM+KM) IGMxIMxJMxKM Same as M07=3 but flux guess is taken from tape. The mag- netic tape containing the flux guess must be assigned to file TAPE14. Sine or cosine guess computed internally by subroutine SINUS. Card input not required at this point but see block EF. notes for further details on
RO (required)	6(I1,I2,E9.4)	IM+1	Radial (X) mesh ordered by magni for $R-\theta-Z$ proble	in cm. The entries must be Itude and must begin with 0.0 ems
TAO (required)	6(I1,I2,E9.4)	JM+1	Rotational (Y) m ordered by magni	mesh in revolutions (or cm) itude.
ZO (required)	6(I1,I2,E9.4)	KM+1	Axial (Z) mesh f	In cm ordered by magnitude.
EF (optional)	6(I1,I2,E9.4)	IGM	Flux energy spec See explanatory	ctrum to be read only if MO7=5. notes.
MO (required)	6(11,12,19)	MO6xIMxJM	Zone number block $R-\theta$ (X-Y) planese entries. The zon in the radial (X) order from left tional (Y) inter- rotational inter- zone numbers for specified.	cks for each of the MO6 unique s. Each block contains IMxJM one numbers for all intervals () direction must be given in to right for the first rota- rval, then for the second rval, and so forth until the r the entire plane have been
IDZNO (required)	6(I1,I2,I9)	КM	Identifies zone each axial inten interval.	number plane to be used for rval starting with the bottom
M2 (required)	6(11,12,19)	IZM	Material numbers	s by zone starting with zone l.
K7 (required)	6(I1,I2,E9.4)	IGM	Fission spectrum with the highest	m ordered by group beginning t energy group.

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#### PROBLEM-DEPENDENT DATA (continued)

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Block		Number of	
Name	Format	Entries	Comments
V7 (required)	6(I1,I2,E9.4)	IGM	Velocities (cm/sec) for each group starting with the highest energy group.
IO (optional)	6(I1,I2,I9)	моі	Mixture numbers labeling the cross-section mixture blocks. The length of each mixture block is n+1 where n is the number of isotopes in the mix. No input is required if MO1=0.
Il (optional)	6(I1,I2,I9)	M01	Mixture commands indicating which isotopes are to be used to make each mixture. The first entry for each mix is 0. Once a mixture has been defined, it may be used as a component of another mixture. However, the material inven- tory tables, which are part of the printed output, are inapplicable for a mixture speci- fication more complex than a mix in a mix (e.g., a mix in a mix in a mix). An isotope cannot be specified more than once in the same mix in burnup calculations. If specified more than once in other calculations, the printed inven- tory table will be incorrect. The procedure used in concentration search problems is dis- cussed in Sec. IV of this report. No input is required for Il if MO1=0.
I2 (optional)	6(I1,I2,E9.4)	M01	Atomic densities $(x \ 10^{24} \ since \ input \ cross$ sections are assumed to be in barns) of the various components of each mixture. The first entry for each mixture must be 0.0, corres- ponding to the 0 in the Il table, followed by the density for each material in the order they are given in the Il table. No input re- quired if M01=0. See Sec. IV of this report for procedure used in concentration searches.
R2 (optional)	6(I1,I2,I9)	M	Radial numbers. These are read and used only in $\delta$ calculations. They indicate which radial mesh interval modifier (block R3) is to be used for each radial (or X) interval.
R3 (optional)	6(I1,I2,E9.4)	IZ	Radial mesh interval modifiers. These are the factors used with the eigenvalue to achieve the increase or decrease in size of the radial (or X) mesh intervals. A different modifier may be specified for each mesh interval. This block is read only in $\delta$ calculations.
TA2 (optional)	6(I1,I2,I9)	ML	Rotational numbers. These are used in the same manner as the R2 block except in the $\theta$ (or Y) direction.
TA3 (optional)	6(11,12,E9.4)	JZ	Rotational mesh interval modifiers. These are used in the same manner as the R3 block except in the $\theta$ (or Y) direction.
Z2 (optional)	6(11,12,19)	KM	Axial numbers. These are used in the same manner as the R2 block except in the axial (Z) direction.
Z3 (optional)	6(I1,I2,E9.4)	KZ	Axial mesh interval modifiers. These are used in the same manner as the R3 block except in the axial (Z) direction.

BURNUP DATA

Word	Variable	Format	Comments
Card 1	(required)		••
1	NCON	16	<ul> <li>Burnup control:</li> <li>= 0: end of problem, read input data for next case. The remaining words on the card are ignored. This card, with NCON=0, must be the last card in the data deck for each problem.</li> <li>= N: read burnup parameters for N burnable isotopes and take time step of DELT. Card 2 below is read only if this option is used.</li> <li>&lt; 0: take time step of DELT using burnup parameters from previous time step. If this option is used, Card 2 below is not required.</li> </ul>
2	NPRT	16	<pre>Print control: = 0: partial print. Fluxes, fission densities, and fission neutron source rates are not printed for the depleted system. = 1: full print. Quantities normally printed after an eigenvalue calculation are printed for the depleted system.</pre>
3	DELT	E12.0	Length (days) of time step (burnup interval).
Card 2	(optional:	repeat for	all burnable isotopes, i.e., N=1,NCON)
1	MATN(N)	16	Material number of the burnable isotope.
2	NBR (N)	16	Control for breeding ratio calculation: = 0: no effect = 1: fertile isotope = 2: fissile isotope
3	LD(N)	16	Decay source control: = 0: no decay source = I: decay source from burnable isotope I
4	LCN (N,1)	16	Capture source #1: = 0: no capture source = I: capture source from burnable isotope I
5	LCN (N,2)	16	Capture source #2: Same options as LCN(N,1).
6 •	LFN(N,1)	16	<pre>Fission source #1: = 0:  no fission source = I: fission source from burnable isotope I</pre>
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12	LFN(N,7)	16	Fission source #7: Same options as LFN(N,1).

#### EXPLANATORY NOTES

<u>Cross Sections</u>. If control parameter A01=1. the cross-section tables are read in the BNWL format. In this format, the cross sections are assumed to be in the order  $\sigma_g^f$ ,  $\sigma_g^a$ ,  $v\sigma_g^f$ ,  $\sigma_{g+g}^t$ ,  $\sigma_{g-1+g}$ ,  $\sigma_{g-2+g}^t$ , ..., and the table for each group starts on a new card. In the LASL format (A01=2), the cross sections are either in the order  $\sigma_g^c$ ,  $\sigma_g^f$ ,  $\sigma_g^s$ ,  $\sigma_g^a$ ,  $v_{g}^f$ ,  $\sigma_{g+g}^{tr}$ ,  $\sigma_{g-1+g}$ ,  $\sigma_{g-2+g}^c$ , ..., or in the order  $\sigma_g^a$ ,  $v\sigma_g^f$ ,  $\sigma_g^{tr}$ ,  $\sigma_{g+g}^{-g+g}$ ,  $\sigma_{g-1+g}^{-g+g}$ ,  $\sigma_{g-2+g}^{-g+g}$ , ..., and the cross sections are read in a continuous block. The code checks the cross-section tables for internal consistency. That is, the transport cross section computed from

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$$\sigma_{g}^{tr} = \sigma_{g}^{a} + \sum_{g'} \sigma_{g \neq g'}$$
 (A-1)

must agree with the  $\sigma_g^{tr}$  specified in the table. If they differ by more than EPS in relative magnitude, either the problem is terminated or a warning is printed, depending on the magnitude of the discrepancy. <u>Flux Guess</u>. Let  $\phi(G,I,J,K)$  denote the flux for group G at radial (or X) mesh point I, rotational (or Y) mesh point J, and axial (Z) mesh point K. If M07=0, the guess

$$\phi(G, I, J, K) = 1.0$$
 (A-2)

is used for all groups and all mesh points. This option is recommended only for cell problems in which all the boundaries are reflected and the energy spectrum is not known. For other problems, option M07=5 should be used if a flux guess from other sources is not available.

If M07=1, the flux guess is synthesized from

$$\phi(G,I,J,K) = X(I)Y(J)Z(K)H(G), \qquad (A-3)$$

where X(I), Y(J), Z(K), and H(G) are, respectively, the radial, rotational, axial, and energy distributions. These distributions can be obtained from one-dimensional calculations using, for example, the DTF-IV code. This flux guess is effective for problems which are nearly separable in both space and energy. The one-dimensional distributions are read from cards as separate blocks in the order X(I), I=1,IM; Y(J), J=1,JM; Z(K), K=1,KM; and H(G), G=1, IGM.

If M07=2, the flux guess is synthesized from

$$\phi(G,I,J,K) = X(G,I)Y(G,J)Z(G,K) , \qquad (A-4)$$

where X(G,I), Y(G,J), and Z(G,K) are, respectively, the group-dependent radial, rotational, and axial distributions. As before, these distributions can be obtained from one-dimensional calculations. This flux guess is effective for problems which are nearly separable in space but not in energy. The distributions are read from cards as separate blocks in the order X(G,I), I=1,IM; Y(G,J), J=1,JM; and Z(G,K), K=1,KM for each group starting with G=1 for regular problems and G=IGM for adjoint problems.

If M07=3, a full flux guess  $\phi(G,I,J,K)$  is read from cards. The fluxes are read in KM blocks of length IMxJM in the order

for each group starting with G=1 in regular problems or G=IGM in adjoint problems. Normally, this guess will be a card flux dump from a previous run and the blocks will already be in the correct order.

The flux guess for M07=4 is the same as for M07=3 except that the fluxes are on magnetic tape rather than on cards. This guess is normally a tape flux dump from a previous run.

If M07=5, the flux guess is synthesized in the same manner as the M07=1 option, Eq. (A-3). However, the one-dimensional spatial distributions X(I), Y(J), and Z(K) are computed internally by subroutine SINUS from sine, cosine, or flat functions depending on the boundary conditions. The energy distribution H(G) (block EF in the input specifications) is read from cards.

#### APPENDIX B

#### STORAGE REQUIREMENTS

The variable-dimensioned arrays in 3DDT are stored in three levels. Four-dimensional arrays, e.g.,  $\phi(x,y,z,g)$ , are stored on disk; threedimensional arrays, e.g.,  $\phi(x,y,z)$  for a particular energy group, are stored in the ECS; and twodimensional arrays, e.g.,  $\phi(x,y)$  for a particular energy group and axial position, are stored in a container block in the fast central memory. Thus, central memory storage requirements are insensitive to the number of energy groups and the number of axial mesh points. Three levels of storage are

required because the combined storage capacities of the fast core and the ECS are not large enough for most problems.

On the CDC 6600 computer, disk files are treated the same as tape files. That is, they may be written, rewound; read, and backspaced by the standard FORTRAN I/O statements. In 3DDT, disk files are used to store fluxes (file NFLUX1), flux constants (file NCXS1), and cross sections (file NCR1). An additional disk file, NSCRAT, is used as a scratch file. Both random access and sequential file access can be used to transfer arrays between fast core and the ECS. Random ECS access permits the manipulation of data arrays through high-speed block transfers. Sequential ECS file access permits tape file simulation to be employed using ECS as the storage medium. For each read or write statement, the time required to transfer n words between fast core and ECS is given approximately by

$$t = a + n/b, \qquad (B-1)$$

where a = access time in seconds and b = transmission rate in words/sec. With random access, a  $\approx$  $1.5 \times 10^{-5}$  and b  $\approx$  6.7 x  $10^{6}$ ; whereas with sequential access, a  $\approx$  8.0 x  $10^{-3}$  and b  $\approx$  6.7 x  $10^{6}$ . Although the transmission rates are the same for the random and sequential modes, the much larger access time for the sequential mode greatly reduces the effective transfer rate.

The effective transfer rate, R, is given by

$$R = n/t = n/(a + n/b)$$
. (B-2)

Equation (B-2) shows that R increases linearly with n for small n and approaches the transmission rate, b, for large n. Effective transfer rates computed from Eq. (B-2) for the two ECS modes are tabulated below for various values of n.

	Effective Transi	fer Rate (words/sec)
n	Random Access	Sequential Access
10	$6.1 \times 10^{5}$	$1.2 \times 10^{3}$
100	$3.3 \times 10^{6}$	$1.2 \times 10^{4}$
1000	$6.1 \times 10^{6}$	$1.2 \times 10^{2}$
10000	6.6 x 10 <sup>6</sup>	$1.1 \times 10^{6}$

Since the arrays transferred between fast core and ECS are of the order of IMxJM, where IM = number of radial (or X) mesh points and JM = number of rotational (or Y) mesh points, n is of the order of 1000 in most problems. The effective transfer rate for n = 1000 using random access is 50 times larger than that using sequential access. For this reason, random access is used in the CDC 6600 version of 3DDT.

