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*Intensity and Correlated Energy of
Gamma Rays Emitted in the
Spontaneous Fission of Californium-252*

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LA-9408-T
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UC-34c
Issued: June 1982

Multiplicity and Correlated Energy of Gamma Rays Emitted in the Spontaneous Fission of Californium-252

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ABSTRACT

An array of eight high-speed plastic scintillation detectors has been used to infer a mathematical model for the emission multiplicity of prompt gammas in the spontaneous fission of ^{252}Cf . Exceptional time resolution and coincidence capability permitted the separation of gammas from fast neutrons over a flight path of approximately 10 cm. About 20 different distribution models were tested; the best (termed "double Poisson") was

$$\left[\begin{array}{l} \text{Prob of emitting} \\ \text{exactly } n \text{ gammas} \end{array} \right] = \frac{A(B)^n e^{-B}}{n!} + \frac{(1-A)(C)^n e^{-C}}{n!}, \quad A < 1.$$

For a gamma threshold of 140 keV, the model with parameters evaluated reads

$$\pi(n) = 0.675 \frac{(6.78)^n}{n!} e^{-6.78} + 0.325 \frac{(9.92)^n}{n!} e^{-9.92}.$$

The average energy of the prompt gammas is inversely related to the number emitted; however, this inverse relationship is not strong and the total gamma energy does increase with increasing gamma number.

An extension of the experiment incorporated a lithium-drifted germanium gamma spectrometer that resolved nearly 100 discrete gammas associated with fission. Of these gammas, some were preferentially associated with fission in which few gammas were emitted. Certain others were more frequent when many gammas were emitted.

MULTIPLICITY AND CORRELATED ENERGY OF GAMMA RAYS
EMITTED IN THE SPONTANEOUS FISSION OF CALIFORNIUM-252

by

Glenn Samuel Brunson, Jr.

INTRODUCTION

The general objective in this study of gamma-ray multiplicity was to add to knowledge of a relatively unexplored aspect of the fission phenomenon. As will appear later, almost no work has been done previously toward determining the probability distribution of the number of gammas emitted in the spontaneous fission of ^{252}Cf .

Measurement of gamma multiplicity is more complicated than the measurement of neutron multiplicity although the two are similar. Neutron multiplicity measurements are customarily made at the center of a single large, hollow scintillation detector containing of the order of 500 liters of liquid scintillator that is viewed by a number of photomultipliers operated in parallel.[1] The scintillator is loaded with a strong neutron absorber such as gadolinium. Prompt (fast) neutrons given off in fission are down-scattered in energy by the hydrogenous liquid until they are lost from the system or captured by the gadolinium. The high energy gammas from neutron capture are detected in the scintillator and recorded as neutron events.

The detector is gated off to ignore the "gamma flash" that occurs when the prompt fission gammas simultaneously strike the scintillator within about one nanosecond after fission. By contrast, the neutrons require several microseconds for thermalization and capture. This slowing down time effectively separates the neutron capture events from each other. Thus, the neutron counter, with suitable gating, can

ignore the prompt gammas and then count several (possibly all) neutrons from a given fission; the dispersion in capture times is large enough that deadtime corrections are small. A typical neutron counter for this purpose may achieve an efficiency of 0.80 or higher.

Counting fission gammas is quite different. Here there can be no time separation because all the prompt gammas arrive at the detector at virtually the same instant. Because there is no time separation to permit the counting of individual gammas, the gamma multiplicity experiment must depend on spatial separation. This means having a substantial number of independent gamma detectors. Geometrically we would like to come, in the aggregate, as close to a 4π detector as possible. Detecting an n-fold coincidence depends, in the first approximation, on the total counter sensitivity to the nth power.

In the case of the single neutron counter discussed above, it is actually possible to count all the neutrons at the highest credible multiplicity, say nine. It is not possible, however, to count all the gammas from a single fission without a quite large number of detectors. If two or more gammas from a fission strike the same detector, we will count them as a single gamma event.

Various values for the average number of gammas emitted in fission of ^{252}Cf have been reported in the literature.[2-6] The number ranges from approximately 7 to 10. The literature on this point is reviewed in detail in a separate section. Here is an indication of the wide range of this work beyond the measurements of average neutron and gamma yields:

1. Rumpold et al.[8] have been able to associate specific gamma energies with fission fragment energy and with estimated half-lives for the decay in question.
2. Val'skii et al.[4] in Russia have investigated the gamma-ray yield as a function of total fragment kinetic energy.
3. Lajtai et al.[9] in Hungary and Skarsvag[10] in Norway have studied the angular distribution of gammas from ^{252}Cf fission relative to the axis of separation of the two fission fragments.
4. Nifenecker et al.[11] in France have found a positive correlation between the number of neutrons emitted and the total gamma energy.
5. Ramamurthy et al.[12] in India have studied the multiplicity of prompt fission gammas from ^{252}Cf . Only the Ramamurthy group seems to have studied the gamma multiplicity explicitly; their experiment is reviewed in the next section.

We have used eight high speed plastic scintillators and a coincidence criterion of about four nanoseconds. Because eight detectors cannot detect the maximum number of gammas that may be emitted in a fission, our approach is simply to construct a mathematical model of the emission multiplicity distribution that is most consistent with the detection multiplicities actually recorded.

We suppose the model to have the following characteristics:

1. $\sum_{n=0}^{\infty} \pi(n) = 1$, where $\pi(n)$ = probability that a fission emits exactly n gammas.

2. $\sum_{n=0}^{\infty} n\pi(n) = \bar{n}$, where $\bar{n}$ = average number of gammas emitted.
3. The distribution is nonnegative.
4. The distribution is in some sense smooth.

We seek an analytical expression for $\pi(n)$, the emission probability distribution, which can be reconciled most closely with $P(M)$, the detected multiplicity distribution. The criterion for judging the best model is a weighted least squares comparison with experimental data.

Once the counter array was constructed and operating, we were able to extend the investigation to the relationship between the number of gammas emitted and the average energy of those gammas.

The experiment depends critically on timing. Practical considerations led to an apparatus with eight gamma scintillation detectors (127 x 127 mm) clustered as tightly as possible around the source. This resulted in a source-to-detector distance of only 10 cm. Over this flight path, gammas are separated from the most energetic fission neutrons by about three nanoseconds, hence experimental timing must achieve time definition that is a fraction of this minimum separation.

I

REVIEW OF RELEVANT LITERATURE

There are five available measurements of $\bar{n}$, the average number of gammas emitted in the spontaneous fission of ^{252}Cf [2-6] and only one reference dealing with the characteristics of the emission probability distribution.[12] The measurements of $\bar{n}$ are summarized in Table 1 insofar as they can be determined from the references.

It has been customary to treat the average gamma yield $\bar{n}$ in the same fashion as $\bar{\nu}$, the average prompt neutron yield. That is, it is tacitly assumed that all gammas are, in principle, detectable just as all neutrons are detectable. In fact, $\bar{n}$ is clearly a function of the gamma detector threshold, although only Verbinski et al.[5] explicitly recognized that fact. As an example of the usual treatment, Hoffman and Hoffman[7] make a statistical combination of the last three values in the table without noting that they refer to different thresholds.

The dependence on threshold is strong, as seen in Table 1. Entries in the table have been drawn by inference where specific information was not available. For example, the threshold for the first entry (Smith)[2] was inferred approximately from the remark that a small photopeak was detected near 60 keV. The threshold for the second entry was inferred from the left-hand-most point in Figure 6 of Reference 3. Thresholds are given in References 4 and 5. The threshold

TABLE 1

SUMMARY OF AVAILABLE DATA ON AVERAGE NUMBER OF GAMMAS ($\bar{n}$)
EMITTED IN SPONTANEOUS FISSION OF ^{252}Cf

Date	$\bar{n}$	σ	Threshold (keV)	$\bar{E}^a$ (MeV)	First Author	Ref.
1956	10.3	?	~40 ?	0.80	A. B. Smith	2
1958	10	?	~40 ?	0.9	H. R. Bowman	3
1969	7.5	1.5	100		G. V. Val'skii	4
1973	7.8	0.3	140	0.88	Verbinski	5
1974	8.32	0.4	85	0.85	F. Pleasonton	6

^aAverage energy per gamma.

for the unpublished experiment by Pleasanton was obtained from the experimenter. It is important to note that all of these results were obtained using sodium iodide detectors.

Despite the fact that the first two measurements in Table 1 were made well more than 20 years ago with relatively primitive equipment, they are in better agreement with later measurements than it at first appears. In Figure 1, $\bar{n}$ is shown as a function of discriminator threshold. The values of $\bar{n}$ for some higher thresholds have been taken from Table I of the Verbinski paper and included in Figure 1. There is an upward trend with decreasing threshold; this is to be expected as more gammas become "eligible" for detection. Moreover, there is another phenomenon that becomes of interest in the region of 40 keV and below. Glendenin et al.[13] have found x rays in the region of 40 keV that are associated with the heavier of the two fission fragments. Thus, when the threshold is lowered to 40 keV, some x rays will be counted indistinguishably with low energy gammas; this phenomenon may have contributed to the larger values obtained for $\bar{n}$ in the earlier measurements.

As another point of comparison, Smith et al.[2] estimated that there were 5.0 gammas per fission between 0.5 and 2.3 MeV. We can estimate from Table I of Reference 5 that Verbinski et al. found 4.1 gammas in the same energy interval. Also, we note the qualitative agreement between References 5 and 6 with respect to the average energy per fission gamma. Reference 6, with a lower threshold, reports a slightly lower average gamma energy, as it should.

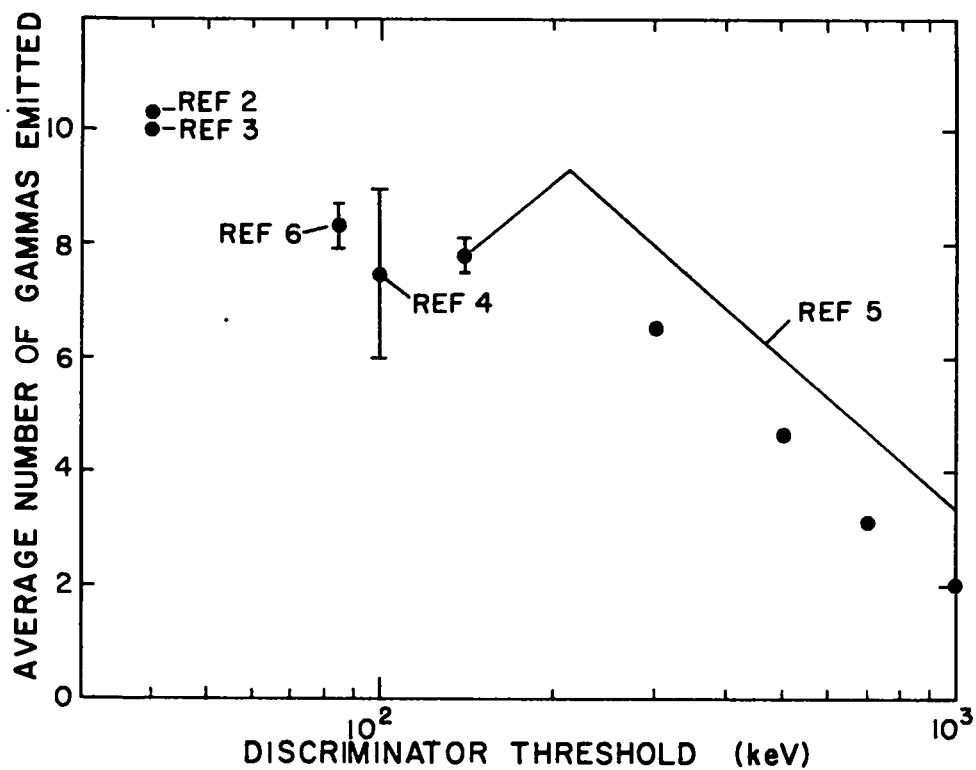


Fig. 1. Graphical summary of available data on the average number of prompt gammas emitted in the spontaneous fission of ^{252}Cf .

We turn now to the work of Ramamurthy et al.,[12] which appears to be the only previous attempt to characterize the emission distribution of gammas from the spontaneous fission of ^{252}Cf . There are several problems with this experiment. The experimenters (in 1977) assumed $\bar{n} = 10.3$, attributing this number to Johansson and Kleinheinz [14] when, in fact, these authors seem merely to have been citing the 1956 results of Smith et al.[2] This overlooks the last three entries in Table 1, all of which were available in 1977. They worked with only three gamma detectors, too few to obtain meaningful statistics when emission multiplicities may easily be 15 or greater.

No mention is made of the gamma threshold nor of the time criterion for coincidence. No mention is made of the threshold level of the initiating fission detector. In the present work, this has been found to be a serious source of bias if not properly adjusted.

In analyzing the results, the Ramamurthy group did not take into account the complexity of the probabilities involved. For example, they write for the singles detection rate per fission in a given counter i

$$C_i = \bar{n} \cdot \Omega_i$$

where Ω_i is the detection efficiency and $\bar{n}$ is the average number of gammas emitted. This does not take into account the fact that the probability of a single detection is actually a compound probability that includes the "miss" probabilities for the other detectors. These "miss" probabilities are not, in general, negligibly close to unity. Moreover, they are nonlinear [of the form $(1 - \Omega_1)^n$] and, therefore, must be evaluated for the emission probability distribution.

Ramamurthy and associates concluded that their data were consistent with a Gaussian gamma emission distribution having $\bar{n} = 10.3$ and a standard deviation of 4.2 ± 0.4 .

II

EXPERIMENTAL EQUIPMENT

The experimental gamma counter array consisted of eight large cylindrical scintillation detectors clustered symmetrically around the californium fission source (see Figure 2). The fission source was electro-deposited on a stainless steel disc and placed in intimate contact with a thin scintillator on a small photomultiplier. This fission detector provided a timing signal on the occurrence of a fission. Each such signal initiated (if the system was in the quiescent state) an analysis cycle that counted the coincident gammas in the surrounding eight detectors.