The use of sequential file access would make 3DDT compatible with computing systems for which ECS is not available. Thus, it is expected that a version using sequential ECS file access will be prepared later. Disk storage could be used in place of the ECS, but only at considerable cost in computing time. The effective transfer rate between fast core and disk is at least two orders of magnitude smaller than that between fast core and the ECS using random access. The use of magnetic tape storage is even less desirable, because effective transfer rates with magnetic tape are several times smaller than with disk.

The three-tier storage arrangement is well suited to the iteration scheme used in 3DDT. During an outer iteration, the three-dimensional arrays for the current group are extracted from the disk and stored in the ECS before entering the inner iteration loop. In the inner iteration loop, the twodimensional arrays for the current axial position are transferred from the ECS to central memory. Since the two-dimensional arrays must be transferred from the ECS to central memory every time the axial position changes, and since the axial mesh may be swept many times for every outer iteration, most of the data transmission is between ECS and central memory. Thus, disk access time and transmission rates are not as critical as those for the ECS.

The variable-dimensioned arrays stored in the container block A in central memory require N  $_{\rm CM}$  storage locations, where

N<sub>Cm</sub> = 17xIMxJM + ITLxMT + 15xML + 2xJMxKM + 2xIMxKM + 4xM01 + 6xIGM + 8xIZM + 2(IP+JP+KP) + 2(IM+JM+KM) + 4xML + 12xIGP + M06 + 11xKM + 2xMAX(IM,JM) + T06(IM+JM+KM+IZ+JZ+KZ). (B-3)

Extended core storage and disk storage requirements are, respectively,

and

In Eqs. (B-3) through B-5),

- IM = number of radial (or X) mesh intervals,
- JM = number of rotational (or Y) mesh intervals,

(B-4)

- KM = number of axial (Z) mesh intervals,
- ITL = cross-section table length,
- MT = total number of materials including mixes,
- ML = number of input material cross sections,
- MO1 = number of mixture specifications,
- IGM = number of energy groups,

- IP = IM+1,
- JP = JM+1,
- KP = KM+1,
- IGP = IGM+1,
- IZM = number of material zones,
- IZ = number of radial (X) zones to be modified (delta option),
- JZ = number of rotational (Y) zones to be modified (delta option),
- KZ = number of axial (Z) zones to be modified (delta option),
- M06 = number of R- $\theta$  (X-Y) planes with unique zone numbers, and

T06 = 1 for delta calculation, = 0 otherwise.

For a fairly large problem, in which IM=JM=KM=M06=IZ=JZ=KZ=30, IGM = 16, ITL = 12, MT = 100, ML = 20, M01 = 200, T06 = 1, and IZM = 80, the storage requirements for the variable arrays are

$$N_{cm} = 23,186,$$
  
 $N_{ecs} = 340,200,$   
 $N_{disk} = 2,611,200.$ 

On a 65K machine, the dimension of the container block A can be as large as 25,000. The maximum ECS memory available to a single user is currently about 800,000 and each machine has access to two disks of about 6,500,000 words each. Thus, the large problem specified above can be accommodated on a 65K machine and much larger problems can be accommodated on a 130K machine. Because of the manner in which arrays are stored in central memory, very large two-dimensional (R-Z or X-Z) problems can be accommodated on a 65K machine. For example, consider an R-Z problem with IM=KM=M06=IZ=KZ=75, IGM = 25, ITL = 15, MT = 200, ML = 40, M01 = 400, IZM = 160, and TO6 = 1. In the dummy direction,  $\theta$ , JM = 3 (at least three mesh points are required in each dimension) and JZ = 0. For this case, the storage requirements are

 $N_{cm} = 24,598,$  $N_{ecs} = 337,500,$  $N_{disk} = 2,606,250.$ 

#### APPENDIX C

#### SIMPLIFIED LOGICAL FLOW CHART

A simp	lified logical flow chart for 3DDT is	Subroutine	Description
shown in Fig routines are description Subroutine	g. C-l. Only external (nonsystem) sub- e included in the flow chart and a brief of these subroutines is given below. <u>Description</u>	S860	Reads cross sections from cards and checks them for consistency, performs adjoint reversals of the cross sections if required, and writes the cross- section tape. \$860 is called by INP and calls ERR02.
3DDT	Main programcontrols the overall flow of the problem. 3DDT calls INP, INIT, FISCAL, S8830, ERRO2, CONSTS, OUTER,	S862	Reads input flux guess and prepares a flux tape. It is called by INP and calls REAG2.
INP	MARCH. Controls the reading and printing of all	SINUS	Subroutine to calculate flux input guess. It is called by INP and calls REAG2.
inpu and var cal REA( POI	Input data other than the burnup data and computes program constants and variable-dimension pointers. INP is called by 3DDT and calls S860, S862, REAG2, REAI2, SINUS, MAPR, ERRO2, and	REAG2	Intermediate subroutine used to read floating-point data. REAG2 is called by INP, SINUS, and S862 and calls LOAD.
ERRO 2	This subroutine prints error messages and is called by 3DDT, INP, S860, INIT, and CNNP	REAI2	Intermediate subroutine used to read integer data. REAI2 is called by INP and calls LOAD.
POINTR	Calculates pointers for arrays in ECS. POINTR is called by INP.	LOAD	Reads either integer or floating-point data in the DTF-IV format. LOAD is called by REAG2 and REAI2.
RETRVR	Retrieves integer data from the con- tainer block A. RETRVR is called by INP.	MAPR	Produces a picture plot by zone and material. MAPR is called by INP.



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Fig. C-1. Simplified logical flow chart.

Subroutine	Description	Subroutine	Description
INIT	Performs adjoint reversals of veloci- ties and fission fractions, mixes cross sections, modifies geometry, calculates	IFLUXN	Subroutine to normalize the fluxes before each group calculation. It is called by INNER, INNER2, and INNERP.
	areas and volumes, and calculates fis- sion neutrons. INIT is called by 3DDT and calls CLEAR and ERRO2.	CNNP	Performs convergence tests, computes a new eigenvalue, and computes new param-
CLEAR	Sets an array of a specified length equal to a given constant. It is called	00050	called by 3DDT and calls ERRO2.
	by INIT, OUTER, S8850, GRAM, and AVERAG.	\$8850	fluxes, total flux, power density, and
FISCAL	Calculates fission sums and performs fission normalization. FISCAL is called by 3DDT.		fission source rate and produces a flux dump on cards or tape. It is called by 3DDT and calls PRT, S8830, S8847,
S8830	Prints the monitor line after each out-		and CLEAR.
	er iteration. It is called by 3DDT and S8850.	S8847	Computes and prints overall neutron balance table. It is called by \$8850.
OUTER	Performs a complete outer iteration. OUTER calls INNER2, INNER, INNERP, and CLEAR and is called by 3DDT	PRT	Subroutine to print any two-dimensional array. It is called by \$8850.
CONSTS	Calculates coefficients for the flux equation. CONSTS is called by 3DDT.	GRAM	Calculates and prints the mass of each material in each zone and the zone volume. It is called by 3DDT and calls
INNER2	Calculates the flux in a specified group		CLEAR.
	when IM.GT.JM and IGE = 0. It is called by OUTER and calls IFLUXN.	INPB	Reads and prints the input burnup data. It is called by 3DDT.
INNER	Calculates the flux in a specified group when IM.LE.JM or IGE = 1. It is called by OUTER and calls IFLUXN.	AVERAG	Calculates one-group zone-averaged fluxes, fission cross sections, and absorption cross sections and the
INNERP	Calculates the flux in a specified group for periodic boundary conditions		breeding ratio. It is called by 3DDT and calls CLEAR.
	in the $\theta$ (or Y) direction. It is called by OUTER and calls IFLUXN.	MARCH	Calculates the time-dependent isotopic concentrations. MARCH is called by 3DDT.

#### APPENDIX D

# VARIABLE-DIMENSIONED ARRAYS AND NONSUBSCRIPTED COMMON VARIABLES

The variable-dimensioned arrays in 3DDT are described in Table D-I. These arrays are stored in the container block A(25000). Nonsubscripted common variables are described in Table D-II. The tables are included for the benefit of those readers who may want to make changes or additions to 3DDT.

#### TABLE D-I

#### DESCRIPTION OF VARIABLE-DIMENSIONED ARRAYS

Variable (Dimension)	Description
AO(JM,KM)	Radial area elements for each rotational and axial mesh interval. For IGE=1, area elements do not include radius factor.
Al(IM,KM)	Rotational area elements for each radial and axial mesh interval.
A2(IM,JM)	Axial area elements for each radial and rotational mesh interval.
ALAM(ML)	Decay constants of the input materials.
ATW(ML)	Atomic weights of the input materials.
ANU(IGM)	Temporary storage for $v$ in S860. Stored in location LK6 in array A.
AXS(ML)	Spectrum-averaged absorption cross section for each input material for the current zone.
BRDT(IZM)	Breeding ratio by zone.
C(ITL,IGM,MT)	Temporary storage of cross sections by position in table, group, and material in subroutine S860. Storage starts at location LN2B in array A.
CO(ITL,MT)	Cross-section array for current group by position in table and material.
CXB(IM,KM)	Back boundary constants by radial and axial position for current group.
CXR (JM, KM)	Right boundary constants by rotational and axial position for current group.
CXS1(IM,JM)	Mesh volume left boundary constants by radial and rotational posi- tion for current group and axial position.
CXS2(IM,JM)	Mesh volume front boundary constants by radial and rotational posi- tion for current group and axial position.
CXS3(IM,JM)	Mesh volume bottom boundary constants by radial and rotational posi- tion for current group and axial position.
CXS3P(IM,JM)	Same as CXS3 but for plane above current axial position.
CXS4(IM,JM)	Mesh volume central constants by radial and rotational position for current group and axial position.
CXT(IM,JM)	Top boundary constants by radial and rotational position for current group.
EO(IGP)	Fission rate for each group, EO(I), and sum over all groups, EO(IGP).
El(IGP)	Fission neutron source for each group, EO(I), and sum over all groups, E1(IGP).
E2(IGP)	Inscatter source for each group, E2(I), and sum over all groups, E2(IGP)
E3(IGP)	Outscatter from each group, E3(I), and sum over all groups, E3(IGP).
E4(IGP)	Absorptions in each group, E4(I), and sum over all groups, E4(IGP).
E5(IGP	Left boundary leakage for each group. E5(I), and sum over all groups, E5(IGP).

#### TABLE D-I (continued)

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Variable (Dimension)	Description
E6(IGP)	Right boundary leakage for each group, E6(I), and sum over all groups, E6(IGP).
E7(IGP)	Front boundary leakage for each group, $E7(I)$ , and sum over all groups, $E7(IGP)$ .
E8(IGP)	Back boundary leakage for each group, E8(I), and sum over all groups, E8(IGP).
E9(IGP)	Bottom boundary leakage for each group, E9(I), and sum over all groups, E9(IGP).
E10(IGP)	Top boundary leakage for each group, E10(I), and sum over all groups, E10(IGP).
Ell(IGP)	Total leakage for each group, Ell(I), and sum over all groups, Ell(IGP).
EF(IGM)	Temporary storage for group distribution of flux in S862 and SINUS. Stored at location LK6 in array A.
FO(IM,JM)	Fission rate (summed over groups) for previous outer iteration by radial and rotational position for current axial position.
F2(IM,JM)	Same as FO but for current outer iteration.
FXS(ML)	Spectrum-averaged fission cross section for each input material and current zone.
HA(IM or JM)	Temporary storage for line inversion.
HOLN(ML)	Names of materials.
IO(M01)	Mix number table.
I1(M01)	Material number table.
I2(M01)	Atom density table.
I3(M01)	Atom density table for GRAM calculation.
IDZNO(KM)	Identifies zone number plane to be used for each axial position.
K(IM)	Temporary storage for radial material numbers in MAPR. Stored at location LCO in array A.
K6(IGM)	Effective fission neutron spectrum.
K7(IGM)	Input fission neutron spectrum.
LAXS(IZM)	Locates array AXS(ML) in ECS for each zone.
LCN(ML,2)	Identifies source isotopes for capture in burnup calculations.
LCXB(IGM)	Locates array CXB(IM,KM) in ECS for each group.
LCXR(IGM)	Locates array CXR(JM,KM) in ECS for each group.
LCXS1(KM)	Locates array CXS1(IM,JM) in ECS for each axial position.
LCXS2(KM)	Locates array CXS2(IM,JM) in ECS for each axial position.
LCXS3(KM)	Locates array CXS3(IM,JM) in ECS for each axial position.
LCXS4(KM)	Locates array CXS4(IM,JM) in ECS for each axial position.
LCXT(IGM)	Locates array CXT(IM,JM) in ECS for each group.
LD(ML)	Identifies source isotopes for decay in burnup calculations.
LECS2(KM)	Locates array S2(IM.JM) in ECS for each axial position.
LFDENO(KM)	Locates array FO(IM,JM) in ECS for each axial position.
LFDEN1(KM)	Locates array F2(IM,JM) in ECS for each axial position.
LFLXO(KM)	Locates array NO(IM,JM) in ECS for each axial position.
LFLX1(KM)	Locates array N2(IM,JM) in ECS for each axial position.
LFN(ML,7)	Identifies source isotopes for fission in burnup calculations.
LFXS(IZM)	Locates array FXS(ML) in ECS for each zone.
LIDZ(MO6)	Locates array MO(IM,JM) in ECS for each unique zone number plane.
LMASS(IZM)	Locates array MASS(ML) in ECS for each zone.

TABLE D-I (continued)

Variable (Dimension)	Description
LMASSP(IZM)	Locates array MASSP(ML) in ECS for each zone.
LVOL(KM)	Locates array VO(IM,JM) in ECS for each axial position.
MO(IM,JM)	Zone numbers by radial and rotational position for current axial position.
M1(IM,JM)	Same as MO(IM,JM) but for plane below current axial position.
M2(IZM)	Material numbers by zone.
MASS(ML)	Current mass of each input material for the current zone.
MASSP(ML)	Material mass for the previous burnup time step for the current zone.
MATN(ML)	Material number for burnable isotopes.
NO(IM,JM)	Flux by radial and rotational position for current group and axial position and previous outer iteration.
N2(IM,JM)	Same as NO(IM,JM) except for current outer iteration.
N2B(IM,JM)	Same as N2(IM,JM) except for plane below current axial position.
N2T(IM,JM)	Same as N2(IM,JM) except for plane above current axial position.
NBR(ML)	Indicates fertile or fissile isotope for breeding ratio calculation.
PA(IM or JM)	Temporary storage for line inversion.
PHIB(IZM)	Average total flux by zone.
RO(IP)	Initial radii (or X mesh).
R1(IP)	Current radii. Also used for temporary storage of radial flux distribution in SINUS.
R2(IM)	Radial zone numbers.
R3(IZ)	Radial zone modifiers.
R4(IM)	Average radii.
R5(IM)	Length of each radial mesh interval.
RF(IM)	Temporary storage for radial flux distribution in S862. Stored at location LRO in array A.
S2(IM,JM)	Total source by radial and rotational position for current group and axial position.
TAO(JP)	Initial thetas (or Y mesh).
TA1(JP)	Current thetas. Also used for temporary storage of rotational flux distribution in SINUS.
TA2(JM)	Rotational zone numbers.
TA3(JZ)	Rotational zone modifiers.
TA4(JM)	Average thetas.
TA5(JM)	Length of each rotational mesh interval.
TF(JM)	Temporary storage of rotational flux distribution in S862. Stored in location LTAO in array A.
VO(IM,JM)	Volume elements by radial and rotational position for current axial position.
V7(IGM)	Neutron velocities by group.
VOL(IZM)	Volume by zone.
ZO(KP)	Initial axial (Z) mesh.
Z1(KP)	Current axial mesh. Also used for temporary storage of axial flux distribution in SINUS.
Z2 (KM)	Axial zone numbers.
Z3(KZ)	Axial zone modifiers.
Z4(KM)	Average axii.

#### TABLE D-I (continued)

Variable<br/>(Dimension)DescriptionZ5(KM)Length of each axial mesh interval.ZF(KM)Temporary storage for axial flux distribution in S862. Stored at<br/>location LZ0 in array A.

#### TABLE D-II

TABLE D-II (continued)

DESCRIPTIO	ON OF NONSUBSCRIPTED COMMON VARIABLES	Variable	
Variable	Description	G07	Inner iteration maximum per group.
A01	Cross-section format indicator.	IDMTPS	Flag for producing flux dump on magnetic tape.
A02	Theory (regular or adjoint) indicator.	IGE	Geometry indicator.
ALA	Lambda for current outer iteration.	IGEP	= IGE + 1.
B01	Left boundary condition.	IGM	Number of energy groups.
во ]	Right boundary condition.	IGP	= IGM + 1.
B03	Front boundary condition.	IGV	Group indicator in OUTER, INNER, INNER2,
B04	Back boundary condition.		and INNERP.
B05	Bottom boundary condition.	IHS	Position of sigma self-scatter in cross-
B06	Top boundary condition.	<b>T</b> 112	Section table.
B07	Flag used for internal computation in FISCAL and INIT.	THE	section table.
CNT	Parabolic interpolation trigger in CNNP.	IHTM1	= IHT - 1.
CVT	Lambda convergence trigger.	II	Inner iteration count for a single group.
DAY	Accumulated burnup time in days.	IM	Number of radial (or X) mesh intervals.
DELT	Burnup interval in days.	MLMI	= IMxJM.
D05	Outer iteration maximum.	IMKM	= IMxKM.
E01	Used for temporary storage in INIT,	104	Eigenvalue type indicator.
	OUTER, INNER, INNER2, INNERP, and	IP	≖ IM + 1.
500	FISCAL. EUT = $I - \lambda$ in CNNP.	IPCUR	Flag for punching atom densities after
EUZ	E02 = $ 1 - \lambda $ in CNNP.	TPELV	Elag for purching flux dump
E03	Absolute value of the difference be-	TTEMP	Temporary storage
	tween lambda for the current outer		Temporary storage
	iteration and that for the previous iteration.	ITEMPO	
EPS	Lambda (total fission source) conver-	I TEMP I	
	gence criterion.	TTEMP 2	Time light (and) for problem
EPSA	Eigenvalue search convergence criterion.	IIIMOF	Since fimit (Sec) for problem.
EQ	Eigenvalue slope in search options.	TTL	Cross-section table length.
ESFL	Extended core storage field length.	12	Number of radial zone modifiers.
EV	Eigenvalue for current outer iteration.	IZM	Number of material zones.
EVM	Eigenvalue modifier.	129	= 12M + 1.
EVP	Previous eigenvalue in search options.	ML	Number of rotational (1) mesh intervals.
EVPP	Two eigenvalues back in search options.	JMKM	= JMXKM.
EVPT	Eigenvalue for previous outer iteration.	JP	
FEF	Useful energy released per fission	JZ	Number of rotational zone modifiers.
201	(205.0 MeV).	KU /	remporary storage for GU/.
GÜB	Inner iteration pointwise flux conver- gence criterion.	км	Number of Axial (2) mesh intervals.

TABLE D-II (continued)

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TABLE D-II (continued)

Variable	Description	Variable	De	scription	
кр	= KM + 1.	ORF	Overrelaxation	factor.	
KPAGE	Page counter for monitor print.	ORFP	ORF for $ 1 - \lambda $	$\leq$ 10EPS.	
KZ	Number of axial zone modifiers.	P02	Outer iteration	count.	
LAH	Upper limit on $ig 1-\lambdaig $ for parabolic	POD	Parameter oscil	lation dampe	r.
	interpolation.	S01	Neutron source	rate or powe	er level.
LAL	Lower limit on $ 1 - \lambda $ for parabolic interpolation.	S02	Parametric eige	nvalue type	indicator.
LAP	Converged lambda for previous eigen-	S03	Parametric eige	nvalue.	
LADD	value.	S04	Inner Iteration $ 1 - \lambda  > 10$ EPS	maximum per	group for
LAFT	back.	SK7	Sum over all gr	oups of K7(I	GM).
LAR	Lambda for previous outer iteration.	Т7	= alpha/velocit	у.	
LC	Total number of inner iterations for a single outer iteration.	T11	Total fission m vious outer ite	eutron sourcerations.	e for pre-
LCP	Total number of R- $ heta$ (or X-Y) iterations	TEMP	Temporary stora	lge.	
	for a single outer iteration.	TEMP1	Temporary stora	lge.	
LFLG	Error flag used in ECS read and write	TEMP 2	Temporary stora	ıge.	
M01	Mixture encoifications table length	TEMP 3	Temporary stora	ige.	
MOE	Transition reduction factor for	TEMP4	Temporary stora	ige.	
MOS	$R-\theta$ iterations.	TI	Current time (s	ystem clock)	•
M06	Number of $R-\theta$ (or X-Y) planes with	т06	$= 1/0 (\delta \text{ calcul})$	lation/no).	
NO7	unique zone numbers.	TSD	$= 1.602 \times 10^{-13}$	FEF (MW-see	c/fission).
MO7	Flux guess indicator.	V11	Total source fo	or the curren	nt group.
MCR	from cards.	The r	emaining nonsubs	cripted comm	on variables
ML	= MCR + MTP.	are pointe	rs for locating	in the cont:	aner block
MT	Total number of materials including mixtures.	A, the var	riable-dimensioned	i arrays spec	cified.
MTP	Number of material cross sections from		Pointer		Pointer
	tape (=0 currently).	Variable	for Array	Variable	IOT ATTAY
NCFC	Trigger for computation of flux	LAO LA2	AO(JM,KM) A2(IM,JM)	LA1 LALAM	ALAM(ML)
	constants.	LATW	ATW (ML)	LAXSC	AXS(ML)
NCON	Burnup control.	LBRDT	BRDT (12M)	LCO	CO(ITL,MT)
NCR1	Disk file (tape simulation) for storing	LCXBC	CXB(IM,KM)	LCXCSP	CXSI(IM,JM)
	cross sections by group, material, and	LCXSC2	CXS2(IM.JM)	LCXSC3	CXS3(IM,JM)
	position in table.	LCXSC4	CXS4(IM,JM)	LCXTC	CXT(IM,JM)
NCXS1	Disk file (tape simulation) for boun-	LEO	EO(IGP)	LE1	E1(IGP)
	dary and central volume constants for	LE2	E2(IGP)	LE3	E3(IGP)
	flux calculation. Constants are stored	LE4	E4(IGP)	LE5	E5(IGP)
	by group and mesh point.	LE6	E6(IGP)	LE7	E7(IGP)
NFLUX1	Disk file (tape simulation) for storing	LE8	E8(IGP)	LE9	E9(IGP)
	current fluxes. Fluxes are stored by	LEIO	E10(1GP)	LEII	EII(IGP)
	group and mesh point.	LFU	FU(IM,JM)		FZ(IM, JM)
NCOTO		LFASC	FAS(ML)	LINZNO	IDZNO(KM)
NGOTU	loop.	LIOLN	IO(MO1)	LIL	11(M01)
NINP	System input file.	LI2	12(MO1)	LI3	I3(M01)
NIM	bystem input file.	LK6	K6(IGM)	LK7	K7(IGM)
NUUT	System output file.	LLAXS	LAXS(LCM)	LTCA63	LCXS2(M)
NPRT	Print control.	LICKSS	LCXS3(KM)	LLCXS4	LCXS4 (KM)
NSCRAT	Dick file (tame simulation) used as a	LLCXB	LCXB(IGM)	LLCXR	LCXR (1GM)
novini	scratch file.	LLCXT	LCXT(IGM)	LLD	LD(ML)
		LLECS2	LECS2(KM)	LLFDNO	LFDENO(KM)
NXCM	= ITL - IHS.	LLFDN1	LFDEN1(KM)	LLFLX0	LFLXO(KM)

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#### TABLE D-II (continued)

TABLE D-II (continued)

	Pointer		Pointer		Pointer		Pointer
Variable	for Array	Variable	for Array	Variable	for Array	Variable	for Array
LLFLX1	LFLX1(KM)	LLFN	LFN(ML,7)	LR2	R2(IM)	LR3	R3(IZ)
LLFXS	LFXS(IZM)	LLVOL	LVOL (KM)	LR4	R4(IM)	LR5	R5(IM)
LLIDZ	LIDZ(MO6)	LLMASS	LMASS (IZM)	LS2	S2(IM,JM)	LTA0	TAO(JP)
LLMASP	LMASSP(IZM)	LMO	MO(IM,JM)	LTA1	TA1(JP)	LTA2	TA2(JM)
LM1	M1(IM,JM)	LM2	M2(IZM)	LTA3	TA3(JZ)	LTA4	TA4(JM)
LMASSC	MASS (ML)	LMASPC	MASSP (ML)	LTA5	TA5(JM)	LVO	VO(IM,JM)
LMATN	MATN (ML)	LN2	N2(IM,JM)	LV7	V7(IGM)	LVOLC	VOL(IZM)
LN2B	N2B(IM.JM)	LN2T	N2T(IM,JM)	LZ0	ZO(KP)	LZl	Z1(KP)
LNBR	NBR(ML)	LNO	NO(IM.JM)	LZ2	Z2(KM)	LZ 3	Z3(KZ)
LPA	PA(IM or JM)	LPHIB	PHIB(IZM)	LZ4	Z4 (KM)	LZ5	25 (KM)
LRO	RO(IP)	LR1	R1(IP)		•		

#### APPENDIX E

#### EXECUTED SAMPLE PROBLEM

Input data and selected computer output for a simple three-group, three-zone, one-step burnup problem in X-Y-Z geometry are given in the following pages. The reactor consists of a cubical core region 80 cm on a side, surrounded on all sides by a blanket region 30 cm thick. Thus, the reactor is a cube 140 cm on a side. The core region is divided into two material zones: an inner cube 40 cm on a side and an outer cubical shell 20 cm thick. Initial compositions of the core and blanket zones are given in Table E-I.