The scintillator on the fission detector was of Pilot U.[15] It was 0.25 mm thick by about 20 mm in diameter, and in intimate contact with the source foil that carries the ^{252}Cf deposit as a 6 mm circular spot at its center. The source strength was about 800 fissions per second. The other side of the scintillator was coupled to an RCA 4886 photomultiplier[16] (3/4 inch nominal diameter) with silicone optical grease. This tube has a photocathode about 13 mm in diameter. The fission fragments had flight paths on the order of microns and impinged on the scintillator in a spot comparable in size to that of the active deposit. This means that most photoelectrons originated in a similarly sized spot on the photocathode, and the envelope of their trajectories

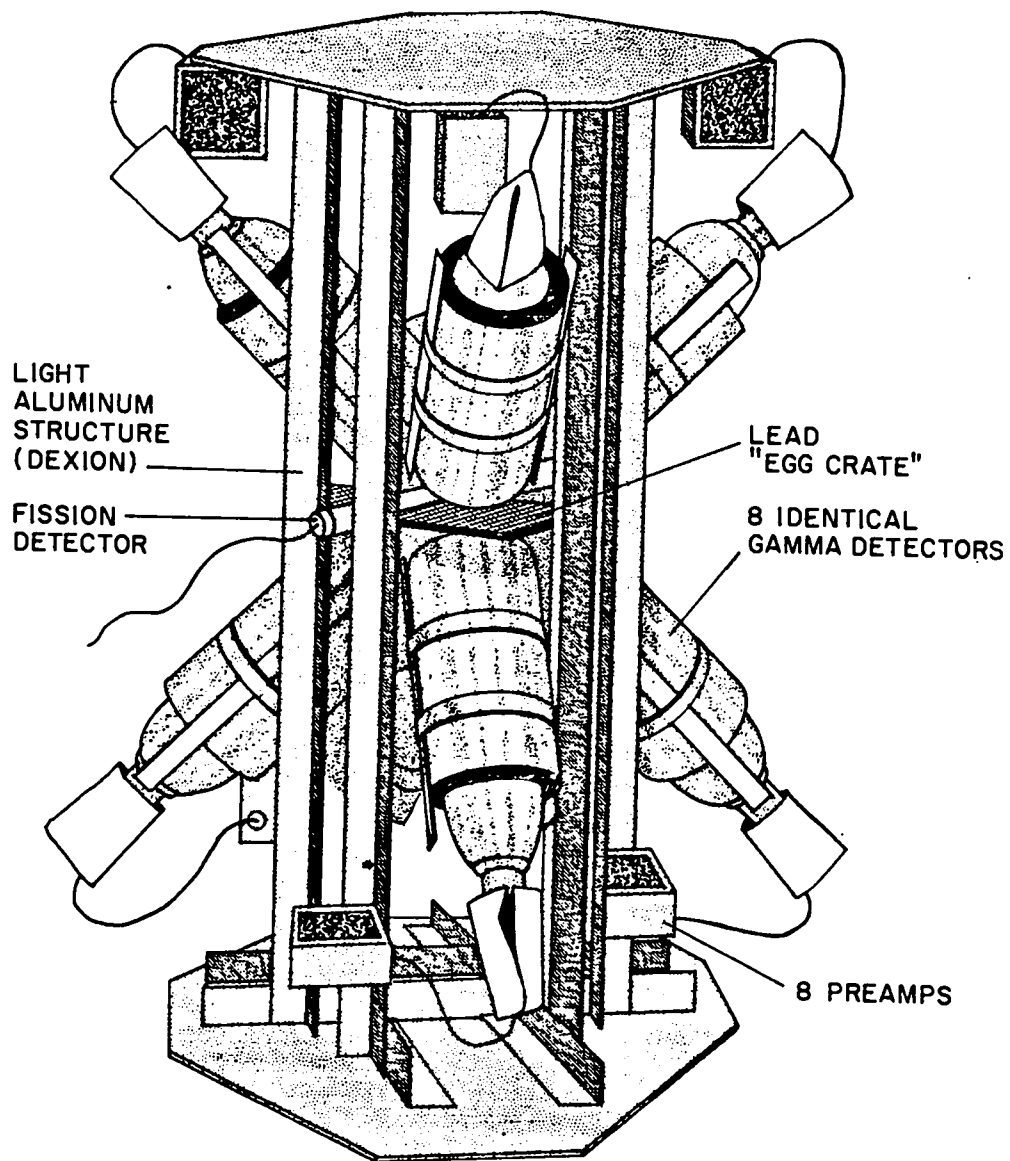


Fig. 2. Perspective of the gamma detector array. The small fission detector assembly extends so that the fissionable material is at the center of the counting geometry.

to the first dynode was a narrow cone in which transit times did not differ much. The result was small time dispersion for the trigger pulses originating in the fission detector. To reduce the amount of scattering material near the source, the fission detector was operated without a mu-metal shield. For this reason, much care was required to assure that the fission detector was returned to the same rotational orientation each time it was removed. Otherwise, the varying effects of the earth's magnetic field could cause perceptible changes in the amplitude of the output. The fission detector was operated at -700 volts bias and the signal did not require amplification before input to the discriminator.

In ^{252}Cf , the principal mode decay is by alpha emission (97%). The other 3% is by spontaneous fission.[17] Taking into account two fission fragments per fission, the emission rate of alphas exceeds that of fission fragments by a factor of about 16. Both were detected by the scintillator. The alpha particles from ^{252}Cf are emitted with energies up to 6.1 MeV; the fission fragments have energies in the range 70-100 MeV. Given the relatively poor energy resolution of plastic scintillators, care was necessary to set the discriminator so as to trigger on an average distribution of fission events. Figure 3 shows the spectrum of pulses occurring in the fission detector. The highest peak represents alpha events and the two lower humps correspond to the bimodal distribution of fission fragment energies. The right (high energy) hump corresponds to the lighter fission fragments because they receive the greater energy in the fission process.

It is surprising to find the approximately 6 MeV alpha particles so close to the approximately 70 MeV heavy fission fragments; this is attributable to the differing response of the scintillator to the two distinct particle types. However, with the possible exception of the rare symmetric fission, every fission event must be represented by one fragment in each of the humps of the energy distribution. Thus, if the discriminator threshold is set exactly between the two humps, the experiment will be triggered on a representative distribution of all fission events.

Suppose, on the other hand, the threshold is set relatively high so as to accept only the upper 10% of the pulses in the upper hump of the fission fragment energy distribution. This will bias the triggering events in two ways:

1. It will tend to select fissions in which the mass partition is highly asymmetric because these are the events that produce the most energetic light fission fragments.
2. And of these fissions, the selected events are further biased toward relatively lower excitation energies because low excitation leaves more kinetic energy available to interact with the scintillator.

The lower excitation energy is evidenced by reduced yield of prompt neutrons and total gamma energy. This effect has been measured by Nifenecker[11] and our results are in qualitative agreement with respect to the gammas.

The eight gamma detectors, 127 x 127-mm cylinders of the plastic scintillator NE102,[15] were mounted on Amperex XP2040[18] photomultipliers with silicone optical grease. The scintillators were wrapped

dry with reflective aluminum foil. The photomultipliers were fitted with mu-metal shields. This photomultiplier has 14 stages and, when operated at -2250 volts bias voltage, did not require amplification before the signal was fed to the discriminator (Phillips 715).[19]

The gamma detectors were symmetrically supported around the fission source on light aluminum structural members (Dexion). Pairs of detectors faced each other along, and were coaxial with, each of the four diagonals of a cube. All were brought toward the center of the cube until their edges touched. The resulting enclosed space was conceptually a regular octahedron with about 200 mm between the faces of opposing scintillators. Each of the scintillators was wrapped on its curved surface with 1.6 mm of lead sheet. In addition, there was an "egg crate" of 3.2-mm lead plates that formed a separate cell for each detector. This egg crate was cut away at the center of the counting geometry to provide space for the fission counter assembly.

The purpose of the lead shielding was to minimize "cross talk" between detectors. It was anticipated (and observed) that, without shielding, a single gamma could, by a Compton or other event, give a spurious coincidence between adjacent detectors. Subsequent measurements indicated that the lead described above substantially reduced the effect in adjacent detectors.

Clearly, if we were to discriminate between gammas and fast neutrons on a flight path of only 10 centimeters, timing was critical everywhere in the experiment. The fission counter as mentioned above contributed a relatively small amount to dispersion in timing. The gamma detectors did exhibit some dispersion from various causes.

The photocathode of the XP2040 is about 110 mm in diameter and is spherical in form to equalize electron flight times to the first dynode. Nevertheless, there is a transit time difference of about one nanosecond between the center and the edge of the photocathode. Even though a given gamma event will create photoelectrons from many points of the photocathode, this transit time difference tends to vary the pulse shape and hence timing, depending on where in the scintillator the event occurred.

Another factor contributing to time dispersion in the gamma detectors is the velocity difference for visible light and gamma rays in the scintillator material. A gamma ray can traverse the thickness of the 127 mm scintillator in about 0.42 nanoseconds while visible light required about 0.67 nanoseconds to travel the same distance. (The index of refraction is 1.58.)[15] Thus, everything else remaining the same, there can be as much as 1/4 nanosecond difference in the time of creation of the first photoelectron depending only on whether the gamma interaction with the scintillator occurred near the front face or near the photocathode. In addition, the shape of the signal pulse and hence the reaction time of the discriminator will also vary with the location of the scintillation event because of the varying paths the photons can follow before interacting with the photocathode.

The composite time resolution as measured with a time to pulse height converter from a fission counter trigger to an associated fission gamma event in one of the gamma detectors was about 1.1 nanoseconds full width at half maximum (see Figure 4). The time scale was

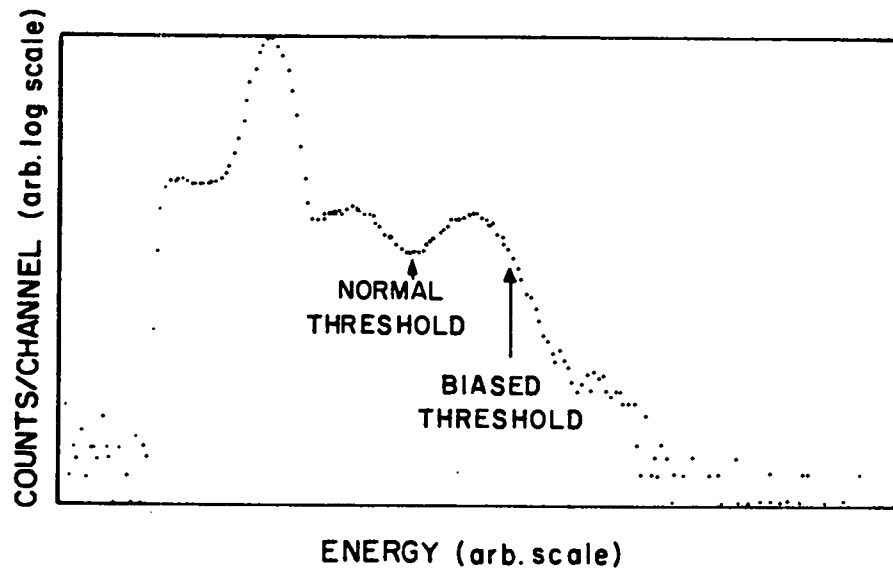


Fig. 3. Spectrum of pulse heights from the fission detector. The arrows indicate the location of the two threshold pulse heights referred to in the text.

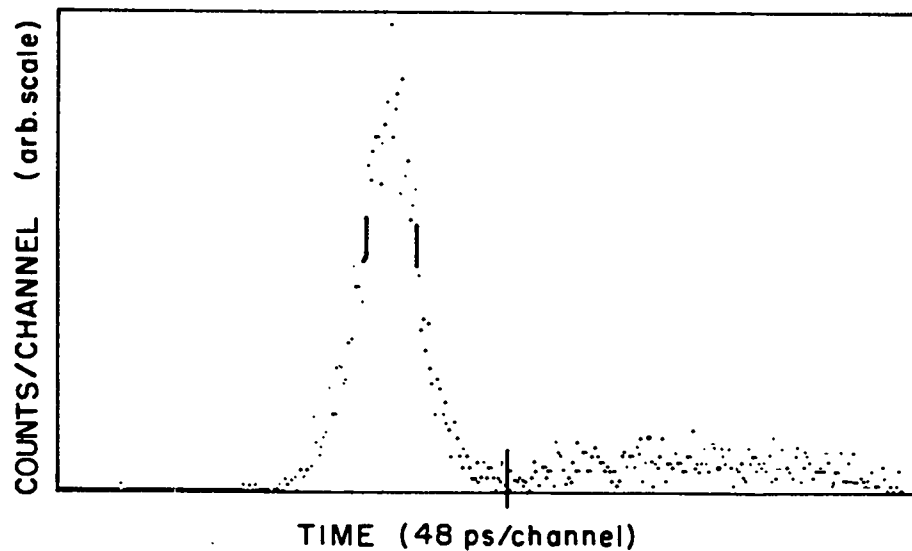


Fig. 4. Time distribution of pulses in a typical gamma detector. This was obtained by triggering a time-to-pulse height converter with a fission detector pulse and observing the arrival of gammas and fast neutrons at one of the eight gamma detectors. The upper two markers indicate a resolution of 1 nsec at half maximum. The lower marker indicates the approximate point of partition between gammas and fast neutrons.

approximately 48 picoseconds per channel and the markers on the peak are separated by approximately 1 nanosecond. The marker on the baseline indicates the approximate boundary between gammas and fast neutrons.

The experiment depended critically on the characteristics of the discriminators. The Phillips Model 715 constant fraction discriminator contains five independent channels in a single NIM module. The threshold has a range from -25 to -1000 mV. The module also has a VETO input that is common to all five channels; this permits an external signal to turn off the channels while analysis of the last event is going on. These discriminators, due to their constant fraction characteristic, have exceptional timing performance. They are said to have less than 150 picoseconds "walk" (time shift) for inputs from -25 mV to -2.5 V. This is predicated on constant pulse shape, which as mentioned above, cannot be assumed under our operating conditions. The minimum width of the discriminator output is 5 nanoseconds and the rise and fall times are approximately 1 nanosecond. The minimum pulse pair separation is 9.5 nanoseconds if the output is at minimum width. There is also a 9.5-nanosecond delay between input and output. The output is -800 mV and compatible with system requirements for counting and gating.

Aside from the detectors and associated power supplied, the experimental equipment (schematic in Figure 5) consisted of standard nuclear electronic modules of two general types, NIM and CAMAC. NIM modules are plugged into a NIM "bin" and partake of a common power source, but all signal and data handling is accomplished externally by

cables to and from each module. CAMAC modules are plugged into a CAMAC "crate," which likewise provides a common power source for the modules. In addition, the crate provides common data paths ("data ways") for reading and controlling such modules by means of a microprocessor. NIM modules are shown as wide rectangles; CAMAC as narrow ones.

The CAMAC modules were operated in a LeCroy 3500[20] data acquisition system that provides a small CAMAC crate. This system is very versatile and provides many options and features not pertinent to this experiment. It can accommodate up to eight CAMAC modules under control of its microprocessor (INTEL 8085). In addition, there is accessible internally an arithmetic processing unit (Advanced Micro Devices APU9511). Besides program memory, there are 16K 3-byte words of data memory.

We will now describe the data acquisition system with reference to Figure 5.

The eight gamma channels were connected to eight of the ten discriminators available in the two Phillips Model 715 discriminators. The analog gamma signals (with attenuation and delay) were also connected to eight inputs of the LeCroy Model 2249A 12-channel integrating analog-to-digital converter (ADC). The eight gamma discriminator outputs were led to eight contiguous inputs (upper byte or lower byte) of the LeCroy Model 2341A 16-channel coincidence register.[20]

The fission counter was connected to one of the two remaining discriminators. That output goes to gate the coincidence register. Another output goes to trigger the one remaining discriminator. The reason for this will be explained later.

Signal delays in the system are adjusted simply by changing cable length. Only cable RG174 was used for signal handling. This cable has a signal propagation velocity of approximately 20 cm/nanosecond. For comparison, the velocity of light is 30.0 cm/nanosecond in free space, and 19 cm/nanosecond in the plastic scintillator.

On the time scale of interest, there were measurable differences in the eight detectors. The lengths of the cables from the gamma discriminators to the coincidence register were adjusted to bring gamma pulses into simultaneity within $\pm 1/10$ nanosecond (or, equivalently, ± 2 cm of cable length).