Because of symmetry, only one octant of the reactor is represented in the calculational model. A reflective boundary condition is applied to the interior boundaries of the octant and a vacuum boundary condition is applied to the exterior boundaries. The spatial mesh contains 14 intervals in each dimension and the initial flux profiles are computed internally from cosine distributions and an input energy distribution.

#### TABLE E-I

INITIAL COMPOSITIONS OF CORE AND BLANKET ZONES

Atom De	nsity in Units of	10 <sup>24</sup>
Core	Core	Blanket
(Zone 1)	(Zone 2)	(Zone 3)
0.008	0.010	0.030
0.0016	0.002	=0
0.0001	0.00012	
0.020	0.020	•••
0.006	0.006	
0.013	0.013	0.0062
=0	=0	=0
	Atom De Core (Zone 1) 0.008 0.0016 0.0001 0.020 0.006 0.013 =0	Atom Density in Units of Core         Core           (Zone 1)         (Zone 2)           0.008         0.010           0.0016         0.002           0.0001         0.00012           0.020         0.020           0.006         0.006           0.013         0.013

In the sample problem, the initial  $k_{eff}$  calculation is followed by a burnup interval of 30 days with the reactor at 1000 MWT total power. Following the burnup interval, a final  $k_{eff}$  is computed for the new material compositions resulting from fuel depletion, breeding, and fission-product buildup.

# INPUT DATA

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JODT SAMPLE	PROBLEM . 3-G	ROUP . 3-ZONE .	x-Y-Z,14X]	4X14 MESH	30 DAY BUR	NHP	CARD 1
2	0	0		3	ł	5	CARD 2
14	14	14	0.0 E.	0.00	E+00 1.0	E-04	CARD 3
1	0	1		Û	1	0	CARD 4
13	21	10		0	Ο	0	CARD 5
0	0.0 E+00	3		6	7	9	CARD 6
-1.25 E+02	2	3		5	30	10	CARD 7
0.0 E+00	1.0 E-03	5.0 E-03	5.0 E-	01 1.0	E+00 1.0	E-03	CARD H
1	0	1	2	270 1.4	E+00	75	CARI) 9
U238C238.05	0.0 U-238	3-GROUP CROS	SS SECTION	IS FOR COR	E.		
2.76600166-11	.83582941-14	4.7403077106	2.76596583	-14.76023	543-15.0169	050552	1
4.11970414220	•	0.	1.59990100	-10.	8,0399	900000	2
1.59999000-10	•	4.1999900000	8.0000000	005.42741	058-10.		٤
6.82093260-10	•	1.20796811+30	6.81993172	-10.	1.2761	5743+1	4
1.20795810+13	.999999000-2	1.18596818-2	C CLOTIN		<b>-</b>		5
PU239C239.05	0.0 PU-239	3-6R00P CR0	SS SECTION	IS FUR CURI	259865 1030	000000	,
1.88391315891	.8035487002	3.2199908291	1.00390717	535.40035	120042.1030	000000	
2.86092519050	•	3. 20000000000	L • 8.39999000	001.07014	762-10	900000	<i>.</i>
1.84000000004	-8099900000	1 06101053+1	0000000 5.5970550 5	6003.00130	742-10. 782051 3663	4736+1	5
3.03310603156	.3040314090	- 00101955+1. - 00220120-2	3.03351033		105001.0043	+730-1	5
1.00102900+10 DU2400240 0E	0 0 PU-240	3-690110 CD04	S SECTION	S FUR COR	F		
1.26693901151	.1508329669	3.68961589611	1.26693349	403.88098	953884.9565	493901	1
3.24001106910		).	2.80000000	-14.99997	000-27.9199	900000	2
2.80000000-11	.58999000-18	1.19999000001	7.86999000	003.90563	684-10.		3
6.83166019-12	.27792271-9	1.09999115+10	6.83158660	-16.83375	196-91.1683	0702+1	4
1.09999042+15	.0000000-25	5.90356248-2		•			5
CARBON 12.01	0.0 CARBON	3-GROUP GROS	SS SECTION	S FUR CORI	Ε		
2.01936907-60	• 2	2.27920041990	)•	0.	2.2792	104199	1
1.99159767030	• (	). 2	2.8421709-	140.	3.5899	300000	2
0. 0	•	3.5899900000	3.15699000	102.87600	730-10.		3
5.29670305-60	• "	+.30379842053	3.0626253-	160.	4.3037	194205	4
4.30379312384	.33000000-10	).			_		5
NA 22.99	0.0 NA	3-GROUP CROS	SS SECTION	S FOR CORI			,
3.16633341-30	•	2.14081418003	3.16868794	-3().	2.74404	+34593	1
2.43232536680	• (	).	9.00990000	-40.	3.1002 770 10	900000	2
9.00030000-40	•	3.10/1900002	2.91129000	002.87600	130-10.	200260	3
1.02053027-30	•	++4/04566263	1.02340251	-30.	4.4/14	500248	4
4.47045949451	•94999010-17	2.09510200-7	C CROTION	C RUD COUR	-		5
FE C 55.85	0.0 FE	3-6K00P CK02	5 5ECTION 5 51737572	5 FUR CURE	2 2020	465662	1
8.52005273-30	•	:•17330727190 \	0.00101000	-20.	2.4300	100000	2
C.UD.194948240	• (	. * * * * * * * * * * * * * * * * * * *	2.35800000	001.258429	594-10.	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	3
1.01454344 20	• 6		-01677373	-20.	3,4604	381377	4
1.01034300-20		2.77443653-3		200	51,504.		5
3.43021210101	• 0 7 7 7 7 0 0 0 - 2 .	,,					

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U23AB238.05 0.0 U-238 3-GROUP CROSS SECTIONS FOR BLANKET 2.10214983-19.65279995-25.00627110142.10212570-12.49176842-15.2164845447 1 4.39923251220. 0. 1.59990100-10. 2 8.0399900000 1.59999000-10. 8.19999000008.0000000005.28680781-10. 3 5.25432310-10. 1.18704180+15.25332211-10. 1.23957502+1 4 1.18703179+13.99999000-27.83562669-2 5 PU239B239.05 0.0 PU-239 3-GROUP CROSS SECTIONS FOR BLANKET 1.85392199141.76157876913.48437062771.85391822855.26578360205.3382888562 1 3.13310604190. 0. 1.83999000001.67014000006.5599900000 2 1.8400000004.80999000008.39999000006.5100000003.00112930-10. 3 2.59306957562.06590039671.07317383+12.59316746095.90846685841.33249058+1 4 1.07318362+15.00000000-25.11478920-2 5 FE 3-GROUP CROSS SECTIONS FOR BLANKET FE B 55.85 0.0 2.21460906559.23030490-30. 9.23190161-30. 2.2238393696 ) 2.4300000000 2 2.07976530500. 0. 1.00101000-20. 1.0000000-20. 2.4400000002.35899000001.31642478-10. 3 2.88873591501.00254823-20. 4 1.00239771-20. 2.8987613974 2.88873742027.09999000-23.19968474-3 5 0.0 FISSION PRODUCTS 3-GROUP CROSS SECTIONS FIS PR 1.0 +00 0.0 +00 6.7 -02 0.0 +00 6.7 -02 6.7 -02 0.0 0.0 +00 0.0 +00 0.0 +00 1.5 -01 0.0 +00 0.0 +00 +00 1.5 -01 0.0 1.5 -01 0.0 +00 0.0 +00 0.0 +00 +00 0.0 +00 3.3 -01 0.0 +00 3.3 -013.3 -01 0.0 +00 0.0 +00 0.0 +00 0.0 3 R0 203 0.0 203 20.0 205 40.0 70.0 203 0.0 203 20.0 205 40.0 70.0 3 TAO 205 40.0 70.0 3 ZO 203 0.0 203 20.0 3 EF 1.0 1.0 1.0 3104 1104 2106 3 M0 PL 1 2106 104 1104 104 1104 2106 3104 1104 2106 3 MO PL 1 2106 3108 2106 3108 2106 3 M0 PL 1 108 33 10 PL 1 108 2106 34 3108 2106 3108 2106 3 MU PL 2 108 2106 3108 2106 3108 2106 3 MO PL ? 108 2106 108 2106 3108 2106 34 33 MU PL 2 MO PL 3 33 4 1D7N0 1104 2106 33 104 12 133 M2 11 0.090 0.014 0.896 3 Κ7 6.7000+08 ν7 2.0432+083 1.2568+09 12105 133 10 108 11108 0 1 2 3 4 5 11' 1 2 3 11 10 0 6 5 4 6 10 U 7 11 9 103 8 11 -03 12 8.0 -03 1.6 -03 1.0 -04 2.0 -02 6.0 0.0 +00 1.3 -02 1.0 -20 0.0 +00 1.0 -02 2.0' -03 1.2 -04 12 1.3 -02 1.0 -20 +00 3.0 -02 12 2.0 -02 6.0 -03 0.0 6.2 -03 1.0 -203 12 1.0 -20

6	0 3	0.0									BURNUP	1
1	1	0	0	0	0	0	0	0	0	0	(IRURNUP	2
2	2	0	1	0	0	0	0	0	0	0	0 RURNUP	2
3	ĩ	0	2	0	0	0	0	0	0	0	0.BURNUP	2
7	ī	0	0	0	0	0	0	0	O	0	0 BURNUP	2
8	2	0	4	0	0	0	0	0	U	0	0BURMUP	2
10	0	0	0	0	1	2	3	4	5	0	OBURNUP	2
0	_										BURNUP	1

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#### BUDT SAMPLE PROBLEM, 3-GROUP, 3-ZONE, X-Y-Z, 14X 14X14 MESH, 30 DAY BURNUP

Λ01 Λ02 16E 12M 104 ΜC7	CROSS SECTION FORMAT INVICATOR (1/2 = ENWL/LASL) THEORY (0/1=REGULAR/ACJCINT) GEONETRY (0/1=X-Y-Z/R-THETA-Z) NUMBER OF MATERIAL ZONES EIGLNVALUE TYPE (1/2/3/4=K/ALPHA/C/DELTA) FLUX GUESS (0/1+2+3/4/5 = NONE/CARCS/TAPE/SINUSOIC)	2 0 3 1 5
IM JM EV EVM EPS	NUMBER OF RADIAL (X) INTERVALS NUMBER OF ROTATIONAL (Y) INTERVALS NUMBER OF AXIAL (Z) INTERVALS FIRST EIGENVALUL GUESS EIGENVALUE MODIFIER LAMBDA CONVERGENCE CRITERION	14 14 0. 0. 1.0CCOE-04
B01 B02 B03 BC4 B05 B06	LEFT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE) RIGHT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE) FRONT BOUNDARY CONDITION (0/1/2 = VAC/REFL/PERIODIC) BACK BOUNDARY CONDITION (0/1/2 = VAC/REFL/PERIODIC) BOTTOM BOUNDARY CONDITION (0/1 = VACUUM/REFLECTIVE) TOP BOUNDARY CONDITION (0/1 = VACUUM/REFLECTIVE)	1 0 1 0 1 0
MT NO1 MCR IZ JZ KZ	TUTAL NUMBER OF MATERIALS INCLUDING MIXES Number of mixture specifications Number of materials from cards Radial Zones (delta option only) Rufational Zones (delta option only) Axial Zones (delta option only)	13 21 10 0 0 0
502 503 1GM 1HT 1H5 1TL	PARAMETRIC EIGENVALUE TYPE (0/1/2=NUNE/K/ALPHA) PARAMETRIC EIGENVALUE NUMBER OF GROUPS PUSITION OF SIGMA TRANSPORT PUSITION OF SIGMA SELF SCATTER CRUSS SECTION TABLE LENGTH	0. 3 6 7 9

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SOL NEUTRUM SOURCE RATE (POSITIVE) ON POWER IN MWT(NEGATIVE)	-1.2500E+02
MOS INNER ITERATION RECUCTION FACTOR FOR R-THEIA (X-Y) PLANE	2
NC6 NUMBER OF R-THETA(X-Y) PLANES HAVING UNIQUE ZONE NOS.	3
SO4 INNER ITER MAX PER GROUP FOR 1-LAMEDA .GT. 10+EPS	5
DOS OUTER ITERATION MAXIMUM	30
GO7 INNER ITERATION MAXIMUM PER GROUP	10
GC5 NOT USED	0.
GC6 INNER ITERATION FLUX CONVERGENCE CRITERION	1.CCC0E-03
LAL LAMBDA LUWER LIMIT	5.CC00E-03
LAH LAMBDA UPPER LIMIT	5.0C00E-01
POD PARAMETER OSCILLATION DAMPER	1.CG00E+C0
EPSA EIGENVALUE SEARCH CONVERGENCE CRITERION	1.CC00E-03
IPFLX PUNCH FLUX DUMP (0/1=YES/NC)	1
IPCUR PUNCH DEWSITIES (0/1=NO/YES)	0
IDMTPS PREPARE SPECIAL FLLX DUMP TAPE (0/1=YES/NO)	1
ITIMOF TIME LIMIT IN SECONDS	270
ORF OVER RELAXATION FACTOR	1.4000E+00
ESFL ECS FIELD LENGTH IN THOUSANDS (CCTAL)	75

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#### THE FOLLOWING NUCLIDES ARE FRCM CARDS

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1	11238C	U-238 3-GENUP CRUSS SECTIONS FOR CORE	
2	PU239C	PU-239 3-GROUP CRESS SECTIONS FOR CORE	
3	PU240C	PI)-240 3-GROUP CHESS SECTIONS FOR CORE	
4	CARBON	CARBON 3-GROUP CRESS SECTIONS FOR CORE	
5	ΝA	NA 3-GROUP CRCSS SECTIONS FOR CORE	
Â	FF C	FE 3-GROUP CRESS SECTIONS FOR CORE	
7	112388	U-238 3-GROUP CRCSS SECTIONS FOR BLANKE	Т
1 0	D1123 QH	PU-239 3-GROUP CRESS SECTIONS FOR BLANKE	Т
0		FE 3-GROUP CRESS SECTIONS FOR BLANKE	Т
3		EISTIN PRODUCTS 3-GROUP CROSS SECTIONS	
10	LT2 BK	FISION INCOURS 5 DAMA BACKS CLEATING	

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MESH BOUNDARIES (RO/TAO/ZO=RADIAL/RCTATICNAL/AXIAL)

RO 15 0. 5.0COOE+01	5.0000E+00 5.50CCE+01	1.0000E+01 6.0000E+01	1.5000L+01 6.5000E+01	2.0000E+01 7.0000E+01	2.5000E+01	3.00C0E+C1	3.5000E+01	4.0000E+01	4.5000E+01
TA 0 15 0. 5.0000E+01	5.C000E+00 5.5C00E+01	1.0000E+01 6.0000E+01	1.5000E+01 6.5000E+01	2.0000E+01 7.0000E+01	2.5000E+01	3.000CE+01	3.5000E+01	4.0000E+01	4.5000E+01
ZC 15 0• 5•000CE+01	5.0000±+00 5.5000E+01	1.0000E+01 6.0000E+01	1.5000E+01 6.5000E+01	2.0000E+01 7.0000E+01	2.5000E+01	3.0000E+01	3.5000E+01	4.0000E+01	4.5000E+01

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 FLUX GUESS
 RF/TF/ZF=RADIAL/RCTATICNAL/AXIAL FROM SINUS, EF=ENERGY FROM CARDS

 RF
 14

 •99863E+00
 •98769E+00
 •96593E+00
 •93358E+00
 •89101E+00
 •83867E+00
 •70711E+00
 •62932E+00
 •54464E+00

 •45399E+00
 •35837E+00
 •25882E+00
 •15644E+00
 •89101E+00
 •83867E+00
 •70711E+00
 •62932E+00
 •54464E+00

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 •77715E+0C
 •70711E+00
 •62932E+00
 •54464

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#### ZUNE NUMBERS FOR UNIQUE R-THETA (X-Y) PLANES

PLANE = 1

MO	196									
	1	1	1	1	2	2	2	2	3	3
	3	3	3	3	1	1	1	1	2	2
	2	2	3	3	3	3	3	3	1	1
	1	1	2	2	2	2	3	3	3	3
	3	3	1	1	1	1	2	2	2	2
	3	3	3	3	3	3	2	2	2	2
	2	2	2	2	3	3	3	3	3	3
	2	2	2	2	2	2	2	2	3	3
	3	3	3	3	2	2	2	2	2	2
	2	2	3	3	3	3	3	3	2	2
	2	2	2	2	2	2	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3
	3	3	3	3	3	3	3	3	3	3

PLANE =

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IDZNO	14									
	1	1	1	1	2	2	2	2	3	3
	3	3	3	3						
MATERIAL	NUMBERS	8 Y ZONE								
M 2	3									
	11	12	13							

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#### FISSION SPECTPUM

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К7 3 8.9600E-01 9.0000E-02 1.4000E-02

#### NEUTRON VELOCITY

V7 3 1.2568E+C9 6.7C00E+08 2.0432E+08

# MIXTURE SPECIFICATIONS (10/11/12=MIX NUMBER/MAT. NUMBER FOR MIX/MATERIAL DENSITY)

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IC	21 11 12 13	11 12	11 12	11 12	11 12	11 12	11 13	11 13	12 13	12 13
I 1	21 0 2 10	1 3	2 4	3 5	4 6	5 1C	6 C	10 7	0 8	1 9
12 0. 2.0000 1.0000	21 E-03 E-20	8.0000E-03 1.2000E-04	1.6000E-03 2.0000E-02	1.0000E-04 6.0000E-03	2.0000C-02 1.3000E-02	6.0000E-03 1.00C0E-20	1.300CE-02 0.	1.0000E-20 3.000CE-02	0. 1.0000E-20	1.0000E-02 6.2000E-03

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ZONE	M	Δ 1	>					Ŕ	AĹ	)1	٨ι	•																			
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THE ZONE MAP ABOVE AND THE MATERIAL MAP BELOW APPLY TO THE FOLLOWING AXIAL POSITIONS K K = 5
$ \begin{array}{llllllllllllllllllllllllllllllllllll$
MATERIAL NAP 131313131313131313131313131313 131313131313131313131313131313
13131313131313131313131313131313 131313131313131313131313131313 1313131313131313131313131313 13131313131313131313131313
131313131313131313131313131313131313 12121212
12121212121212121213131313131 1212121212
12121212121212121313131313 1212121212121
ZI)NE MAP 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3
3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3
3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3
3 3
3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3
APPLY TO THE FOLLOWING AXIAL POSITIONS K K = 9 K = 10
K = 11 K = 12 K = 13
K = 14

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MATERIAL MAP
13131313131313131313131313131313
12131313131313131313131313131313
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# T I M E = 0.000 D A Y S

м	IXTURE NUMBER	MIX CEMMAND	MATERIAL ATOMIC DENSITY
1	11	С	0•
2	11	1	.8000000E−02
3	11	2	•1600000E-02
4	11	3	.1000000E-03
5	11	4	.20000000E−01
6	11	5	.6000000E-02
7	11	6	-13000000E-01
ห	11	10	.1000000E-19
9	12	С	0.
10	12	1	• 10000000 E-01
11	12	2	.2000000E-02
12	12	3	• 12000000E-03
13	12	4	.2000000E-01
14	12	5	.6000U000E-02
15	12	6	13000000E-01
16	12	10	-10000000E-19
17	13	С	0.
18	13	7	.3000000E-01
19	13	દ	.10C0000E-19
20	13	ç	.62000000E-02
21	13	10	.1000000E-19
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CRUSS-S	ECTION EDIT						
GROUP	1 CROSS-SECTIO	INS			( ) ) 0 705 ( 00	0	0
MAT 1	1.83583E-01	2.76597E-01	4.76024E-01	5.01691E+00	4.11970E+00		0.
MAT 2	1.80355E+00	1.88391E+00	5.48D 36 E+00	5.10390E+00	2.86093E+00	0.	0.
MAT 3	1.15083E+00	1.26693E+00	3.880992+00	4.95655E+00	3.24001E+00	C.	0.
MAT 4	0.	0.	0.	2.27920E+00	1.99160E+00	C.	0.
MAT 5	0.	3.16869E-03	0.	2.74404E+00	2.43233E+00	C •	0.
MAT 6	0.	8.51728E-03	0.	2.20209E+00	2.06395E+00	C•	0.
MAT 7	9.65280E-02	2.1C213E-01	2.49177E-01	5.21648E+00	4.39923E+00	С.	0.
MAT 6	1.76158E+00	1.85392E+00	5.26578E+00	5.33829E+00	3.13311E+00	С.	0.
MAT 9	0.	9.23030E-03	0.	2.22384E+00	2.07977E+00	С.	0.
MAT 1C	0.	6.7CC00E-02	0.	6.700C0E-02	0.	С.	0.
MAT 11	4.469426-03	5.48345E-03	1.29649E-02	1.39473E-01	1.19116E-01	С.	0.
MAT 12	5.58103E-03	6.81555E-03	1.61867E-02	1.51647E-01	1.285656-01	С.	0.
MAT 13	2.89584E-03	6.36360E-03	7.47531E-03	1.70282E-01	1.44872E-01	0.	0.
GROUP	2 CRUSS-SECTIO	DNS					
MAT 1	0.	1.599996-01	0.	8.19999E+00	8.0C000E+00	5.42741E-01	0.
MAT 2	1.67014E+00	1.84CCOE+00	4.80999E+00	8.39999E+00	6.51000E+00	3.08137E-01	0.
MAT 3	4.999976-02	2.8CCC0E-01	1.58999E-01	8.19999E+00	7.86999E+00	3.90564E-01	0.
MAT 4	0.	0.	0.	3.58999E+00	3.15699E+00	2.87601E-01	0.
MAT 5	0.	9.000306-04	0.	3.10719E+00	2.91129E+00	2.876C1E-01	0.
MAT 6	0.	1.0000000-02	0.	2.440006+00	2.35899E+00	1.25843E-01	0.
MAT 7	0.	1.59999E-01	0.	8.199995+00	8.0000CE+00	5.28681E-01	0.
MAT 8	1.670141+00	1.84C00E+00	4.80999E+00	8.39999E+00	6.5100CE+00	3.00113E-01	0.
MAT S	0.	1.0CCC0E-02	0.	2.44000E+00	2.35899E+00	1.31642E-01	0.
MAT 1C	0.	1.5C000E-01	0.	1.500C0E-01	0.	С.	0.
MAT 11	2.67722E-03	4.38739E-03	7.71188E-03	2.02023E-01	1.86477E-01	1.398762-02	0.
MAT 12	3.346286-03	5.448996-03	9.63906E-03	2.21947E-01	2.05239E-01	1.52041E-02	0.
MAT 13	1.67014E-20	4.86197E-03	4.80999E-20	2.61128E-01	2.54626E-01	1.66766E-02	0.
GAUP5	3 CROSS-SECTIO						
MAT 1	0.	6.81993E-01	0.	1.276176+01	1.20796E+01	3.99999E-02	7.78597E-02
MAT 2	2-30463E+00	3.03328E+00	6.59124E+00	1.36435E+01	1.06103E+01	5.00CC0E-02	5.092296-02
NAT 3	2.277925-09	6.83159E-01	6.83375E-09	1.16831E+01	1.09999E+01	5.000000-02	5.90356E-02
MAT A	0	3.062635-16	0.	4.30380E+00	4.30379E+00	4.33COCE-01	0.
	0.	1 023405-03	0.	4.47148F+00	4.47046E+00	1.94999E-01	2.09510E-02
MAT 2	0	1.016775-02	0.	3.46044F+00	3.45027E+00	7.09999E-02	3.77444E-03
MAT C	0	5 253325-01	0.	1.23958E+01 '	1.18703E+01	3.999995-02	7.83563E-02
MAT 0		2 563176+00	5.008475+00	1.33249E+01	1.07318E+01	5.00C00E-02	5.11479E-02
MAI C	2.003905700	1 002555-02	0.	2.89876E+00	2.88874E+00	7.099991-02	3.19968E-03
	0	3 300005-01	0.	3.30000E-01	0.	C.	0.
FAI IC	V • 2 607618-03	1 061585-02	1.054605-02	2.829825-01	2.72465E-01	1.11580E-02	8.85032E-04
	5.001416-03	1 310685-02	1.318255-02	3-14196E-01	3.01089E-01	1.12590E-02	1.06230E-03
MAI 12	4.009206-00	1 582216-02	5.90847E-20	3.898455-01	3.74020E-01	1.64020E-03	2.37053E-03
MAI 13	2.005905-20	1.002216-02	J. 700 41 E-20	J.0707JL 01	51110202 01	110.0102 09	