III

DATA ACQUISITION

With the LeCroy 3500, data acquisition procedures can be programmed either in FORTRAN or in assembly language. In the interest of reducing deadtime, all of these experiments acquired data through assembly language subroutines. For convenience, the overall data acquisition was managed by a FORTRAN program that read operator input (for example, count time) from the keyboard, cleared data memory, and then called the subroutine for actual data acquisition. At the end of the specified count time, control was returned to the FORTRAN calling program. This then performed preliminary analysis as may have been programmed and wrote out the results on the printer. At program end, the data as well as some interim results remained protected in data memory for the use of a subsequent program. The data acquisition software is presented in detail in Appendix C.

We now describe data acquisition with respect to a single fission event. (Refer to Figure 5.) The data cycle starts from the quiescent condition when a fission occurs in the fission counter and causes a signal pulse greater than the threshold of the associated discriminator. The discriminator has multiple outputs, one of which goes to the interrupt register and signals the waiting program that an analyzable event has occurred. The acquisition program then, by means of the output register, turns on an inhibit signal that shuts off

further input until the inhibit is revoked at the end of data processing. The output of the fission discriminator also triggers a gate that acts as another inhibit. The two inhibit signals are combined by a logical OR (in the fan-in/fan-out module). The reason for the apparent duplication is that the hard wired INHIBIT (through the gate generator) can be applied much more quickly than the software inhibit; on the other hand, it is necessary to have software control of the INHIBIT during initialization and output.

The output of the fission discriminator (a logic pulse of approximately 8 nanoseconds duration) also goes through a delay line to gate the coincidence register. The width and arrival time of the gate signal were adjusted to accept any pulse arriving within 3 nanoseconds after the peak.

All, some, or none of the eight gamma scintillation detectors may have detected a gamma at the instant of fission. After the gate has closed, we are left with a byte of data in the coincidence register. We will refer to this byte as a "bit pattern". This bit pattern contains a "1" corresponding to each channel that had a signal pulse in coincidence with the gate and a "0" for each of the other channels. For example, if only detectors 1, 3, and 4 had coincident gamma signals, the bit pattern would be 0000 1101. This has the numerical value of 13_{10} and the data acquisition program would accordingly increment by 1 data memory 13, indicating one more occurrence of an event having this bit pattern. It requires 256 memory locations (0-255) to accept all the possible values of the bit pattern.

Earlier we alluded to a second discriminator in the fission detector channel. Its purpose is to provide a separately adjustable gate to control the LeCroy 2249A ADC.[20] The gate is about 15 nanoseconds wide to permit proper integration of the analog pulses from the detectors.

At the instant the gates close, we have in the coincidence register a bit pattern as described above. Every "1" in that bit pattern indicates the occurrence of a gamma pulse that met:

- 1) the energy criterion of being greater in amplitude than the gamma discriminator threshold; and
- 2) the time criterion by arriving while the gate was open at the coincidence register.

We have in each of the 8 ADC channels in use a 10-bit number corresponding to the integrated amplitude of the signal that arrived while the ADC was gated on. Because the ADC gate is wide enough to permit integration of the photomultiplier pulses, there is the possibility that a neutron signal will be accepted. Even if there is no signal at all, there is an unavoidable "pedestal" due to the integrating characteristic of the ADC. The pedestals and neutron pulses have small effect, because the data acquisition software reads only those ADC channels corresponding to bona fide gamma events as identified by the above two criteria. The evaluation of these biases will be discussed later.

The program sums (in our hypothetical fission event) the ADC values in channels 1, 3, and 4, (these are 10-bit numbers) and divides by 3 to obtain the average amplitude for the three gammas recorded.

This is also divided by 4 to scale the number down to one byte (8 bits) for economy of handling time and data memory. Suppose further, that the above steps result in the number 48. This scaled average amplitude is stored in the data memory in a block reserved for multiplicity 3. That is, data memory location 48 in block 3 is incremented by 1.

This involved storage operation takes place after each event and the input is protected from additional signals as long as necessary by the inhibit control signal as described earlier. Concurrently, the program counts the number of analysis cycles for later comparison with the scaler count of initiating signals. At the end of data storage, the necessary clear signals are generated and the INHIBIT is removed to wait for the next fission event.

At the end of the specified count time, control returns to the FORTRAN calling program. This then prints out preliminary results and allows the operator to (1) quit, (2) repeat, (3) change count time and repeat, or (4) output a statistical summary of the energy data. This summary includes for each block (multiplicity) a total count, maximum ordinate, average energy (first moment), and the sigma of the distribution. After the option 4 printout, the operator is offered the choice of the first three above options or calling by block number for plotting any of the eight spectra.

Figure 6 is an example of the output obtained from each data run. The options mentioned in the last paragraph appear only on the video display and not on the printout. We will explain the printed data in detail because there are a large number of internal checks to detect inconsistencies. Line 1 gives the number of cycles as

determined by a running tally of the number of times the analysis program was initiated. Also in the first line is the counting time. In this case, the time is to be adjusted by a factor of 10 (it should be 21,627) because a ten times lower than normal clock rate was selected. For comparison with the number of analysis cycles, Line 2 gives the number of gates as counted by a scaler external to the analysis program.

Line 3 gives the number of gamma events counted by each of the eight detectors as determined from the stored bit patterns; (failure of a channel is obvious); and the sum of these eight numbers (total gamma events) is shown in Line 9. Line 4 is the probability of an event per cycle for each detector (Line 3 divided by the number of cycles). Line 5 gives the relative efficiencies of the eight detectors as estimated from the "miss" probabilities (to be explained later). Line 6 gives the product of all eight "miss" probabilities.

Line 7 gives the number of times each multiplicity (0 to 8) was observed. These nine numbers were obtained from the bit patterns; the sum of Line 7 entries is given in Line 11 as a third check on the number of cycles (against Lines 1 and 2). Also, the entries in Line 7 are each multiplied by their respective multiplicities and summed to obtain the total number of gammas (Line 10) for comparison with Line 9.

Line 8 gives the respective probabilities of the multiplicities (0 to 8) obtained by dividing Line 7 by Line 2. These are the results carried over to the program that infers the emission multiplicity.

Line 12 is a meaningless combination of all those cycles in which at least one gamma was detected. However, it will be seen that

the first data entry in Line 12 plus the first entry in Line 7 equals Line 2 as it should (number of cycles with gammas plus the number of cycles without gammas equals the total number of cycles). The entries in the COUNTS column (in Lines 13 through 20) correspond respectively to the last eight entries in Line 7. As mentioned above, Line 7 was obtained from the stored bit pattern; the data under COUNTS was obtained by summing in each of the spectral blocks.

"CMAX" is the maximum ordinate in each of the spectral blocks; the program uses this number to establish the vertical scale on output plots. "EBAR" is the first moment (average energy) of a given block in arbitrary units. "SIGMA" is the root-mean-square deviation about "EBAR". Repeating in other words what was described earlier, each datum stored in spectral block M is actually the average energy of M different signals coming from M different detectors. Thus, in block M, the dispersion about the mean should go about as $1/\sqrt{M}$. Looking at the last column we see that σ for $M = 4$ is about half that for $M = 1$, and that σ for $M = 8$ is about half that for $M = 2$. The difference between blocks is illustrated in Figure 7.

```
[1] TOTAL CYCLES = 5224227. TIME = 2162.7 SECONDS
[2] TOTAL INPUT GATES = 5224227.

CHANNEL DISTR. OF GAMMAS AND GAMMA SENSITIVITY

[3] 1452571. 1463088. 1288786. 1286922. 1480169. 1383126. 1357006. 1325103.
[4] .278045 .280058 .246694 .246337 .283328 .264752 .259752 .253646
[5] .040980 .041331 .035633 .035574 .041904 .038685 .037833 .036799

[6] GAMMAS:(PRQ = .085908)

[7] 521887. 1300161. 1523772. 1106780. 542585. 182100. 41052. 5546. 344.
[8] .099897 .248871 .291674 .211855 .103859 .034857 .007858 .001062 .000066

[9] TOTAL GAMMAS FROM DETECTORS = 11036771.
[10] TOTAL GAMMAS FROM DAT.MTR. = 11036771.
[11] TOTAL CYCLES FROM DAT.MTR. = 5224227.

BLOCK          COUNTS          CMAX          EBAR          SIGMA

[12] 0          4702340.0        75397.0        51.10313        29.35766
[13] 1          1300161.0        26121.0        51.78431        39.90958
[14] 2          1523772.0        25003.0        51.18028        27.77851
[15] 3          1106780.0        20621.0        50.70624        22.26257
[16] 4          542585.0         11413.0        50.43822        18.96043
[17] 5          182100.0         4341.0         50.25068        16.69513
[18] 6           41052.0         1108.0         50.07118        15.14287
[19] 7           5546.0          167.0         50.24432        13.99011
[20] 8            344.0           14.0         49.25581        12.66420
```

Fig. 6. Typical output from data acquisition program.

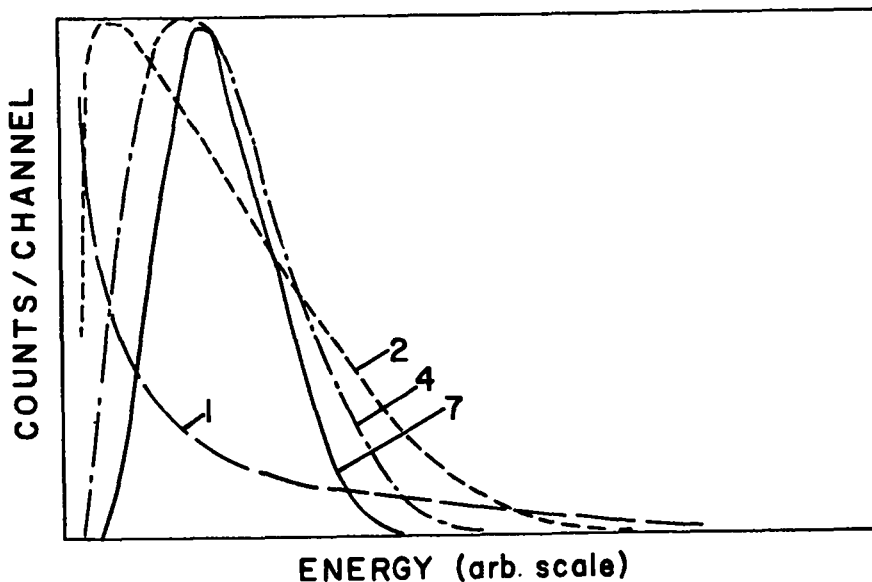


Fig. 7. Comparison of stored spectra for various multiplicities.

IV

OPERATIONAL CONSIDERATIONS

General

We wish to detect "fission" gammas with a high efficiency. Some consideration needs to be given to what is meant by "fission" gammas. The idea of absolute simultaneity falters when we consider the fission phenomenon on the time scale of interest here. For example, Rumpold et al.[8] have been able to classify several specific fission gammas as to whether they are emitted less than 1 picosecond after fission, between 1 picosecond and 2 nanoseconds, or after 2 nanoseconds. Bowman et al.[3] have done a similar experiment and most of the lines they studied, in general not the same as those of Rumpold[4] are assigned half lives less than 1/2 nanosecond. In general, our results accord qualitatively with those of Bowman. It appears that for the purposes of our experiment, effectively all of the fission gammas appear in less than half a nanosecond. This is based on the observation that our time resolution curves taken between the fission counter and one of the gamma detectors was not substantially wider than the time resolution between two of the gamma detectors when the source was ^{60}Co . This source emits two gamma rays that are reported to be connected by a half life of approximately 1 picosecond.

Energy Calibration

Establishing an energy scale for the discriminators on the eight main counting channels is complicated by the poor energy resolution of plastic scintillators. As Nardi[21] points out, the principle gamma interaction with these materials is by Compton scattering. For a mono-energetic source such as ^{137}Cs , the resulting spectrum is a poorly resolved peak corresponding to the energy of the Compton electrons. The indistinct "Compton edge" on the right hand shoulder of this peak corresponds to the maximum energy transfer to an electron in the scintillator. For the 662 keV gamma from ^{137}Cs , this maximum energy is about 480 keV. The ^{137}Cs line was used to define the energy scale in the main scintillator channels. Because Compton scattering[22] occurs with a range of energy transfers from the maximum down to zero, it is possible for any given discriminator setting to miss even high energy gammas. The relationship between the shoulder (Compton edge) of the spectrum is not linear with the peak gamma energy.

The primary emphasis in this experiment is on high efficiency in the gamma detectors. For this reason, as well as the general characteristics of plastic scintillators, we have depended on others' results for values of $\bar{n}$, the average number of gammas emitted in fission of ^{252}Cf . These are the last two entries in Table 1.

Energy Threshold (keV)	$\bar{n}$	Source	Corresponding Compton edge (keV)
140	7.80 ± 0.3	Verbinski[5]	50
85	8.32 ± 0.4	Pleasanton[6]	21

The measurements reported here were all made at a gamma threshold corresponding to one or the other of the two above Compton edge energies.

Timing

The LeCroy Model 2341A coincidence register provides 16 inputs. We used only eight at a time, either the top or bottom half. The gate must precede data pulses by at least 3 nanoseconds. The exact opening time of the gate is not critical but the closing time is. On the time scale of interest here, the transition time of the discriminator output (approximately 1 nanosecond) and response time (turn-off) of the coincidence register are quite significant when we undertake to set the gate so that the effective "turn-off time" is fixed to about 1/2 nanosecond.

All set-up operations involving timing required the use of auxiliary equipment not used in the experiment proper. This consisted of a time to pulse height converter (TPC) that was used in conjunction with the normal pulse height analysis function of the LeCroy system. The fission counter signal (after discrimination) started the time interval that was terminated by the discriminator signal from one of the gamma channels. The result of this analysis is shown in Figure 4. The peak on the left represents the distribution of gamma events in time while the hump on the right similarly represents the fast neutron events. The time scale is approximately 48 picoseconds per channel. By adjusting cable lengths, we brought all eight channels into simultaneity within approximately two channels. (One centimeter of cable corresponds to about one channel.)

Without changing the zero time reference in the TPC, we then brought the gate signal from another output of the fission counter discriminator to the "stop" input of the TPC. By adjusting this cable length, we were able to set the leading edge of the gate at a point some 6 nanoseconds ahead of the gamma peak. The exact position is not critical.

What is critical is the end of the gate because it must be positioned to exclude fast neutrons while losing as few gammas as possible. This sensitive timing depends on the interaction between the fall time of the gate signal and the turn-off characteristic of the coincidence register.

The first step was to bring a copy of the fission counter discriminator pulse from a separate simultaneous output to the "stop" input of the TPC (always maintaining the same zero time reference). By adjusting the cable length, the leading edge of this "decoy" signal was placed exactly at the point on the time scale at which we wanted to close the gate against neutron pulses. With all cable lengths fixed, this decoy signal was led to the coincidence register in place of one or another of the gamma input signals. The regular data acquisition program was run. In the channel with the substitute signal, it was immediately obvious if the gate was missing all the decoy pulses (gate too short) or counting a decoy pulse at every cycle (gate too long.) By trial and error, the duration of the (discriminator output) signal was adjusted to count about half of the decoy pulses. This corresponded to having the trailing edge of the gate at the chosen cut-off position. This method gives a proper cut-off time

taking into account the real, combined characteristics of the various components.