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TIME (MINUTES)	OUTER I TERATIONS	IN. IT. 30 LCAP	IN. IT. 20 LOOP	E IGENVALUE SLOPE	EIGENVALUE	LAMBDA
.10	0	n	0	0.	0.	0.
. 31	1	15	420	0.	.83482885F+00	.83482885E+C0
• 4 7	2	15	4 20	0.	•97368630E+CO	.11663304E+C1
.64	3	15	4 20	0.	.10188209E+01	10463543E+01
.80	4	15	420	0.	.1034C800F+01	.1C149773E+01
.96	5	15	420	0.	.10404645E+C1	10061740E+01
1.12	6	15	4 20	0.	.10433985E+01	10028200E+01
1.29	7	15	4 20	0.	.10448276E+01	1C013696E+01
1.45	8	15	419	0.	10455372E+01	10006791E+01
1.50	9	3	91	0.	10455627E+01	10000244E+01
1.55	10	3	85	0.	•10456522E+01	•1C000856E+01
FINAL NEUTRON H	BALANCE TABLE					

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GROUP	FISS SOURCE	IN-SCATTER	CUT-SCATTER	ABSORPTION	LEFT LEAK	RIGHT LEAK	FRONT LEAK	BACK LEAK
1	9.32941E+18	3.2768CE+05	6.59828E+18	2.55785E+18	0.	5.79768E+16	0.	5.80457E+16
2	9.37106E+17	6.060026+18	3.64069E+18	2.73898E+18	0.	2.057CCE+17	0.	2.05924E+17
3	1.45772E+17	4.17945E+18	4.44266E+14	4.23067E+18	0.	3.13536E+16	0.	3.13822E+16.
4	1.C4123E+19	1.023956+19	1.02394E+19	9.527496+18	0.	2.9503CE+17	0.	2.95352E+17
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GRUUP	BUTTOM LEAK	<b>JUP LEAK</b>	TOTAL LEAK
1	0.	5.80387E+16	1.740616+17
2	0.	2.05909E+17	6.175346+17
3	с.	3.13834E+16	9.41192£+16
4	0.	2.95332E+17	8.857146+17

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#### FLUX FCR GROUP 1

IZCNTAL PLANE	K= 1							
1	2	3	4	5	6	7	8	9
8.46611E+15	8.33843E+15	8.09129E+15	7.74263E+15	7.30375E+15	6.52439E+15	5.47077E+15	4.16639E+15	2.51690E+15
8.33904E+15	8.21122E+15	7.96363E+15	7.61418E+15	7.17509E+15	6.4C379E+15	5.36632E+15	4.08515E+15	2.46694E+15
8.04307E+15	7+96479E+15	7.71553E+15	7.36233E+15	6.91880E+15	6.16222E+15	5.15696E+15	3.92245E+15	2.36695E+15
7.74585E+15	7.61678E+15	7.36377E+15	6.99954E+15	6.53436E+15	5.79737E+15	4.84124E+15	3.67/74E+15	2.21676E+15
7.3C844E+15	7.17.17E+15	6.92173E+15	6.53584E+15	6.00804E+15	5.3C185E+15	4.41600E+15	3.35003E+15	2.01628E+15
6.53015F+15	6.408996+15	6.16632E+15	5.80007E+15	5.303116+15	4.66355E+15	3.87694E+15	2.93820E+15	1.76583E+15
5.476796+15	5.37184E+15	5.16152E+15	4 • 84457E+15	4.41803E+15	3.87781E+15	3.22080E+15	2.44116E+15	1.46695E+15
4.17181E+15	4.090166+15	3.92669E+15	3.68098E+15	3.352208+15	2.939396+15	2.44160E+15	1.85536E+15	1.12431E+15
2.520618+15	2.47040E+15	2.36991E+15	2.21908E+15	2.01791E+15	1.76682E+15	1.46743E+15	1.12446E+15	7.53468E+14
1.485966+15	1.45590E+15	1.395dlE+15	1.30588E+15	1.186568+15	1.039046+15	8.66206E+14	6.74924E+14	4.81670E+14
8.7193E+14	8.54135E+14	8.18533E+14	7.65440E+14	6.95452E+14	6.09849E+14	5.11269E+14	4.C4903E+14	3.00176E+14
5.05249E+14	4.94828E+14	4.74086E+14	4.43280E+14	4.02941E+14	3.541010+14	2.98646E+14	2.39784E+14	1.82339E+14
2.81912E+14	2.76073E+14	2.644765+14	2.47317E+14	2.24981E+14	1.98161E+14	1.68020E+14	1.36341E+14	1.05491E+14
1.38623E+14	1.35747E+14	1.30043E+14	1.21628E+14	1.10720E+14	9.76937E+13	8.31472E+13	6.79373E+13	5.31246E+13
	IZCN TAL PLANE 1 8.46611E+15 8.33904E+15 8.04307E+15 7.74585E+15 7.3C844E+15 6.53015F+15 5.47679E+15 4.17181E+15 2.52061E+15 1.48546E+15 8.71983E+14 5.05249E+14 1.38623E+14	IZCNTAL PLANEK=12 $1$ $2$ $8.46611E+15$ $8.3394E+15$ $8.3394E+15$ $8.33904E+15$ $8.21122E+15$ $8.09307E+15$ $7.96479E+15$ $7.74585E+15$ $7.61678E+15$ $7.30844E+15$ $7.17917E+15$ $6.53015F+15$ $6.40899E+15$ $5.47679E+15$ $5.37184E+15$ $4.17181E+15$ $4.09016E+15$ $2.52061E+15$ $2.47040E+15$ $8.71932E+14$ $8.54135E+14$ $5.05249E+14$ $4.94828E+14$ $2.81912E+14$ $2.76073E+14$ $1.38623E+14$ $1.35747E+14$	IZCNTAL PLANEK=123 $1$ $2$ $3$ $8.46011E+15$ $8.33943E+15$ $8.09129E+15$ $8.33904E+15$ $8.21122E+15$ $7.96363E+15$ $8.09307E+15$ $7.96479E+15$ $7.71553E+15$ $7.74585E+15$ $7.61678E+15$ $7.36377E+15$ $7.30844E+15$ $7.17917E+15$ $6.92173E+15$ $6.53015E+15$ $6.40899E+15$ $6.16632E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.17101E+15$ $4.09016E+15$ $3.92669E+15$ $2.52061E+15$ $2.470401E+15$ $2.36991E+15$ $1.40596E+15$ $1.45590E+15$ $1.395d1E+15$ $8.71943E+14$ $8.54135E+14$ $8.18533E+14$ $5.05249E+14$ $4.94828E+14$ $4.74086E+14$ $2.81912E+14$ $2.76073E+14$ $2.64476E+14$ $1.38623E+14$ $1.35747E+14$ $1.30043E+14$	IZCN TAL PLANEK=12341234 $8.46011E+15$ $8.33043E+15$ $8.09129E+15$ $7.74263E+15$ $8.33904E+15$ $8.21122E+15$ $7.96363E+15$ $7.61418E+15$ $8.09307E+15$ $7.96479E+15$ $7.71553E+15$ $7.36233E+15$ $7.74585E+15$ $7.61678E+15$ $7.36377E+15$ $6.999954E+15$ $7.3C844E+15$ $7.17917E+15$ $6.92173E+15$ $6.53584E+15$ $6.53015E+15$ $6.40899E+15$ $6.16632E+15$ $5.80007E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.17101E+15$ $4.09016E+15$ $3.92669E+15$ $3.68098E+15$ $2.52061E+15$ $2.47040E+15$ $2.36991E+15$ $2.21708E+15$ $1.40596E+15$ $1.45590E+15$ $1.395d1E+15$ $1.30588E+15$ $8.71943E+14$ $8.54135E+14$ $8.18533E+14$ $7.65440E+14$ $2.81912E+14$ $2.76073E+14$ $2.64476E+14$ $2.47317E+14$ $1.38623E+14$ $1.35747E+14$ $1.30043E+14$ $1.21628E+14$	IZCNTAL PLANEK=1345 $1$ $2$ $3$ $4$ $5$ $8.46011E+15$ $8.33043E+15$ $8.09129E+15$ $7.74263E+15$ $7.30375E+15$ $8.33904E+15$ $8.21122E+15$ $7.96363E+15$ $7.61418E+15$ $7.17509E+15$ $8.09307E+15$ $7.96479E+15$ $7.71553E+15$ $7.36233E+15$ $6.91880E+15$ $7.74585E+15$ $7.61678E+15$ $7.36377E+15$ $6.99954E+15$ $6.53436E+15$ $7.30375E+15$ $6.1678E+15$ $7.36377E+15$ $6.99954E+15$ $6.53436E+15$ $7.30311E+15$ $6.40899E+15$ $6.16632E+15$ $5.80007E+15$ $5.30311E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.41803E+15$ $4.17181E+15$ $4.09016E+15$ $3.92669E+15$ $3.68098E+15$ $3.35220E+15$ $2.52061E+15$ $2.47040E+15$ $1.39561E+15$ $2.21708E+15$ $2.01791E+15$ $1.48596E+15$ $1.45590E+15$ $1.39561E+15$ $1.30588E+15$ $1.18656E+15$ $8.71943E+14$ $8.54135E+14$ $8.18533E+14$ $7.65440E+14$ $6.95452E+14$ $5.05249E+14$ $4.94828E+14$ $4.74086E+14$ $4.43280E+14$ $4.02941E+14$ $2.81912E+14$ $2.76073E+14$ $2.64476E+14$ $2.47317E+14$ $2.24981E+14$ $1.38623E+14$ $1.35747E+14$ $1.30043E+14$ $1.21628E+14$ $1.10720E+14$	IZCNTAL PLANEK=134561234568.46611E+158.33043E+158.09129E+157.74263E+157.30375E+156.52439E+158.33904E+158.21122E+157.96363E+157.61418E+157.17509E+156.46379E+158.09307E+157.96479E+157.71553E+157.36233E+156.91880E+156.1622E+157.74585E+157.61678E+157.36377E+156.99954E+156.53436E+155.79737E+157.30315E+156.40899E+156.16632E+156.99954E+156.00804E+155.30185E+157.3031E+156.40899E+156.16632E+155.80007E+155.30311E+154.66355E+155.47679E+155.37184E+155.16152E+154.84457E+154.41803E+153.87781E+154.17181E+154.09016E+153.92669E+153.68098E+153.35220E+152.93939E+152.52061E+152.47040E+152.36991E+152.21708E+152.01791E+151.76682E+151.48596E+151.45590E+151.395d1E+151.30588E+151.18656E+151.63904E+158.71943E+148.54135E+148.18533E+147.65440E+146.95452E+146.09849E+142.81912E+142.76073E+142.664476E+142.47317E+142.24981E+141.92161E+141.38623E+141.35747E+141.30043E+141.21628E+141.10720E+149.76937E+13	IZCN TAL PLANEK=1234567 $8.46011E+15$ $8.33043E+15$ $8.09129E+15$ $7.74263E+15$ $7.30375E+15$ $6.52439E+15$ $5.47077E+15$ $8.33904E+15$ $8.21122E+15$ $7.96363E+15$ $7.61418E+15$ $7.17509E+15$ $6.40379E+15$ $5.36632E+15$ $8.09307E+15$ $7.96479E+15$ $7.71553E+15$ $7.36233E+15$ $6.91880E+15$ $6.16222E+15$ $5.15696E+15$ $7.74585E+15$ $7.61678E+15$ $7.36377E+15$ $6.99954E+15$ $6.53436E+15$ $5.79737E+15$ $4.84124E+15$ $7.3031E+15$ $7.17417E+15$ $6.92173E+15$ $6.53584E+15$ $5.30185E+15$ $4.41600E+15$ $7.3031E+15$ $6.40899E+15$ $6.16632E+15$ $5.80007E+15$ $5.3018E+15$ $4.41600E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.41803E+15$ $3.87781E+15$ $3.22080E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.41803E+15$ $3.87781E+15$ $3.22080E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.41803E+15$ $3.87781E+15$ $3.22080E+15$ $5.2061E+15$ $2.47040E+15$ $3.92669E+15$ $3.68098E+15$ $3.35220E+15$ $2.93939E+15$ $2.44160E+15$ $2.52061E+15$ $2.47040E+15$ $2.36991E+15$ $2.21708E+15$ $2.01791E+15$ $1.76682E+15$ $1.46743E+15$ $1.40596E+15$ $1.45590E+15$ $1.395d1E+15$ $1.30588E+15$ $1.18656E+15$ $1.60294E+15$ $8.71943E+14$ $8.54135$	IZCN TAL PLANEK =12345678 $8.46611E+15$ $8.33043E+15$ $8.09129E+15$ $7.74263E+15$ $7.30375E+15$ $6.52439E+15$ $5.47077E+15$ $4.16639E+15$ $8.33904E+15$ $8.21122E+15$ $7.96363E+15$ $7.61418E+15$ $7.17509E+15$ $6.462379E+15$ $5.36632E+15$ $4.08515E+15$ $8.09307E+15$ $7.96479E+15$ $7.71553E+15$ $7.61418E+15$ $7.17509E+15$ $6.462379E+15$ $5.36632E+15$ $4.08515E+15$ $7.74585E+15$ $7.61678E+15$ $7.36277E+15$ $6.999954E+15$ $6.16222E+15$ $5.15696E+15$ $3.92245E+15$ $7.30844E+15$ $7.17417E+15$ $6.92173E+15$ $6.53584E+15$ $6.00804E+15$ $5.3018E+15$ $4.41600E+15$ $3.35003E+15$ $6.53015F+15$ $6.40899E+15$ $6.16632E+15$ $5.80007E+15$ $5.30311E+15$ $4.66355E+15$ $3.87694E+15$ $2.93820E+15$ $5.47679E+15$ $5.37184E+15$ $5.16152E+15$ $4.84457E+15$ $4.41803E+15$ $3.87781E+15$ $3.22080E+15$ $2.93820E+15$ $5.47679E+15$ $2.47140E+15$ $3.92669E+15$ $3.68098E+15$ $3.35220E+15$ $2.93939E+15$ $2.44116E+15$ $2.52061E+15$ $2.47140E+15$ $2.36991E+15$ $2.21708E+15$ $2.01791E+15$ $1.76682E+15$ $1.46743E+15$ $1.12446E+15$ $2.52061E+15$ $2.47140E+15$ $1.395d1E+15$ $1.30588E+15$ $1.8653E+15$ $1.03904E+15$ $1.46743E+15$ $1.12446E+15$ $2.44116E+15$ $1.395d1E+15$ $1.395d1E+15$ $1.30588E+15$ $1.6090E+1$