Detector Efficiency

It would be highly misleading to attempt to measure counter efficiency by the usual technique of counting gammas first with the source present and then with it absent. We need the detection efficiency for prompt fission gammas. The straightforward technique of subtracting gross background from gross count would include gammas from two other origins extraneous to our purpose. First, it would include decay gammas from long-lived fission products accumulating since the fabrication of the ^{252}Cf source. Second, it would include the numerous gammas associated with alpha decay of ^{252}Cf . (This is a very substantial source because there are about 30 times as many alpha decays as spontaneous fissions.)

It might be supposed that the gamma counter could be gated by fissions and that could lead to a reasonable measurement of the counter efficiency. One would simply divide the coincident count rate by the trigger rate getting the probability of a count for each fission, then divide that by π to get the efficiency. However, this raises another problem. At the efficiencies involved in this experiment (approximately 0.035), it is easily possible for two or more gammas from the same fission to strike the same detector resulting in a lower estimate of the efficiency than the true value.

There is, however, a relatively clean way to estimate the efficiency of the detectors under the actual conditions of the experiment.

We accept from the literature that ^{252}Cf emits 7.80[5] or 8.32[6] gammas per fission, depending on the threshold, and we assume for this purpose that this is the average of a Poisson distribution.

When a fission occurs and emits n gammas, the probability of not registering a count is the probability that all gammas miss the detector (or traverse it without registering.) This miss probability is:

$$P_{\text{miss}} = (1-\epsilon)^n \quad (1)$$

where ϵ is the efficiency for detecting a single fission gamma at the current discriminator setting. Summing up the P_{miss} for all values of n under the Poisson distribution we have:

$$P_{\text{miss}} = \sum_{n=0}^{\infty} \frac{(\bar{n})^n e^{-\bar{n}}}{n!} (1-\epsilon)^n = e^{-\bar{n}} \sum_{n=0}^{\infty} \frac{[\bar{n} (1-\epsilon)]^n}{n!} = e^{-\bar{n}} \left(e^{\bar{n} (1-\epsilon)} \right) = e^{-\bar{n}\epsilon} \quad (2)$$

$$\text{so} \quad \ln(P_{\text{miss}}) = -\bar{n}\epsilon \quad (3)$$

$$\text{and} \quad \epsilon = \frac{-\ln(P_{\text{miss}})}{\bar{n}} \quad (4)$$

The choice of $\bar{n}$ is not critical because these results have been used only to obtain the relative efficiencies of the eight gamma detectors.

INTERPRETATION OF DATA

The interpretation of the multiplicity data is based on the following equation derived by Van der Werf[23]. Refer to Appendix B. Some symbols have been substituted for convenience. Assuming an array of N identical detectors of average efficiency ϵ counts a shower of G gammas, we have the following probability of triggering exactly M detectors (that is, the emission multiplicity is G and the detection multiplicity is M):*

$$P_N(G,M) = \binom{N}{M} \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} [1 - (N - \ell)\epsilon]^G . \quad (5)$$

*For historic reasons, we have to this point used the symbol $\bar{\pi}$ to represent the average number of gammas emitted in fission. Henceforth, for reasons of clarity and mnemonic convenience, we will use the following symbols:

$\bar{G}$	average number of gammas emitted
G	number of gammas emitted in a specific event, "emission multiplicity"
$\pi(G)$	probability that exactly G gammas are emitted
N	number of gamma detectors in the array
M	number of detectors triggered in a given event, "detection multiplicity"
$P(M)$	probability that exactly M detectors are triggered
M'	number of gammas actually striking the M triggered detectors, "impact multiplicity"

A careful distinction must be made between M the number of detectors triggered and M' the actual number of gammas striking those M detectors; (obviously $G \geq M' \geq M$).

We assume a hypothetical emission distribution (a model) and some starting values for the parameters of the distribution ($C_1, C_2, \dots$)

$$\pi(G) = \pi(G, C_1, C_2, \dots) \quad (6)$$

and calculate the expected probability of triggering exactly M detectors

$$P_N(M) = \sum_{G=M}^{\infty} \pi(G) P_N(G, M) \quad (7)$$

Written explicitly:

$$P_N(M) = \sum_{G=M}^{\infty} \pi(G) \binom{N}{M} \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} [\ell - (N - \ell)\epsilon]^G. \quad (8)$$

We have from the experiment a set of observed probabilities for the distribution of M corrected for false counts caused by scattering.

$$P_N^O(M), \quad (0 \leq M \leq N).$$

$\bar{G}(M)$ average emission multiplicity for a given M

$\bar{M}'(M)$ average impact multiplicity for a given M

$\bar{E}(M)$ average gamma energy per triggered detector for a given M

$\bar{E}'(M)$ average energy per gamma for a given M

ϵ effective average detector efficiency for fission gammas

The numerical procedure is shown in the flow chart, Figure 8. The program adjusts the parameters of the chosen model to obtain a "best fit" by the criterion of minimum weighted least squares. The minimization of the weighted squares is performed by a comprehensive routine written by J. P. Chandler of Indiana University in 1966. The brevity of the flowchart may be misleading. The data analysis program includes many options and the FORTRAN listing requires some 20 pages single spaced. For that reason, it has not been included.

The output from this computation is extensive and includes the experimental data $P_N^O(M)$, the "best fit" results from the model $P(M)$, the absolute difference between the two, and the relative deviation, as well as the converged value of the summed square deviations.

After the problem has converged, the program computes the matrix of the two dimensional probabilities:

$$P_N(G, M) ,$$

the probability that in a given fission exactly G gammas are emitted and exactly M detectors are triggered. Finally, the program calculates an answer to the Bayesian question:

If M detectors were triggered, what was the average number of gammas emitted in the fission event? The calculation is:

$$\bar{G}(M) = \frac{\sum_{G=M}^{\infty} G \cdot P(G, M)}{\sum_{G=M}^{\infty} P(G, M)} . \quad (9)$$

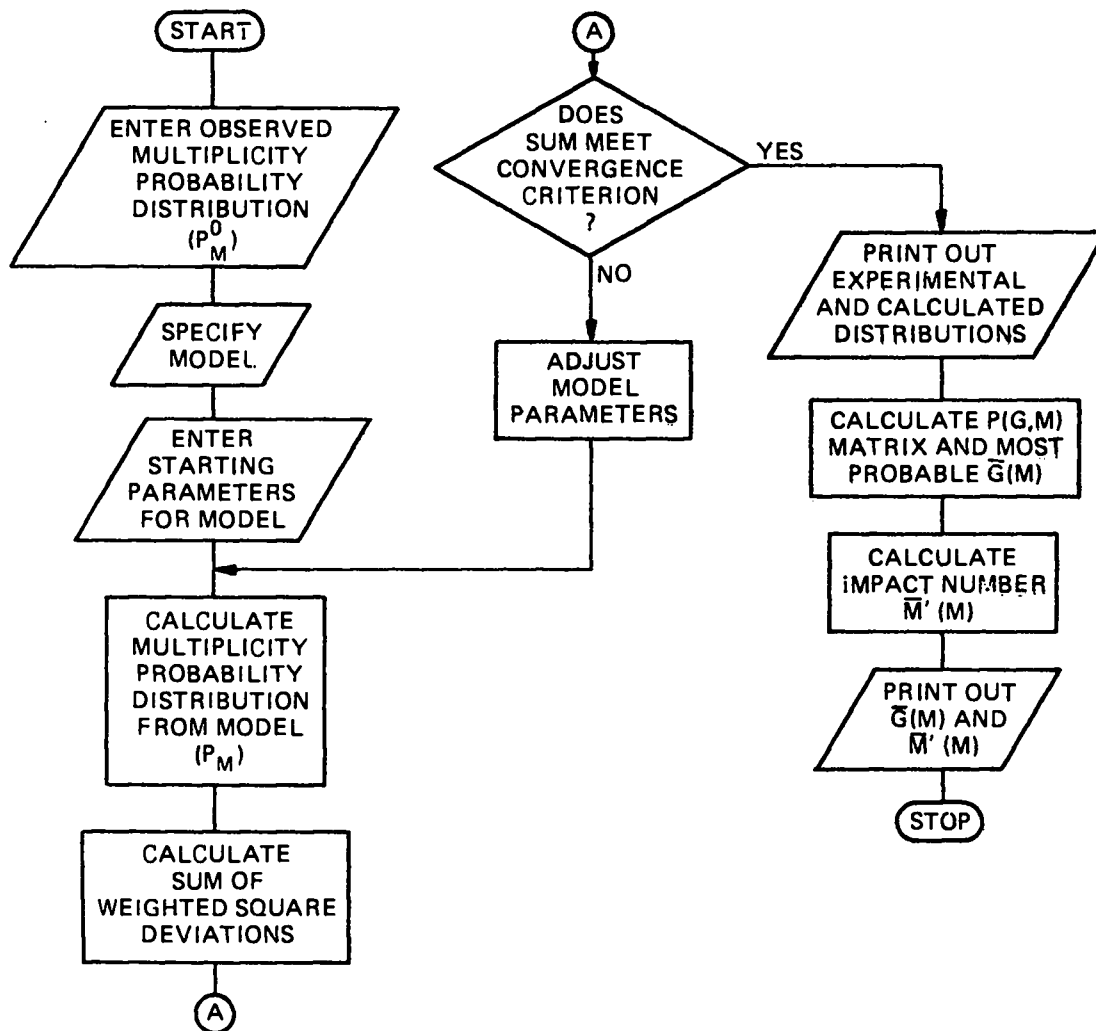


Fig. 8. Flow chart for the data analysis program.

The program also computes the average impact multiplicity, $\bar{M}'(M)$, as a function of detection multiplicity. As shown in Equation 38 in Appendix B, we have

$$P_N(G, M', M) = \binom{N}{M} \binom{G}{M'} (1 - N\epsilon)^{G-M'} \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} (\ell\epsilon)^{M'} \quad , \quad (10)$$

which gives the probability that out of a shower of G gammas on N detectors, exactly M' of the gammas trigger M detectors. Using the inferred emission distribution $\pi(G)$ we can calculate $\bar{M}'(M)$.

$$\bar{M}'(M) = \frac{\sum_{M'=M}^{\infty} \sum_{G=M}^{\infty} M' \cdot \pi(G) \cdot P_N(G, M', M)}{\sum_{M'=M}^{\infty} \sum_{G=M}^{\infty} \pi(G) \cdot P_N(G, M', M)} \quad . \quad (11)$$

The purpose is to achieve a better correlation of energy with multiplicity. As previously described, the data acquisition program yields $\bar{E}(M)$, the total gamma energy of the M -fold event divided by M . However, on the average, this total energy is the response to $\bar{M}'$ gammas that takes into account the fact that detectors frequently experience multiple gamma impacts from the same fission. Hence $\bar{E}$ from the data acquisition program must be corrected to

$$E'(M) = \bar{E}(M) \frac{M}{\bar{M}'(M)} \quad (12)$$

VI

RESULTS

In Table 2 we present six sets of data enumerated below:

Run Number	Gamma Discriminator Threshold (keV)	Average No. of Gammas	Fission Discriminator Setting*
S2	140	7.80	Optimum
S3	140	7.80	Optimum
S4	140	7.80?	Biased
S6	85	8.32	Optimum
S7	85	8.32	Optimum
S8	85	8.32?	Biased

In Table 2 there are typically four entries for each run under each multiplicity:

1. the number of times that multiplicity was observed (run data);
2. the estimated number of events that are "demoted" to the next lower multiplicity;

*"Optimum" refers to a setting between the humps of the fission fragment energy spectrum that is considered to be the best available for triggering on a representative distribution of fission events. "Biased" means a discriminator setting that accepts approximately the upper 10% of all pulses in the high energy hump of the fission fragment spectrum.

TABLE 2

RAW DATA SHOWN WITH CORRECTION FOR FALSE COUNTS AND RESULTING CORRECTED DATA

		----- Detection Multiplicity (M) -----									
Total Events		0	1	2	3	4	5	6	7	8	
<u>Run S2 at 140 keV</u>											
raw data	9,415,495	1,173,314	2,647,957	2,782,169	1,781,587	761,227	221,502	42,617	4,873	248	
correction to M-1				-54,867	-98,332	-76,552	-33,314	-8,558	-1,228	-76	
correction from M+1			54,867	98,332	76,552	33,314	8,558	1,228	76		
corrected data		1,173,314	2,702,824	2,825,634	1,759,807	717,989	196,746	35,287	3,721	172	
<u>Run S3 at 140 keV</u>											
raw data	12,314,425	1,478,350	3,417,192	3,651,645	2,373,541	1,027,015	301,176	58,349	6,775	382	
correction to M-1				-70,806	-129,106	-102,073	-45,005	11,650	-1,684	-105	
correction from M+1			70,806	129,106	102,073	45,005	11,650	1,684	105		
corrected data		1,478,350	3,487,998	3,709,945	2,346,508	969,947	267,821	48,383	5,196	277	
<u>Run S4 at 140 keV</u>											
raw data	3,117,564	429,021	919,589	912,595	554,510	225,474	62,980	11,981	1,344	70	
correction to M-1				-19,216	-32,216	-23,753	-9,815	-2,418	-345	-21	
correction from M+1			19,054	32,216	23,753	9,815	2,418	345	21		
corrected data		429,021	938,643	925,757	546,047	211,536	55,583	9,908	1,020	59	
<u>Run S6 at 140 keV</u>											
raw data	16,341,425	1,628,434	4,060,708	4,772,499	3,465,441	1,695,870	572,080	127,924	17,384	1,089	
correction to M-1				-135,105	-270,940	-236,761	115,917	-33,805	-5,501	-396	
correction from M+1			135,105	270,940	236,761	115,917	33,805	5,501	396		
corrected data		1,628,434	4,195,813	4,908,334	3,431,262	1,575,026	489,968	99,620	12,279	693	
<u>Run S7 at 140 keV</u>											
raw data	5,224,227	521,887	1,300,161	1,523,772	1,106,780	542,585	182,100	41,052	5,546	344	
correction to M-1				-48,258	-86,499	-75,618	-37,098	-10,746	-1,771	-126	
correction from M+1			43,258	86,499	75,618	37,098	10,746	1,771	126		
corrected data		521,887	1,343,419	1,567,013	1,095,899	504,065	155,708	32,077	3,901	218	
<u>Run S8 at 140 keV</u>											
raw data	2,769,793	294,488	713,202	809,830	568,964	271,289	89,365	19,818	2,660	177	
correction to M-1				-23,729	-45,928	-38,765	-18,472	-5,254	-851	-60	
correction from M+1			23,729	45,928	38,765	18,472	5,254	851	60		
corrected data		294,488	736,931	832,029	561,801	250,996	76,147	15,415	1,869	117	

3. the estimated number of events that are "retrieved" from the next higher multiplicity; and
4. the corrected number of times the multiplicity occurred.

Items 2 and 3 above are corrections for gammas that, through a Compton scattering, happen to trigger two detectors, thereby promoting an M-fold event to M+1. The details of this correction are given in Section VII. The last item is entered for each multiplicity in the interpretative computation described earlier.

Multiplicity

We have summarized in Appendix A the many different models tested against the experimental data. The criterion for fitting was the weighted least square deviation between the detection multiplicity data recorded and the detection multiplicity data calculated from a given emission model. Of the many models tried, the following "double Poisson" was the best:

$$\pi(G) = c_2 \frac{(c_3)^G e^{-c_3}}{G!} + (1 - c_2) \frac{(c_4)^G e^{-c_4}}{G!} \quad (13)$$

Next best was the skewed pseudo-Gaussian:

$$\pi(G) = c_1(G)^{c_2} e^{-c_3(G - c_4)^2} \quad G \geq 1$$

$$\pi(0) = c_5 e^{-c_3(-c_4)^2} \quad (14)$$

All other models were vastly inferior. They resulted in values for χ^2 two and three orders of magnitude greater than did the two above models.

The double Poisson was used to obtain all the subsequent reported results except in one case specifically noted. Figure 9 is an example of the output of the program CGA01 that was used to infer the emission gamma distribution. Referring to Figure 9, we see "FINAL VALUES OF X(1)" These are respectively, the efficiency, C2, C3, and C4.

The block marked (1) contains the $\pi(G)$. The columns in block (2) are, respectively:

- a. ("FOLD") - detection multiplicity.
- b. ("EXP") - experimental probability (with previously described correction) of observing that multiplicity.
- c. ("MODEL") - calculated probability (from model) of observing that multiplicity.
- d. ("DEV") - difference between preceding two columns.
- e. ("RDEV") - relative difference.
- f. ("CHI") - difference between model and experiment expressed in sigmas (RMS deviations) on the original experimental data.

"CHISQ" is absolute χ^2 ; "BIG CHI" is χ^2 normalized to the number of events in the run. This is the value to be used for determining goodness of fit. The final two lines of numerical output are, respectively, the average emission number ($\bar{G}(M)$, $M = 1$ to 8, reading left to right) and the average impact number ($\bar{M}'(M)$, $M = 1$ to 8).

The results with the 140 keV discriminator setting are given in Figure 10. It was not possible on this vertical scale to distinguish between the experimental results and the calculated results

```

BEGIN MINIMIZATION. . . .

TERMINATED WHEN THE DIMENSIONS OF THE SIMPLEX BECAME AS SMALL AS THE DELMIN(J).

      289 FUNCTION COMPUTATIONS.

      FINAL VALUE OF CHISQ = .1342339620E-06

      FINAL VALUES OF X(I). . . .
      .3457000000E-01 .6700110865E+00 .6760730498E+01 .9910121346E+01

CHECK ON PTOT      1.000000
EXPECTED N =      7.799988
SIGMA =      3.161151

[1] .000793 .005410 .018542 .042631 .074148 .104408
     .124499 .129945 .121834 .104755 .083937 .063404
     .045479 .031090 .020279 .012619 .007487 .004233
     .002282 .001173 .000575 .000269 .000121 .000052
     .000021 .000008 .000003 .000001 .000000 .000000

SINGLE COUNTER MISS PROBABILITY = .76463851

[2] FOLD      EXP      MODEL      DEV      ROEV      CHI
     0      .1246152 .1245823 -.0000329 -.0002643 -.2863166
     1      .2870613 .2871595 .0000982 .0003420 .5622867
     2      .3001047 .3000159 -.0000888 -.0002961 -.4976595
     3      .1869054 .1869063 .0000009 .0000047 .0062080
     4      .0762561 .0762860 .0000299 .0003919 .3321019
     5      .0208960 .0208971 .0000011 .0000547 .0242460
     6      .0037478 .0037375 -.0000103 -.0027396 -.5146383
     7      .0003952 .0003963 .0000011 .0027955 .1705248
     8      .0000183 .0000189 .0000007 .0357835 .4692969

CHISQ = .1342340E-06

BIG CHI = 1.2638791
WHAT NEXT?? (1)TRY NEW MODEL (2) INPUT NEW
DATA (3)CHANGE EFFICIENCY (4)QUIT (5)SIFT
(6) SIFT END.
? 6

AVERAGE GAMMA IMPUT FOR A GIVEN MULTIPLICITY

      6.579      7.929      9.325      10.747      12.172      13.576      14.943      16.267

IMPACT MULTIPLICITY VS DETECTION MULTIPLICITY
      1.136      2.282      3.441      4.613      5.796      6.986      8.180      9.373

```

Fig. 9. Typical computer output from the data analysis program.

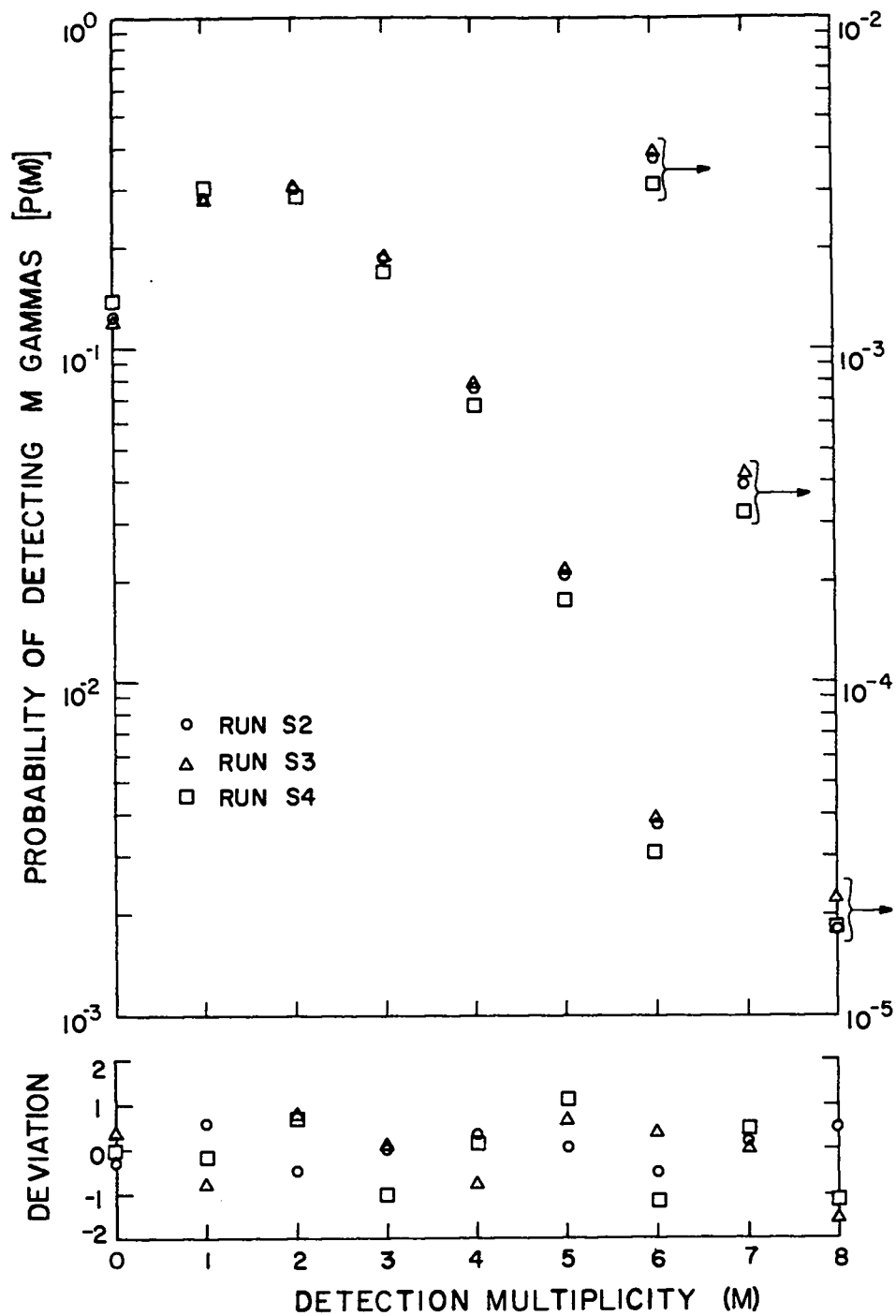


Fig. 10. Experimental results for measurements at 140 keV threshold. The points in the upper graph represent both experiment and model since they differ by too small an amount to show on this scale. The lower graph shows the difference between experiment and model in units of standard deviation.

obtained with the model. Instead, the corresponding deviations (Model-Experiment) are plotted in units of standard deviation in the lower graph.

Figure 11 is the corresponding graph for the data obtained with the 85 keV threshold. The runs with the fission detector at a higher threshold, S4 (Figure 10) and S8 (Figure 11), show relatively lower probabilities at higher multiplicities. This can be attributed to fewer gammas or lower efficiency. Table 3 summarizes the results of fitting the model to the six sets of data.

We start with the nine probabilities for detection multiplicities 0-8 inclusive. This represents only eight degrees of freedom because the probabilities are normalized to unity. Because three parameters are used in the model, the resulting values for χ^2 are considered on the basis of 5 degrees of freedom. From a standard statistical table, we see that χ^2 for 5 degrees of freedom should exceed 6.63 only 25% of the time. In Table 3, χ^2 exceeds 6.63 in only one of the six cases (S6). For S6, χ^2 is 7.49, which by interpolation should be exceeded only 20% of the time; this still indicates a very reasonable fit. Hence, it is unlikely that we can improve the results obtained from this data by choosing a better model. However, it is entirely possible that another experiment using more detectors could improve the model because the greater number of detectors would afford more data points, but not proportionately so.

Figure 12 shows the emission distribution inferred from data sets S2, S3, and S4 (at 140 keV threshold.) Figure 13 gives the same results for S6, S7, and S8 (85 keV threshold).

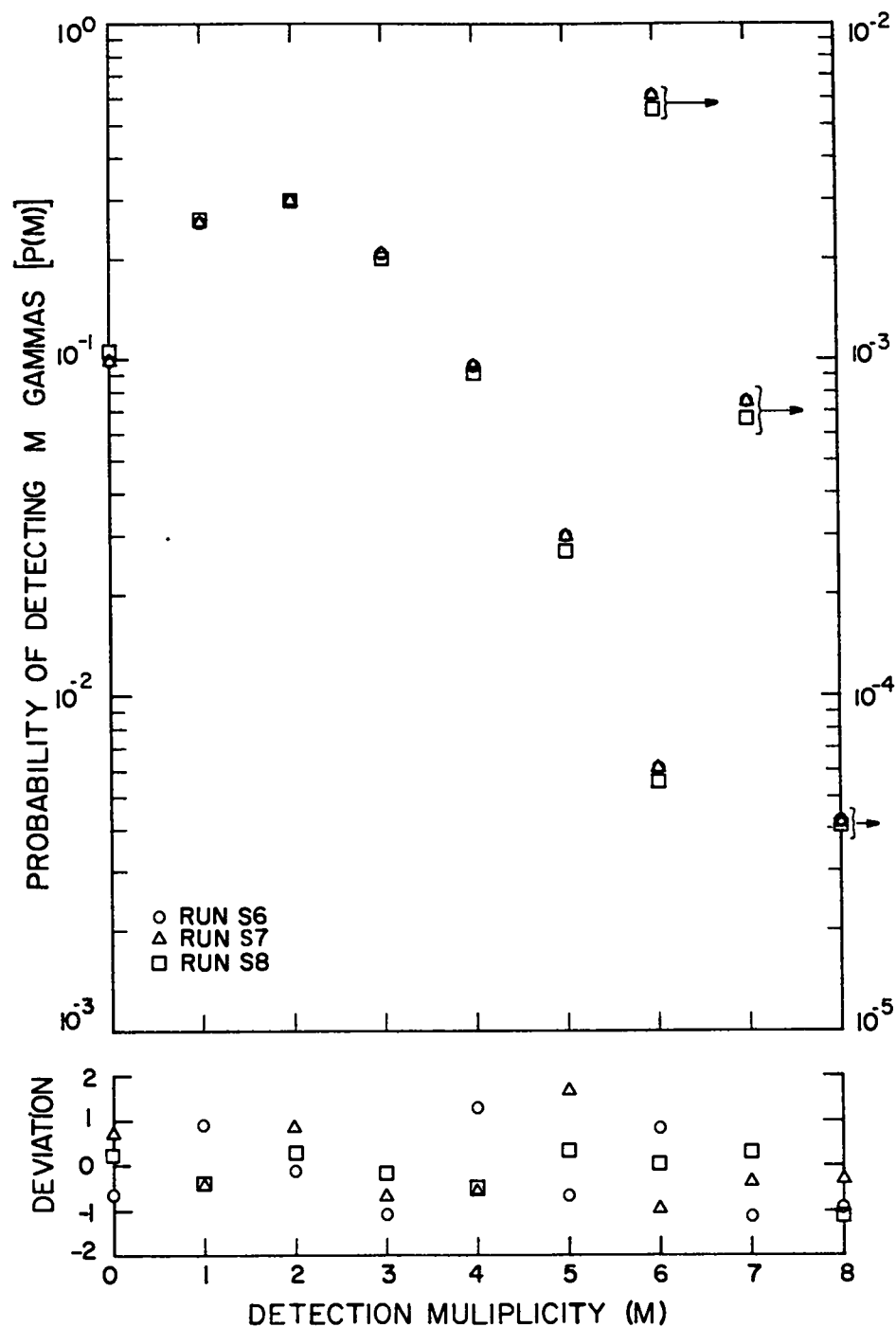


Fig. 11. Experimental results for measurements at 85 keV threshold. The points in the upper graph represent both experiment and model since they differ by too small an amount to show on this scale. The lower graph shows the difference between experiment and model in units of standard deviation.