10 11 12 13 14 1 1.48348E+15 8.70362E+14 5.04208E+14 2.81286E+14 1.38281E+14 2 1.45358E+15 8.52615E+14 4.93849E+14 2.75483E+14 1.35426E+14 3 1.39381E+15 8.17203E+14 4.73225E+14 2.63954E+14 1.297616+14 4 1.304282+15 7.643622+14 4.425745+14 2.468855+14 1.213962+14 5 1.1854CE+15 6.94649E+14 4.02407E+14 2.24650E+14 1.10543E+14 6 1.03829E+15 6.09307E+14 3.53731E+14 1.97927E+14 9.75710E+13 7 8.65790E+14 5.10944E+14 2.98414E+14 1.67868E+14 8-30703++13 8 6.747376+14 4.04734E+14 2.39655E+14 1.36251E+14 6.78948E+13 9 4.81609E+14 3.00101E+14 1.82276E+14 1.05443E+14 5.31050E+13 10 3.24834E+14 2.10747E+14 1.31849E+14 7.76830E+13 3.97342E+13 11 2.107716+14 1.416466+14 9.11577E+13 5.49837E+13 2.84290E+13 12 1.31875E+14 9.11640E+13 6.01043E+13 3.69411E+13 1.93308E+13 13 7.79058E+13 5.49976E+13 3.69446E+13 2.30467E+13 1.21783E+13 14 3.97407E+13 2.84288E+13 1.93288E+13 1.21755E+13 6.47540E+12

FLUX FCR GROUP 2

HURIZONTAL PLANE K= 1 1 2 3 4 5 6 7 8 7,562856+15 7,45326E+15 7,24023E+15 6,93701E+15 6,54357E+15 5,95868E+15 5,26476E+15 4,56559E+15 3,98460E+15 2 7.45349E+15 7.34371E+15 7.13022E+15 6.82637E+15 6.43298E+15 5.85306E+15 5.16836E+15 4.48032E+15 3.90935E+15 3 7.2413CE+15 7.13105E+15 6.91619E+15 6.60978E+15 6.21405E+15 5.64329E+15 4.97698E+15 4.31127E+15 3.76040E+15 4 6.93920E+15 6.82830E+15 6.61088E+15 6.29762E+15 5.88991E+15 5.33205E+15 4.69385E+15 4.06196E+15 3.54125E+15 5 6.54703E+15 6.43615E+15 6.21638E+15 5.89111E+15 5.45828E+15 4.92354E+15 4.32545E+15 3.73921E+15 3.25833E+15 6 5.963236+15 5.85730E+15 5.64672E+15 5.33437E+15 4.92466E+15 4.43056E+15 3.88596E+15 3.35604E+15 2.92254E+15 5.17324E+15 4.98111E+15 4.69695E+15 5.26996E+15 7 4.32742E+15 3.88685E+15 3.4C477E+15 2.93630E+15 2.55085E+15 8 4·57107E+15 4·48549E+15 4·31576E+15 4·06550E+15 3·74171E+15 3·35753E+15 2·93693E+15 2.52254E+15 2.16579E+15 9 3.99015E+15 3.91461E+15 3.76503E+15 3.54503E+15 3.26115E+15 2.92442E+15 2.55191E+15 2.16622E+15 1.78319E+15 10 3.16639E+15 3.10615E+15 2.98707E+15 2.81233E+15 2.5K741E+15 2.32084E+15 2.C2472E+15 1.71402E+15 1.40317E+15 11 2.36218L+15 2.31/22E+15 2.22846E+15 2.09848E+15 1.93156E+15 1.73405E+15 1.51460E+15 1.28374E+15 1.05264E+15 12 1.65279E+15 1.62139E+15 1.55947E+15 1.46896E+15 1.35297E+15 1.21598E+15 1.C6394E+15 9.03988E+14 7.43764E+14 1.04513E+15 1.02533E+15 9.86315E+14 9.29366E+14 8.56498E+14 7.7C581E+14 6.75334E+14 5.7516CE+14 4.74721E+14 13 14 5-14421E+14 5-04691E+14 4-85540E+14 4-57604E+14 4-21897E+14 3-75838E+14 3-33249E+14 2-84259E+14 2-35102E+14 10 11 12 13 14 1 3.1614CE+15 2.35808E+15 1.64967E+15 1.04306E+15 5.13363E+14 2 3.10142E+15 2.31332E+15 1.61842E+15 1.02335E+15 5.03683E+14 3 2.98285E+15 2.22496E+15 1.55679E+15 9.84525E+14 4.84622E+14 4 2-8C892L+15 2-C9552E+15 1-46667E+15 9-27822E+14 4-56810E+14 5 2.58469E+15 1.92921E+15 1.35112E+15 8.55238E+14 4.21245E+14 6 2.31891E+15 1.73232E+15 1.21459E+15 7.69610E+14 3.79332E+14 7 2.02350E+15 1.51343E+15 1.06296E+15 6.74631E+14 3.32878E+14 1.71336E+15 1.28J03E+15 9.03357E+14 5.74684E+14 8 2.840046+14 9 1.40292E+15 1.05227E+15 7.4340UE+14 4.74423E+14 2.349395+14 10 1.10459E+15 8.31646E+14 5.90197E+14 3.78107E+14 1.87700E+14 11 8.31787E+14 6.29414E+14 4.48936E+14 2.88793E+14 1.437316+14 12 5.90376E+14 4.48996E+14 3.21783E+14 2.07789E+14 1.03661E+14 13 3.78277E+14 2.88874E+14 2.07821E+14 1.34609E+14 6.72830E+13 14 1.87797E+14 1.43781E+14 1.03685E+14 6.72873E+13 3.36765E+13

FLUX FCR GROUP 3

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HUR	IZONTAL PLANE	K= 1							
	1	2	3	4	5	6	7	8	9
1	7.210766+15	7.07446E+15	6.79610E+15	6.36149E+15	5.70520E+15	4.995676+15	4.20445E+15	3.27838E+15	1.98829E+15
2	7.07461E+15	6.94051E+15	6.66685E+15	6.24001E+15	5.59620£+15	4.9CC09E+15	4.12381E+15	3.21528E+15	1.94974E+15
3	6.79684E+15	6.66742E+15	6.40372E+15	5.99347E+15	5.37628E+15	4.7CE54E+15	3.96286E+15	3.08961E+15	1.87301E+15
4	6.36299E+15	6.24133E+15	5.99421E+15	5.61186E+15	5.040546+15	4.42C42E+15	3.72256E+15	2.90259E+15	1.75887E+15
5	5.70753E+15	5.59832E+15	5.37784E+15	5.04134E+15	4.57953E+15	4.03681E+15	3.40512E+15	2.65619E+15	1.60841E+15
6	4.99058E+15	4.90288E+15	4.71079E+15	4.42193E+15	4.03753E+15	3.56922E+15	3.C1454E+15	2.35246E+15	1.42263E+15
7	4.20773E+15	4.12687E+15	3.96544E+15	3.72448E+15	3.40632E+15	3.01507E+15	2.54859E+15	1.9898CE+15	1.20112E+15
8	3.28139E+15	3.21810E+15	3.09205E+15	2.90450E+15	2.65750E+15	2.35321E+15	1.99010E+15	1.55558E+15	9.40753E+14
9	1.990296+15	1.95162E+15	1.87466E+15	1.76019E+15	1.60937E+15	1.42324E+15	1.20145E+15	9.40865E+14	6.37431E+14
10	1.18518C+15	1.16202E+15	1.11597E+15	1.04758E+15	9.57678E+14	8.47297E+14	7.17696E+14	5.71449E+14	4.17771E+14
11	7.1550CE+14	7.01479E+14	6.73639E+14	6.32401E+14	5.78440E+14	5.12762E+14	4.36960E+14	3.54027E+14	2.70227E+14
12	4.28576E+14	4.20178E+14	4.03533E+14	3.78955E+14	3.46964E+14	3.08373E+14	2.64462E+14	2.17332E+14	1.70340E+14
13	2.40932E+14	2.36221E+14	2.26897E+14	2.13168E+14	1.95382E+14	1.74C78E+14	1.50073E+14	1.24577E+14	9.92426E+13
14	1.03668E+14	1.01645E+14	9.76454E+13	9.17684E+13	8.41779E+13	7.51258E+13	6.49816E+13	5.42607E+1 <b>3</b>	4.36109E+13
	10	11	12	13	14				
1	1.183H3E+15	7.14568E+14	4.27950E+14	2.40548E+14	1.03498E+14				
2	1.160752+15	7.00597E+14	4.19585 <u></u> E+14	2.35856E+14	1.01483E+14				
3	1.114846+15	6.72852E+14	4.03000E+14	2.26567E+14	9.749810+13				
4	1.04665E+15	6.31742E+14	3.78502E+14	2.12885E+14	9.164048+13				
5	9.56974E+14	5.77926E+14	3.46603E+14	1.95153E+14	8.40722E+13				
6	8.46814E+14	5.12391E+14	3.08103E+14	1.73904E+14	7.504296+13				
7	7•17403E+14	4.36717E+14	2.64276E+14	1•49948E+14	6.491996+13				
8	5.71302E+14	3.53885E+14	2.17213E+14	1.24494E+14	5.42173E+13				
9	4.177L8E+14	2.70157E+14	1.70273E+14	9.91913E+13	4.35820E+13				
10	2.91301E+14	1.96836E+14	1.27774E+14	7.584136+13	3.366755+13				
11	1.96861E+14	1.37866E+14	9.188126+13	5.55058E+13	2.48870E+13				
12	1.27807E+14	9.18932E+13	6.25216E+13	3.83198E+13	1.73274E+13				
13	7.58698E+13	5.55196E+13	3.83244E+13	2.37372E+13	1.08030E+13				
14	3.36855E+13	2.48975E+13	1.73327E+13	1.08050E+13	4.94004E+12				