TABLE 3

RESULTS OBTAINED BY FITTING THE DOUBLE POISSON MODEL^a
TO CORRECTED DATA (from Table 2)

Run Number	Gamma Threshold (keV)	ϵ	C_2	C_3	C_4	$\bar{G}$	σ	χ^2
S2	140	0.0346	0.6700	6.761	9.910	7.80	3.161	1.26
S3	140	0.0351	0.6979	6.862	9.965	7.80	3.135	5.20
S4	140	0.0351	0.7831	6.540	10.099	7.31	3.076	6.07
S6	85	0.0361	0.6874	7.217	10.748	8.32	3.316	7.49
S7	85	0.0361	0.6781	7.186	10.710	8.32	3.321	5.63
S8	85	0.0361	0.7578	7.157	11.038	8.10	3.296	2.21

$$a\pi_G = C_2 \frac{C_3^G}{G!} e^{-C_3} + (1 - C_2) \frac{C_4^G}{G!} e^{-C_4}$$

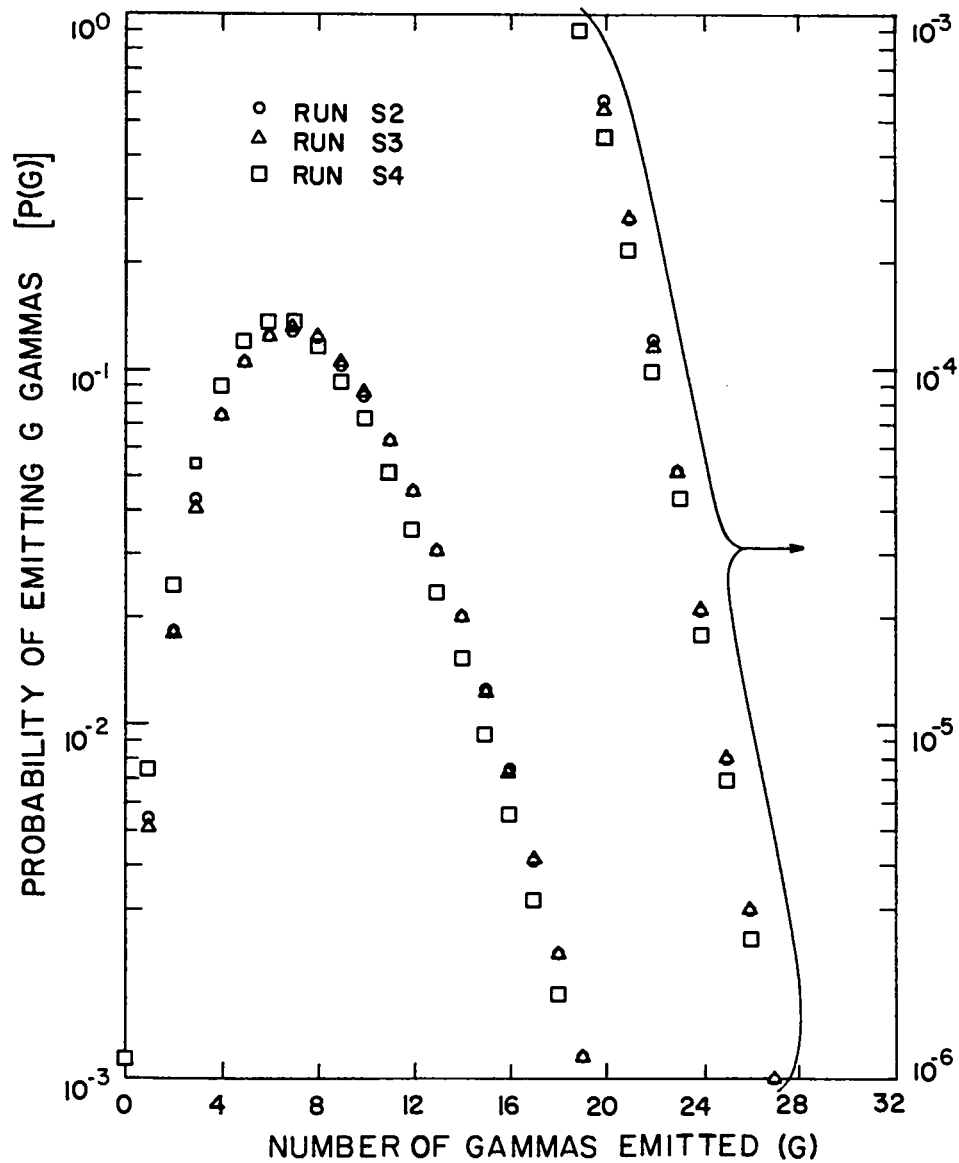


Fig. 12. Emission distribution inferred from the data in Fig. 10(140 keV threshold).

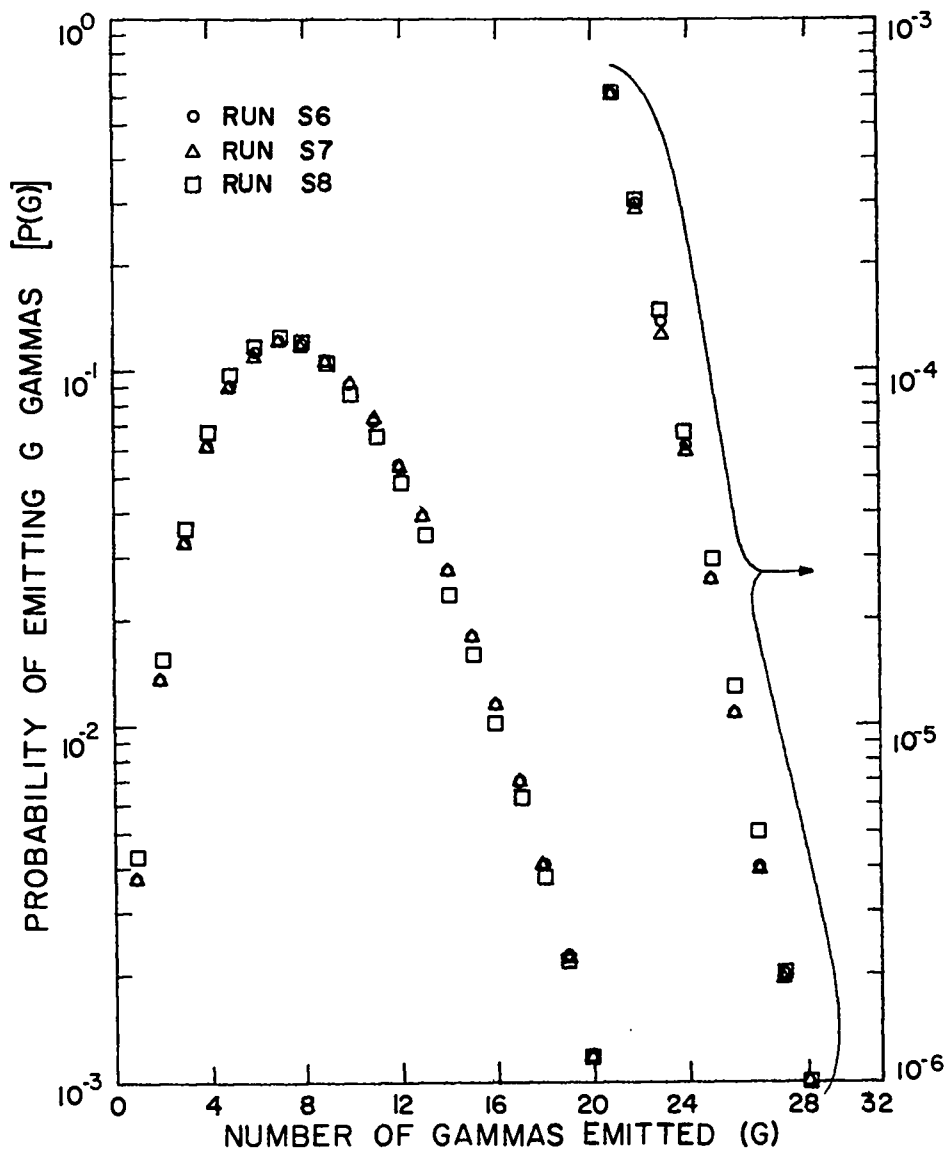


Fig. 13. Emission distribution inferred from the data in Fig. 11(85 keV threshold).

Not only is the double Poisson the best model statistically, it is much easier to handle mathematically. Assume the model

$$\pi(G) = C_2 \frac{(C_3)^G}{G!} e^{-C_3} + (1 - C_2) \frac{(C_4)^G}{G!} e^{-C_4} \quad (15)$$

and combine with $P(G,M)$ to calculate $P(M)$, the probability of detecting an M -fold coincidence.

$$\begin{aligned} P(M) &= \sum_{G=M}^{\infty} \pi(G) P(G,M) \\ &= \sum_{G=M}^{\infty} C_2 \frac{(C_3)^G}{G!} e^{-C_3} P(G,M) + \\ &\quad \sum_{G=M}^{\infty} (1 - C_2) \frac{(C_4)^G}{G!} e^{-C_4} P(G,M) \end{aligned} \quad (16)$$

The two above summations are identical except for constants, so we will evaluate only the first one and insert different constants afterward to obtain the second term.

$$P(M) = \sum_{G=0}^{\infty} C_2 \frac{(C_3)^G}{G!} e^{-C_3} \binom{N}{M} \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} [1 - (N - \ell)\epsilon]^G + \dots \quad (17)$$

We have inserted the Van der Werf expression for $P(G,M)$ and set the lower limit on G to 0 because $P(G,M) = 0$ for $G < M$. Changing the order of summation and rearranging gives:

$$P(M) = C_2 e^{-C_3} \binom{N}{M} (-1)^M \sum_{\ell=0}^M (-1)^{\ell} \binom{M}{\ell} \sum_{G=0}^{\infty} \frac{C_3 [1 - (N - \ell)\epsilon]^G}{G!} \quad (18)$$

The second summation is simply $e^{C_3 [1 - (N - \ell)\epsilon]}$, which, when inserted yields:

$$P(M) = C_2 \binom{N}{M} e^{-C_3 N} (-1)^M \sum_{\ell=0}^M (-1)^\ell \binom{M}{\ell} e^{C_3 \ell} + \dots \quad (19)$$

The summation is just the binomial expansion of $(1 - e^{C_3 \epsilon})^M$; taking into account the term $(-1)^M$, we have

$$P(M) = \binom{N}{M} C_2 e^{-C_3 N \epsilon} (e^{C_3 \epsilon} - 1)^M + (1 - C_2) e^{-C_4 N \epsilon} (e^{C_4 \epsilon} - 1)^M \quad (20)$$

where the second term has been obtained by parallelism with the first. Because C_3 and C_4 are respectively the means of the two distributions

$$\bar{G} = C_2 C_3 + (1 - C_2) C_4. \quad (21)$$

Figure 14 shows the agreement between inferred distributions for run S3 as obtained from two models, double Poisson and skewed pseudo-Gaussian. The agreement is very good between these two models.

At this point, we digress to describe some collateral tests made to verify the efficacy of the computer program. If we set $C_2 = 1$ in the above equation, we can calculate the distribution of detection multiplicities for a hypothetical Poisson emission distribution of mean C_3 . This was done to obtain

$$P(M) \text{ for } 0 \leq M \leq 8.$$

These nine probabilities were then randomized to simulate the scatter that would be expected in a representative run of approximately 10^7 events. We fed these artificial data into the fitting program.

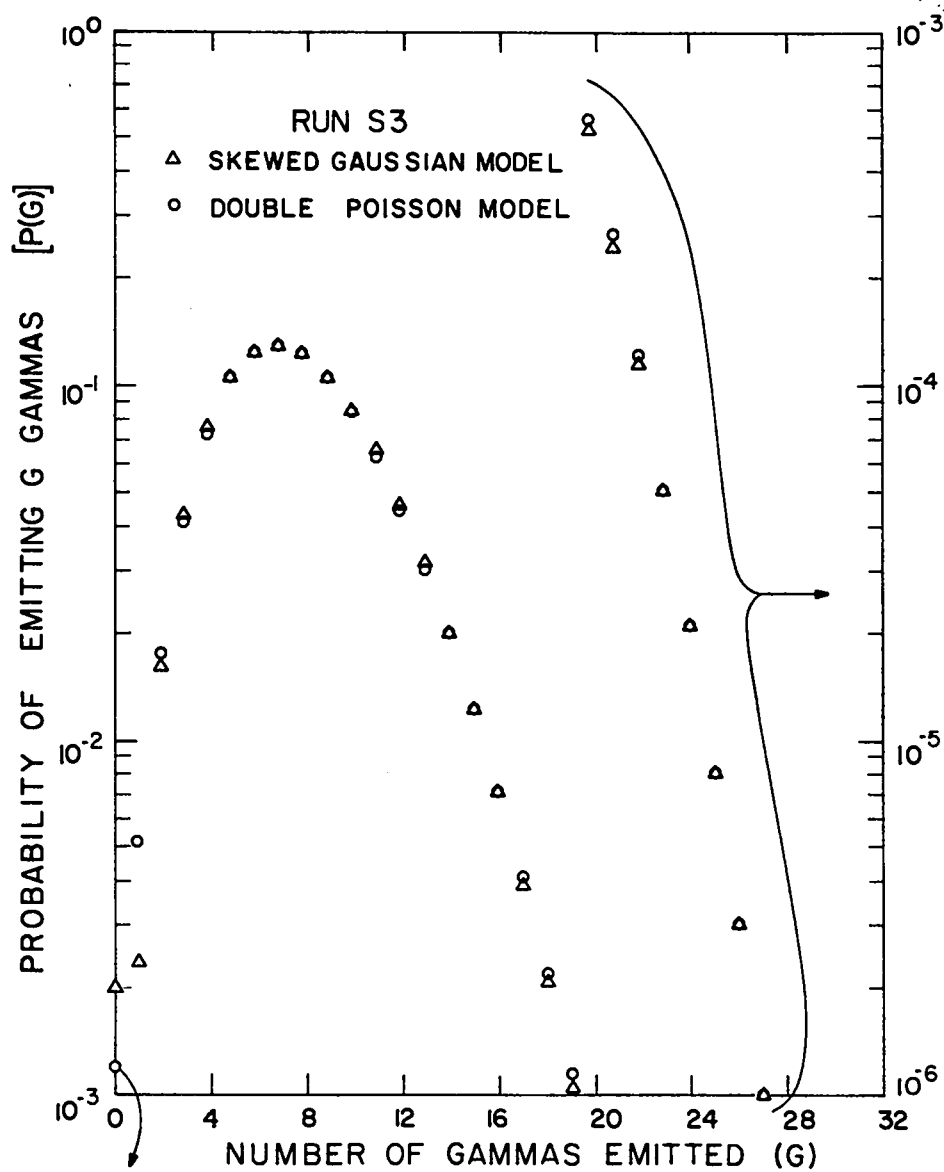


Fig. 14. Comparison of emission distributions for Run S3 as inferred from the double Poisson and skewed pseudo-Gaussian models.

Fitting with the Poisson did, in fact, give us back the original emission distribution within reasonable statistics. More important, when the artificial data was fitted with the skewed pseudo-Gaussian model, the inferred emission distribution was about as close to the input distribution as was the result of the previous fit with the Poisson model.

Energy

It will be recalled from the discussion of the model fitting program that after the computation has converged the program calculates the following:

$\bar{M}'(M)$ the average gamma impact multiplicity as a function
of the detected multiplicity M .

$\bar{G}(M)$ the average emission multiplicity as a function
of the detected multiplicity M .

The data collection program provides $\bar{E}$, which is the sum energy of M triggered detectors divided by M . We then calculate $\bar{E}' = \bar{E} \cdot \frac{M}{\bar{M}'}$ to compensate for the fact that on the average some of the M detectors will have experienced multiple impacts. Table 4 shows these results for the six sets of data. As seen in Figure 15, the average gamma energy (in arbitrary units) decreases for increasing average emission multiplicity. In the cases of runs S4 and S8 in which the fission detector was biased toward fissions of high kinetic energy (that is, low excitation)[11], the effect is more pronounced. The dependence of individual gamma energy on gamma multiplicity is a very interesting phenomenon and we wish to examine the reliability of this observation with particular care.

TABLE 4

DETERMINATION OF AVERAGE GAMMA ENERGY ASSOCIATED WITH MOST PROBABLE EMISSION NUMBER $G(\bar{M})$

Run Number	Threshold keV	M^a	$G(M)^b$	$M'(M)^c$	$E(M)^d$	E'^e	Run Number	Threshold keV	M^a	$G(M)^b$	$M'(M)^c$	$E(M)^d$	E'^e
S2	140	1	6.58	1.136	56.64	49.86	S6	85	1	6.80	1.150	51.93	45.16
		2	7.93	2.282	56.01	49.09			2	8.18	2.312	51.30	44.38
		3	9.32	3.441	55.56	48.44			3	9.63	3.490	50.83	43.69
		4	10.75	4.613	55.27	47.92			4	11.11	4.683	50.49	43.13
		5	12.17	5.796	55.09	47.52			5	12.61	5.893	50.35	42.72
		6	13.58	6.986	54.91	47.16			6	14.08	7.110	50.23	42.39
		7	14.94	8.180	55.62	47.60			7	15.52	8.333	50.27	42.23
		8	16.27	9.373	54.25	46.30			8	16.90	9.956	50.57	42.33
S3	140	1	6.56	1.138	56.41	49.57	S7	85	1	6.80	1.150	51.78	45.02
		2	7.89	2.286	55.79	48.81			2	8.19	2.312	51.18	44.27
		3	9.27	3.446	55.34	48.18			3	9.63	3.490	50.71	43.59
		4	10.68	4.618	55.04	47.67			4	11.12	4.684	50.44	43.82
		5	12.10	5.804	54.93	47.32			5	12.61	5.893	50.25	42.63
		6	13.51	6.996	54.88	47.87			6	14.08	7.109	50.07	42.26
		7	14.01	8.194	54.85	46.86			7	15.51	8.331	50.24	42.21
		8	16.24	9.393	55.40	47.18			8	16.88	9.553	49.26	41.25
S4 Assume ϵ constant	140	1	6.20	1.129	55.41	49.08	S8 Assume ϵ constant	85	1	6.65	1.146	50.49	44.06
		2	7.53	2.268	54.42	47.99			2	8.01	2.303	49.44	42.93
		3	8.93	3.421	53.59	46.89			3	9.47	3.476	48.51	41.87
		4	10.40	4.591	52.94	46.12			4	10.99	4.669	47.91	41.05
		5	11.91	5.781	52.46	45.37			5	12.56	5.883	47.51	40.38
		6	13.41	6.982	52.26	44.90			6	14.11	7.110	47.15	39.79
		7	14.88	8.191	51.79	44.26			7	15.62	8.344	46.32	38.86
		8	16.29	9.401	48.86	41.58			8	17.07	9.580	46.93	39.19

^aDetection multiplicity.^bMost probable emission number for M-fold detection.^cMost probable impact number for M-fold detection.^dAverage energy per detector in the M triggered detector S.^eAverage energy per gamma $E' = \frac{M}{M'(M)} E(M)$

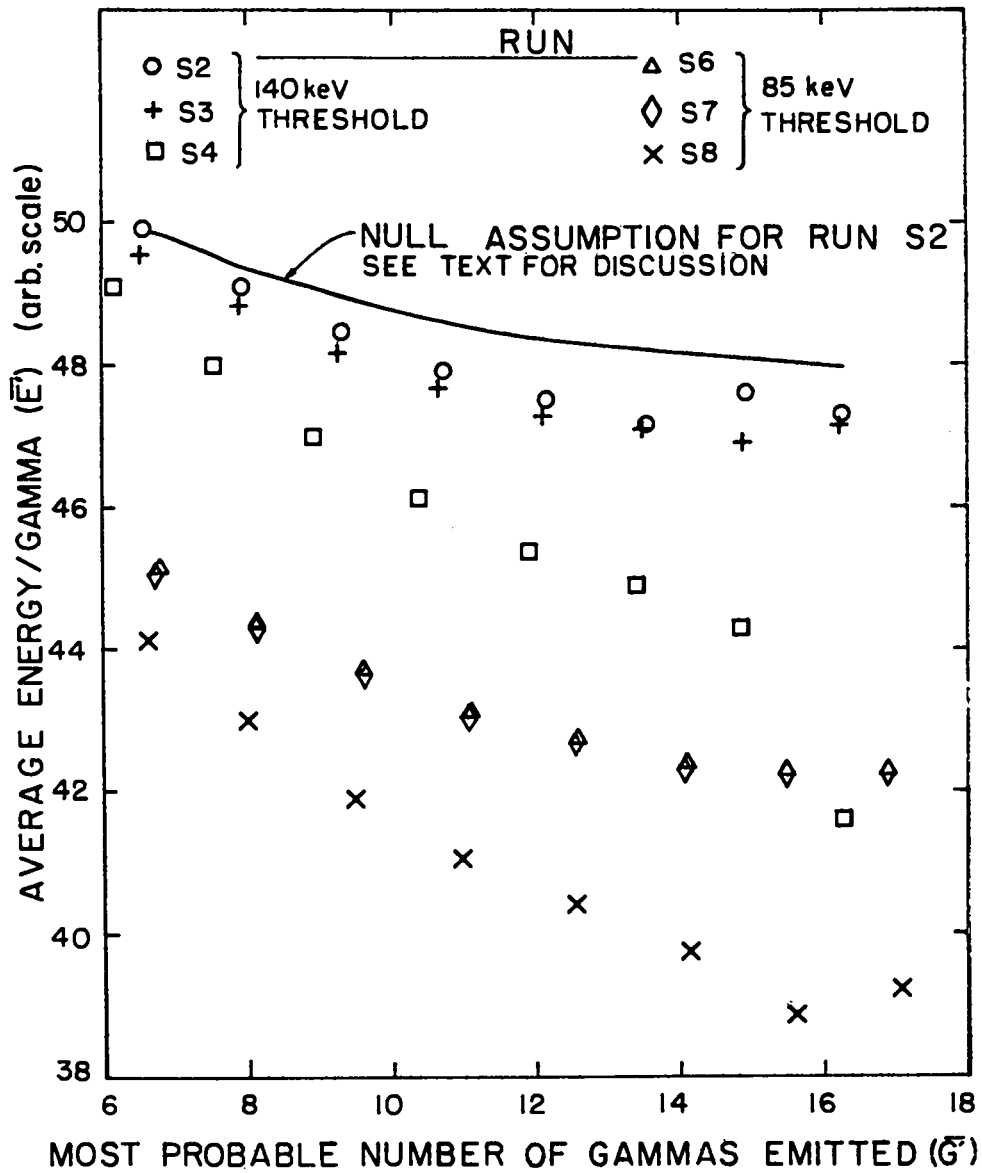


Fig. 15. Average energy of fission gammas as a function of the most probable emission number.

Is it real? For the moment, let us assume that it is not, that the average detected gamma energy is E_0 regardless of the emission multiplicity. In Table 2, we see that a substantial number of events under each multiplicity actually belong in the next lower multiplicity but have been "promoted" by a scattering. This might have an effect on apparent energies somewhat like what has been observed in Figure 15. To be specific, let us examine the data for multiplicity 2 in Run S2 (Table 2). We see that for 2,782,169 events of multiplicity 2, an estimated 54,867 events belonged in multiplicity 1.

At this point, we make two extremely conservative assumptions:

1. The scattered gamma can trigger the spurious count with zero energy input; that is, a single gamma which erroneously triggers a second detector still contributed only E_0 to the energy sum. This obviously ignores the minimum energy deposition required to trigger the second detector.

2. The erroneously counted gamma had an average energy to begin with. (Obviously, a high energy gamma has a better chance of depositing enough energy to trigger the discriminator in the first detector and then scattering with enough remaining energy to trigger a second detector.)

Continuing with our assumptions, we estimate that the data acquisition program accumulated energy as follows for multiplicity 2:

From genuine 2-fold events:

$$2 \times E_0 \times (2,782,169 - 54,867) = 5.4546 \cdot 10^6 E_0$$

From misplaced 1-fold events:

$$1 \times E_0 \times 54,867 = 0.0549 \cdot 10^6 E_0$$

$$5.5095 \cdot 10^6 E_0$$

The data acquisition system assumes all 2-fold events are genuine and divides the energy sum by $2 \times 2,782,169$ to find that $\bar{E}(2) = 0.990 E_0$. The solid line in Figure 15 shows the results of this kind of computation for all multiplicities in Run S2. Despite the severity of the assumptions, the null hypothesis ($\bar{E}_0$ constant) differs significantly from the actual data. We conclude that the effect is real.

VII

SOURCES OF BIAS

There are several sources of spurious signals: for example, cosmic ray showers and fast neutrons. Occasionally, a genuine fission gamma will look like two fission gammas through a Compton scattering that enables it to trigger two separate detectors. Fast neutrons can easily cause scintillation signals in the hydrogenous plastic scintillators through proton recoils, but it is possible to discriminate against them by time of flight. The ^{252}Cf source foil is located at the center of the counter array. This gives a minimum flight distance of about 10 cm to the face of any of the gamma detectors; flight path to the edge of the detector face is around 11 cm. Thus, even though they may penetrate some centimeters into the scintillator before interacting, the gammas traveling at 30 cm/nanosecond will interact, if at all, in less than approximately 2/3 nanosecond. On the other hand, a 4-MeV neutron travels only 2.8 cm/nanosecond. There are few neutrons over 4 MeV. The estimated mean free path for a neutron in the scintillator is about 2 cm. Taking this all into account, we can see that about 4-nanoseconds minimum flight time is required for a 4-MeV neutron and nearly 3 nanoseconds for an 8-MeV neutron. The time separation of the radiations can be seen in Figure 4. The channels are about 48 picoseconds wide.

Aside from the fast neutrons, bias can be expected from the following causes:

1. In the fission detector:
 - a. False triggers from external causes such as cosmic rays or environmental radiation (e.g. ^{40}K).
 - b. False triggers from internal causes such as alphas, betas, neutrons, or nonfission gammas emitted by the ^{252}Cf source.
2. In the eight gamma detectors:
 - a. False counts from external causes such as cosmic rays or environmental radiation (e.g. ^{40}K).
 - b. False counts from nonfission gammas emitted by the fission products in the source or from actual fission gammas that happen not to originate in the trigger fission.
 - c. False multiplicities arising from complex scattering events in the detectors or adjacent structure.
3. Simplification for convenience in analysis:
 - a. Assumption that all gamma detectors have same detection efficiency.
 - b. Assumption that gammas are detected with same efficiency regardless of the multiplicity of the event.

We will take these in turn and discuss the experiments or analysis used to determine their influence on experimental results.

To evaluate 1a above, we removed the ^{252}Cf source and operated the otherwise normal experiment during two overnight runs for a total of about 30 hours. Not a single trigger was observed. This is surprising at first but plausible after some consideration. The scintillator on the fission detector is only 1/4 mm thick. Moreover, the detector is operated at such a high threshold as to reject approximately 6 MeV alphas and approximately 70 MeV fission fragments. Even

if a suitably large energy were transferred to an electron in the scintillator, the mean free path for an electron of such energy is so long that most of the energy would be lost from the scintillator. We can ignore false triggers from this source.

Internal false triggers may occur due to alpha pile-up. In view of the very high threshold, no other source seems plausible. We can estimate an upper limit on the rate of occurrence. The total alpha emission from the ^{252}Cf is $24,000 \text{ sec}^{-1}$, of which one-half, or less, impinge directly on the scintillator. Because of the extremely narrow pulses in the Pilot U scintillator and associated electronics, two pulses must occur within less than approximately 5 nanoseconds for effective pile-up to occur. These numbers lead to an estimate of approximately 0.7 alpha pile-ups per second, or less than one false trigger in a thousand.

A simple measurement was sufficient to evaluate the combined background from items 2a and 2b above. They are both simply counts from events not related to the trigger event. We substituted a signal from an electronic pulse generator for the trigger signal from the fission detector. Everything else was left exactly as for the actual experiment; the ^{252}Cf source remained in its normal position. With the gamma threshold set at 85 keV (the lower of the two used), we counted (in all eight detectors combined) about one event in 10^4 gates. This is a reflection of the stringency of the timing criterion and is small enough to be neglected.

The most important source of error was the scattering effect of real fission gammas (2c above). A Compton scattering in one detector can produce a scattered gamma capable of triggering another detector. Early experiments indicated that was happening with statistically detectable frequency between adjacent detectors. We have previously described the shielding used to minimize this problem. However, it is geometrically impossible to shield nonadjacent detectors from each other without obstructing the detector's "view" of the source.

Another potential source of spurious coincidences is pair production in the lead shielding or other structure. The simultaneous annihilation gammas might also cause a spurious coincidence. Measurements of the gamma spectrum associated with fission show, as would be expected, lead x-rays (at 72.8, 75.0, 84.3, and 87.3 keV) in coincidence with fission; hence these must have been created by genuine fission gammas. Do these x rays constitute a source of false coincidences? It is barely possible that the last two might trigger a detector at the lower of the two thresholds used (85 keV); if the parent gamma happened to trigger another detector, then the multiplicity would be erroneously enhanced ("promoted") by the spurious additional count.

Fortunately, there were simple ancillary experiments to measure the extent to which the californium data is biased by the above processes. We removed the ^{252}Cf source from the fission detector and substituted a source of ^{54}Mn . This isotope decays by electron capture and emits a single gamma of 835 keV (no betas). We increased the gain and decreased the threshold in the fission detector so that

almost any scattering event in the scintillator would trigger an analysis cycle. A second, subsidiary experiment showed that the discriminator was in fact so low that further lowering of the threshold had very small effect on count rate and measured gamma-ray energy. This confirms the above assertion concerning "almost any scattering event."

The experiment was then operated in a normal fashion. Each countable event began with a scattering in the fission detector scintillator. The scattered gamma might then strike one of the eight gamma detectors and be recorded as a multiplicity of 1. Occasionally, the gamma might be re-scattered and trigger one of the other seven detectors, registering a multiplicity of 2. Because the source emits a single gamma, the only way, ignoring background, to observe multiplicity 2 is by means of such a multiple scattering.

If we divide the number of times multiplicity 2 was observed by the number of times multiplicity 1 or 2 was observed, we have a direct measurement of the probability that a scattered gamma from one detector will trigger one of the other seven detectors. Division by 7 then gives the average probability that the scattered gamma will trigger a specified one of the other detectors.

The experiment was then repeated with a ^{60}Co source. Cobalt-60 emits (within approximately 1 picoseconds of each other) two gammas of 1170 and 1330 keV. For the purpose of this experiment, we take them both as 1250 keV. Cobalt-60 also emits betas, but these were suppressed with a cover of polyethylene. With this source, we may expect to observe a substantial number of events with a genuine multiplicity 2, and occasionally there will be an event of multiplicity 3 caused by

a scattering as already described. We divide the number of times multiplicity 3 was observed by the number of times multiplicity 2 or 3 was observed. This gives us the probability that one of the two detectors with genuine events scatters a countable gamma into one of the other six detectors. Dividing this probability by 12 (the number of combinations), we obtain the probability of scattering from one detector to another specified detector.

This is straightforward for the two sources, but the results must be adjusted for energy dependence before they can be used to correct counting data from ^{252}Cf in the experiment proper. The basis for correction is contained in the data from each run. Each operation of the experiment yields an average energy (on an arbitrary scale) for the events counted in the eight gamma detectors. The energies for ^{252}Cf lie conveniently between those for ^{54}Mn and ^{60}Co , so we can get the desired correction probability by simple interpolation. The above described procedure to obtain the corrections probabilities for ^{252}Cf is demonstrated in Table 5.