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#### POWER DENSITY (MWT/LITER)

HOR	IZON TAL PLANE	κ= 1							
	1	2	3	4	5	6	7	8	9
1	2.780820+00	2.73593E+CC	2.64722E+00	2.51675E+00	2.92140E+00	2.6C687E+0C	2.21773E+C0	1.76164E+00	2.39363E-01
2	2.73696L+00	2.69141E+CC	2.60316E+00	2.473466+00	2.86916±+00	2.55869E+CC	2.17578E+CO	1.72782E+00	2.34612E-01
3	2.64766E+00	2.60347E+CC	2.51606E+00	2.38759E+00	2.76484E+00	2.46237E+0C	2.C9202E+C0	1.6604CE+00	2.25102E-01
4	2.51760E+00	2.47417E+CC	2.38799E+D0	2.26068£+00	2.60793E+00	2.31768E+CC	1.96666E+CO	1.55984E+00	2.10819E-01
5	2.922996+00	2.87058E+CC	2.76587E+00	2.60846±+00	2.39425E+00	2.12389E+CC	1.8C018E+C0	1.427C1E+00	1.91753E-01
6	2.6C818E+C0	2.56053E+CC	2.46383E+00	2.31866E+00	2.12436E+00	1.88195E+CC	1.59396E+C0	1.26344E+00	1.67935E-01
7	2.2190CE+00	2.17780E+CC	2.09370E+00	1.96790£+00	1.80095E+00	1.59429E+CC	1.35C29E+C0	1.07132E+0D	1.39511E-01
8	1.763696+00	1.729746+00	1.662045+00	1.56112F+00	1.42788E+00	1.26394E+0C	1.07151E+C0	8.527506-01	1.06924E-01
S	2.397166-01	2.34440E-C1	2.25384E-01	2.11040E-01	1.919086-01	1.6EC29E-C1	1.39556E-C1	1.06939E-01	7.16565E-02
10	1.41318E-01	1.38459E-C1	1.32745E-01	1.24193E-01	1.128450-01	9.881526-02	8.23781E-C2	6.41868E-02	4.580796-02
11	8.29276E-02	8.12302E-C2	7.78444E-02	7.27951E-02	6.61391E-02	5.7558CE-C2	4.86228E-C2	3.85072E-02	2.85474E-02
12	4.80504E-02	4.70593E-C2	4.50866F-02	4.21569E-02	3.832066-02	3.36758E-02	2.84019E-02	2.2804CE-02	1.73408E-02
13	2.68105E-C2	2.62552E-C2	2.51522E-02	2.35204E-02	2.13962E-02	1.88456E-02	1.59791E-C2	1.29663E-02	1.00324E-02
14	1.31833E-02	1.29098E-C2	1.23674E-02	1.15671E-02	1.05297E-02	9.29C89E-03	7.907495-03	6.46099E-03	5.05227E-03

	10	11	12	13	14
1	1.410936-01	8.27/34E-C2	4.79514L-02	2.675101-02	1.315096-02
2	1.38239E-01	8.10856E-C2	4.69652E-02	2.61991E-02	1.28793E-02
3	1.325546-01	7.77179E-C2	4.50047E-02	2.51026E-02	1.23405E-02
4	1.24040E-01	7.26925E-C2	4.208936-02	2.34744E-02	1.15450-02
5	1.12734E-01	6.60627E-C2	3.826986-02	2.136476-02	1.05129E-02
6	9.87438E-02	5.79465E-C2	3.36407E-02	1.882336-02	9.279230-03
7	8.23386E-02	4.85419E-C2	2.83799E-02	1.59647E-02	7.90017E-03
в	6.41690E-02	3.84911E-C2	2.27917E-02	1.295/8E-02	6.45695E-03
9	4.58021E-02	2.85403E-C2	1.73343E-02	1.00279E-02	5.05041E-03
10	3.C8924E-02	2.00425E-C2	1.25391E-02	7.40685E-03	3.77881E-03
11	2.004488-02	1.34709E-C2	8.66930E-03	5.22955E-03	Z.70367E-03
12	1.25416E-02	8.66990E-C3	5.71610년-03	3.51318E-03	1.83840E-03
13	7.4C901E-03	5.23039E-C3	3.51352E-03	2.19179E-03	1.15819E-03
14	3.77944E-C3	2.70365E-C3	1.83821E-03	1.15792E-03	6.15825E-04

MATER IAL	ATOMIC WT.
U238C	2.38050E+02
PU239C	2.39050F+02
PU24CC	2.400508+02
CARBON	1.20100E+01
NA	2.299C0E+U1
FE C	5.585002+01
U238H	2.38050E+02
PU2398	2.39050E+02
FE 8	5.585002+01
FIS PR	1.00000E+00

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ZUNE =	1	VOLUME= 8.00000E+CO	LITERS
MATERI	AL	KILCGRAMS	
U238	С	2.529596+01	
PU239	c	5.08043E+00	
PU240	C	3.18855E-01	
CARBO	N	3.19054E+00	
N	۸	1+83224E+U0	
FE	С	9.64402E+00	
U238	B	C.	
PU239	В	C.	
FE	В	0.	
FIS P	R	1.32829E-19	

7:3:48= 2	VOLUME = 5.60000E+C1	LITERS
MATER IAL	KILUGRAMS	
U238C	2.21339E+02 ·	
102390	4.44537F+1)1	
PU24CC	2.67838(+00	
CARBON	2.23338E+01	
NA	1.28257E+01	
FE C	6.75C81E+01	
U23&B	С.	
PU2398	С.	
FE B	C.	
FIS PR	9.29800L-19	

ZUNE= 3	VULUME = 2.7900CE+C2	LITERS
MATERIAL	KILUGRAMS	
U238C	0.	
PU239C	С.	
PU24CC	С.	
CAPBON	G •	
NA	0.	
FE C	0.	
U 2 3 88	3.30823E+03	
PU2398	1.10/378-15	
FE B	1.604062+02	
FIS PR	4.63240E-18	

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		ZONE	1	FLUX	=2.0524E+16	VOLUME =8.CC	CCE+CC LITERS	
8URNA8LE ISOTOPE NO.	MATERIAL NO.	NAMĚ		AT UM C ENSITY	FISSION RATE	ABSORPTION RATE	SIG MA FISSION	SIGMA Abscrpticn
1 2 3 4 5 6	1 2 3 7 8 10	U238C PU239C PU24CC 11238B PU2398 FIS PR		8.000E-03 1.600E-03 1.000E-04 0. 0. 1.000E-20	8.865E+16 5.021E+17 7.218E+15 C. 0. 0.	4.739E+17 5.825E+17 1.256E+16 0. 0. 2.856E-01	6.749E-02 1.911E+00 4.396E-C1 3.549E-02 1.824E+00 0.	3.608E-01 2.217t+0C 7.648E-01 2.890E-01 2.073E+00 1.740E-01
		ZONE	2	FLUX	=1.15705+16	VULUME =5.60	DCCE+C1 LITERS	
BURNABLE ISUTGPE NO •	MATERIAL NO.	NAME		AT OM CENSITY	FISSION RATE	ABSUPPTION RATE	SIG MA FISSIUN	SIGMA ABSCRPTICN
1 2 3 4 5 6	1 2 3 7 8 10	ป2380 PU2390 PU2400 U2388 PU2398 FIS PR		1.000E-02 2.000E-03 1.200E-04 0. 0. 1.000E-20	4 • 285E + 17 2 • 457E + 18 3 • 363E + 16 0 • 0 •	2.257E+18 2.838E+18 5.82CE+16 0. 0. 1.1C5E+0C	6.614E-02 1.896E+00 4.326E-01 3.478E-02 1.814E+00 C.	3.484E-01 2.190E+00 7.486E-01 2.805E-01 2.056E+00 1.706E-01
		ZONE	3	FLUX	=1.5677E+15	VOLUME =2.79	OCE+C2 LITERS	
BURNABLE ISOTOPE NO.	MA TER IAL NO.	NAME		AT GM CENSITY	FISSIDN Rate	A & SURPTION RATE	SIG MA FISSION	SIGMA A8SCRPTICN
1 2 3 4 5 6	1 2 3 7 8 10	U238C PU239C PU240C U2388 PU2398 FIS PR		0. C. 0. 3.000E-02 1.000E-20 1.000E-20	0. 0. 2.894E+17 7.728E+C0 0.	0. 0. 3.169E+18 8.694E+0C 7.241E-01	4.195E-02 1.822E+00 2.919E-01 2.206E-02 1.767E+00 0.	2.867E-01 2.079E+00 5.828E-01 2.415E-01 1.988E+00 1.656E-01
		2CNE= 1	BREE	DING RATIC=	•1142			
		ZCNE= 2	вкее	DING RATIC=	•5417			
		ZCNE= 3	RREE	CING RATIO=	.8418			

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	MIXTURE NUMBER	MIX CEMMAND	MATERIAL ATOMIC DENSITY
1	11	С	0.
2	11	1	.78460912E-02
3	11	2	15331915E-02
4	11	3	12117540E-03
5	11	4	.2000C000E-01
6	11	5	.6000000E-02
7	11	6	.1300000E-01
8	11	10	.19091724E-03
9	12	с	0.
10	12	1	•98953719E-02
11	12	2	.19540337E-02
12	12	3	-1346h536E-03
13	12	- 4	2000000E-01
14	12	5	.6000000E-02
15	12	6	• 1 3000000 E-0 1
16	12	10	-13434613E-03
17	13	C C	0.
1.8	13	7	299703735-01
19	13	פ	-26812362E-04
20	13	0	-62000000E=02
20	13	10	28015673E=05
21	10	10	•20013013C-03

TIME (MINUTES)	CUTER I TERATIONS	IN. IT. 30 LCCP	IN. IT. 2D LCOP	E IGENVALUE SLOPE	EIGENVALUE	LAMBDA
1.68	0	0	0	0.	.10456522E+C1	0.
1.80	1	6	174	0.	.10362425E+C1	1C362425E+C1
1.87	2	4	1 30	0.	.10356349E+C1	•99941370Ľ+CC
1.92	3	3	102	0.	.10351833E+C1	•99956389E+CC
1.98	4	3	116	0.	.10347714E+C1	.99960211E+CC
2.05	5	3	129	0.	.10344213E+01	•99966171E+CC
2.11	6	3	133	0.	.10341211E+C1	•99970974E+CC
2.18	7	3	130	0.	.10338677E+C1	•999755C4E+CO
2.24	8	3	112	0.	.10336587E+C1	•99979781E+CC
2.29	9	3	106	0.	.10334829E+C1	•99982988E+CC
2.35	10	3		0.	.10333371E+C1	•99985900E+CC
2.40	11	3	98	0.	.10332156E+C1	•99988238E+CC
2.45	12	3	84	0.	.10331129E+01	•99990064E+CC

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MATERIAL	ATGMIC WT.
U238C	2.38050E+02
PU239C	2.39050E+02
PU24CC	2.40050E+02
CARBON	1.2010CE+01
NΛ	2.29000E+01
FE C	5.58500E+01
U2388	2.38050F+02
PU2398	2.390506+02
FE B	5.58500E+01
FIS PR	1.00CC0E+00

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ZONE= 1	VOLUME = 8.0000CE+CC	LITERS
MATERIAL	KILUGRAMS	
U238C	2.48117E+U1	
PU239C	4.884176+00	
PU24CC	3.86374E-01	
CAK80N	3.19C54E+00	
NA	1.83224E+00	
FE C	9.644C2E+00	
U238N	C.	
P U2 3 95	С.	
FE B	C.	
FIS PR	2.53593£-03	

ZONE= 2	VULUME = 5.6000CE+C1	LITERS
MATERIAL	KILUGRAMS	
U238C	2.19023E+02	
PU239C	4.343218+01	
PU24CC	3.00616E+00	
CARBON	2.233386+01	
NA	1.28257E+01	
FE C	6.75C81E+01	
U238B	C.	
IU2398	C •	
FE B	G.	
FIS PR	1.24915E-02	

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ZUNE= 3	VULUME = 2.79000E+C2	LITERS
MATERIAL	KILOGRAMS	
U238C	С.	
PU239C	C.	
PU24CC	С.	
CARBON	0.	
NA	С.	
FE C	0.	
U238B	3•30496E+IJ3	
PU239B	2.96915E+00	
FE 8	1.60406E+D2	
FIS PR	1.29780E-03	