We restate for the sake of clarity, the "false count probability" estimated for ^{252}Cf is the probability that a genuine fission gamma event in one detector will be associated with a spurious count by scattering into another specified detector.

These probabilities are not small; corrections were made by iteration. Suppose n_3 genuine 3-fold coincidences occurred. There are 15 ways that a scattering event could falsely "promote" a 3-fold to a 4-fold coincidence (scattering from any of the 3 properly triggered detectors to any of the 5 detectors that should not have been

TABLE 5
DETERMINATION OF FALSE COUNT PROBABILITY

	⁵⁴ Mn		⁶⁰ Co		²⁵² Cf	
Threshold (keV)	85	140	85	140	85	140
Counting time (sec)	3,276.8	3,276.8	3,276.8	3,276.8		
Number of singles	470,518 (470,423) ^a	399,492 (399,403)	—	—	—	—
Number of doubles	17,000 (16,951)	7,743 (7,696)	156,625 (156,576)	127,499 (127,452)	—	—
Number of triples	—	—	8,844 (8,834)	4,738 (4,728)	—	—
Fraction of false doubles ^b	0.03478	0.01890	—	—	—	—
Fraction of false triples ^c	—	—	0.0534	0.03577	—	—
False count probability ^d	0.0050	0.0027	0.0044	0.0030	(0.0045) ^e	(0.0029) ^e
Energy (arbitrary units)	24.8	28.1	48.4	52.8	43.5	48.5

^aCount corrected for background in parentheses.

^b(Net doubles)/(Net singles plus net doubles).

^c(Net triples)/(Net doubles plus net triples).

^dIn a specified detector.

^eInferred probability of a false count in a specified detector obtained by interpolation on energy.

triggered.) For example, if the discriminators were set at 140 keV, we would estimate that a certain number, $Q = 15 \times 0.0029 \times n_3$, of the 3-fold multiplicities have been erroneously counted as 4-fold multiplicities. Because we do not know n_3 , the number of genuine 3-fold events, we use n'_3 , the number of observed 3-fold events, as a reasonable estimate of n_3 . We calculate $Q' = 15 \times 0.0029 \times n'_3$ and "retrieve" Q' events from the number of 4-fold events observed. That is, we diminish the 4-fold count by Q' and increase the 3-fold count by the same number. This is done sequentially for all multiplicities except 0 multiplicity, which, of course, remains unchanged. After four iterations, we have a stable set of numbers n_M , which is the best estimate of the true distribution we would have observed if there had been no intercommunication between detectors. The extent of these corrections is shown in Table 2.

For convenience in analysis, it was assumed that all the gamma detectors had the same efficiency (item 3a above). In reality, the efficiencies varied up to 7% from the mean. Van de Werf[25] has also derived a general form of Eq. 5 for coincidences between detectors of unequal efficiency. For the purpose of examining this question, the probabilities of the eight detection multiplicities were calculated for a Poisson emission distribution under two assumptions:

1. Equal efficiencies for all gamma detectors (using Eq. 5)
2. Actual unequal efficiencies (using the general equation)

Consider the differences between the actual efficiencies and the average efficiency as a set of eight perturbations on the average efficiency. To first order in these perturbations, there is no difference

between 1 and 2 above. To second order, the differences in calculated detection multiplicities were all less than 1/3%. This was quite small compared with the corrections made for scattering (Table 2), and no corrections were made for this effect.

The average gamma energy varies with the multiplicity, so the detection efficiency must also vary with multiplicity. In applying Eq. 5, we tacitly assume efficiency does not depend on multiplicity (item 3b on page 60). Measurements with ^{60}Co and ^{54}Mn imply that the efficiency changes very little in the narrow energy region containing all the average energies associated with the various multiplicities. In the worst case, Run S8, the spread is only approximately 15%, and typically is much smaller. Based on the cobalt and manganese source measurements, we estimate a maximum difference in efficiency of 2% between high and low multiplicity events. The effect of this assumed difference can be estimated from Eq. 20. The term with C_4 in it contributes almost all of the probability of high multiplicity events, while the term with C_3 contributes almost all of the probability of low multiplicity events. If we assume the efficiencies in the two terms differ by 2%, we can estimate the effect on C_3 and C_4 , given that the fitting process holds the left hand side constant at the experimental value. This leads to an estimate of 1 1/3% uncertainty for C_3 and 2/3% for C_4 (still considering Run S8). The uncertainties for Run S4 are about the same. For the remaining runs, they are reduced by about one half.

VIII

CONCLUSIONS AND DISCUSSION

The foregoing work leads to the following mathematical models for the multiplicity of gammas emitted in the spontaneous fission of ^{252}Cf :

For gamma threshold at 140 keV:

$$\pi(G) = 0.675 \frac{(6.78)^G e^{-6.78}}{G!} + 0.325 \frac{(9.92)^G e^{-9.92}}{G!} \quad (22)$$

$(G \geq 0)$

For gamma threshold at 85 keV:

$$\pi(G) = 0.682 \frac{(7.20)^G e^{-7.20}}{G!} + 0.318 \frac{(10.71)^G e^{-10.72}}{G!} \quad (23)$$

$(G \geq 0)$

No fundamental physical significance is attributed to these models. They are simply mathematical representations judged most suitable by the χ^2 criterion.

There is a marked decrease in average gamma energy with increasing multiplicity. Nevertheless, the total gamma energy is greater for higher multiplicity; that is, the decrease in individual gamma energy is not large enough to offset the increase in the number of gammas.

An attempt was made to correlate neutron multiplicity with gamma multiplicity. A second 16-channel coincidence register was separately gated to observe discriminator pulses from the gamma detectors in the time window appropriate to the arrival of fast neutrons. The data acquisition program was adapted from the one described earlier.

Preliminary runs were made and the system functioned as described. The problem arises in the discriminator deadtime (nominally 9-1/2 nanoseconds). If a detector detects a gamma, the deadtime in that detector channel extends well into the time window in which neutrons should be detected. As a result, the deadtime corrections were large and could not be estimated with any degree of confidence. The data appeared not worth further analysis.

The correlation of gamma multiplicity with neutron multiplicity could be measured effectively by a fairly simple extension of the present experiment. If the entire gamma detector array in Figure 2 were enclosed in a so-called "barrel counter" (a large cavity neutron counter capable of accepting a full-sized 55-gallon drum), it would be possible with small software changes to count the neutrons after thermalization in the barrel counter walls. This would permit correlated measurements of gammas and neutrons, although it might be anticipated that the results would confirm those of Nifenecker.[11]

An attempt was also made to discover if any gamma energies were preferentially associated with low, moderate or high emission multiplicities. This required the addition of a lithium-drifted germanium detector (GeLi) to the detector array and substantial modification of the data acquisition program. After each fission event in the fission detector, the GeLi detector was queried. If it had detected a gamma, then the multiplicity detected by the eight plastic detectors was determined. Depending on whether the multiplicity was "low" (0-1), "moderate" (2-4), or "high" (5-8), the pulse height determined by the GeLi was stored in one of three 4096-channel spectra.

Here are some tentative results as an indication of the possibilities of this line of research. There were a total of nearly 100 resolvable gamma energies. These included energies agreeing with many of the 24 listed by Rumpold[8] and the 31 listed by Bowman et al.[24] (These two references do not find many lines in common.) The following gamma energies show preferential emission:

With Low Multiplicity (keV)	With High Multiplicity (keV)
145.2	454.5
147.8	470.6
150.0	475.1
165.4	572.7
166.9	578.5
	585.9
290.1	
329.4	
592.1	

The above results come from a single 5-day run. This work requires handling large masses of data and deserves a separate, sustained exploration because of the potential value to the interpretation of nuclear structure.

APPENDIX A

POSSIBLE MODELS FOR GAMMA EMISSION DISTRIBUTION

We have tested some twenty different analytical models. Their inclusion depends to a varying extent on mathematical convenience and plausibility with respect to the fission process. Here are the principal candidates to represent the multiplicity distribution of prompt fission gammas:

1. Poisson distribution

$$\pi(n) = \frac{(\bar{n})^n e^{-\bar{n}}}{n!} \quad (24)$$

This is mathematically the most convenient; its applicability is doubtful because it assumes independent events. The emission of gammas from a single fission cannot be considered independent because they draw on the same store of excitation energy.

2. Double Poisson

$$\pi(n) = A \frac{(\bar{n}_1)^n e^{-\bar{n}_1}}{n!} + (1 - A) \frac{(\bar{n}_2)^n e^{-\bar{n}_2}}{n!} \quad (25)$$

This was suggested by Gordon Hansen as having more flexibility (3 parameters) while still being analytically convenient.

3. Pseudo-Gaussian distribution

$$\pi(n) = C_1 e^{-C_2 (n - \bar{n})^2} \quad (26)$$

This is termed pseudo-Gaussian because it adapts the Gaussian form to a discrete variable, the number of gammas emitted. This is a reasonable possibility because nature sometimes works this way.

4. Skewed pseudo-Gaussian

$$\begin{aligned}\pi(n) &= C_1(n)^{C_3} e^{-C_2(n - \bar{C}_4)^2} & (n \geq 1) \\ \pi(0) &= C_5 e^{-C_2(-C_4)^2}\end{aligned}\tag{27}$$

This distribution provides more parameters to allow more flexibility in fitting. From a practical standpoint, setting $\pi(0)$ free makes it possible to accommodate errors which might arise from even a small fraction of false fission triggers.

5. Pseudo-Cauchy

$$\pi(n) = \frac{C_1}{C_2 + (n - \bar{n})^2}\tag{28}$$

Nature sometimes acts this way, as for example, in the density of nuclear resonance levels.

6. Skewed pseudo-Cauchy

$$\begin{aligned}\pi(n) &= \frac{C_1(n)^{C_3}}{C_2 + (n - \bar{n})^2} & (n \geq 1) \\ \pi(0) &= C_4\end{aligned}\tag{29}$$

This is simply an effort to add more flexibility to the foregoing.

7. Pseudo-lognormal

$$\pi(n) = \frac{C_1 e^{-C_2 (\ln N - C_3)^2}}{n} \quad (30)$$

This is a little farfetched since the lognormal is more often associated with economic or biological distributions. However, it has been applied to the distribution of fragment sizes after a rock crushing operation.

8. A large number of very simple models were tried in the beginning and quickly discarded. They included:

- a. Level distributions of varying lengths.
- b. Ramp distributions of both positive and negative slope and varying lengths.

9. Finally, there was something termed a "broken stick" distribution for want of a better name. In the Hoffman paper,[7] we see that the total gamma energy in a single fission is roughly equal to the energy required to boil off one prompt neutron (approximately 7 MeV).

This leads to the following conjecture: suppose we visualize the total excitation energy of the fission fragments as a stick some 30 or so units (MeV) long. Suppose further that the fission divides the energy (breaks the stick) at some statistically distributed point. The fission fragments then de-excite by emitting prompt neutrons so long as the excitation energy is greater than that necessary to emit a neutron. Thereafter, gammas are emitted. On the average, the energy "left over" for gammas in each fission fragment would amount to about

one half that necessary to emit a neutron, so this would give a total gamma yield equivalent to that required for one neutron. This assumption leads to a distribution in the approximate form of an inverted V (ramp up, ramp down). It was soon found to be a poor candidate and was discarded.

APPENDIX B

VAN DER WERF'S EQUATION

With the permission of the author,[25] we reproduce here (with different symbols) his derivation of the equation used in calculating the number (M) of detectors triggered by a shower of G gammas on an array of N identical counters of efficiency ϵ .

$$P_N(G, M) = \text{Prob. } \{M' \text{ gammas hit somewhere in the } N \text{ detectors}\} \\ \times \text{Prob. } \{M' \text{ gammas are distributed in exactly } M \\ \text{detectors.}\} \quad (31)$$

The first probability is simply the binomial distribution of hits and misses where the probability of a hit is $N\epsilon$.

Prob {that N detectors are hit by M' gammas} =

$$\binom{G}{M'} (N\epsilon)^{M'} (1 - N\epsilon)^{G - M'} \quad (32)$$

The second probability corresponds to that of finding exactly M boxes occupied if M' objects are distributed over N boxes.

Prob {M' gammas are distributed in exactly M detectors} =

$$\binom{N}{M} \frac{M!}{N^{M'}} \mathcal{L}_{M'}^{(M)} \quad (33)$$

$\mathcal{L}_{M'}^{(M)}$ is Stirling's number of the second kind, which gives the number of ways M' objects can be partitioned into M nonempty subsets. In terms of a physical model, it is the number of ways M' gammas can

trigger M interchangeable detectors. To calculate the total number of ways that M' gammas can trigger M detectors out of an N -detector array, the Stirling number is multiplied by:

1. $M!$ to account for the fact that the detectors are identifiable
2. $\binom{N}{M}$ to account for the number of ways that M detectors can be chosen from the N detectors in the array

When this product is divided by the total number of ways that the M' gammas can impinge on N detectors ($N^{M'}$), we have an analytical expression (Eq. 33) for the second probability in Eq. 31. Multiplying the two probabilities, we have:

$$P_N(G, M) = \sum_{M'=M}^G \binom{G}{M'} \epsilon^{M'} (1 - M\epsilon)^{G-M'} \binom{N}{M} M! \mathcal{S}_{M'}^{(M)} \quad (34)$$

Explicitly,

$$\mathcal{S}_{M'}^{(M)} = \frac{1}{M!} \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} \ell^{M'} \quad (35)$$

Because

$$\mathcal{S}_{M'}^{(M)} = 0 \text{ for } M' < M$$

the two expressions can be combined and the summation allowed to run from $M' = 0$ instead of $M' = M$. Rearranging, we have

$$P_N(G, M) = \sum_{\ell=0}^M (-1)^{M-\ell} \binom{M}{\ell} \binom{N}{M} \left[\sum_{M'=0}^G \binom{G}{M'} (\ell\epsilon)^{M'} (1 - N\epsilon)^{G-M'} \right] \quad (36)$$

The term in brackets is the binomial expansion of

$[1 - (N - l)\epsilon]^G$ so

$$P_N(G, M) = \binom{N}{M} \sum_{l=0}^M (-1)^{M-l} \binom{M}{l} [1 - (N - l)\epsilon]^G \quad (37)$$

which is the equation we have used to infer the emission distribution of gammas.

Equation 36 can be rearranged to express the probability that from a shower of G gammas on N detectors, exactly M' of those gammas trigger exactly M detectors.

$$P_N(G, M', M) = \binom{N}{M} \binom{G}{M'} (1 - N\epsilon)^{G-M'} \sum_{l=0}^M (-1)^{M-l} \binom{M}{l} (l\epsilon)^{M'} \quad (38)$$

This equation we have used to calculate the average impact multiplicity $\bar{M}'$.

APPENDIX C

FLOW CHARTS AND PROGRAM LISTINGS FOR THE DATA ACQUISITION PROGRAM

The experiment acquired data by means of two mutually dependent computer programs in the LeCroy 3500. A FORTRAN program, GEN, accepted keyboard commands to set a count time and begin data acquisition. This it did by calling an assembly language subprogram, ACQU1, contained in a listing file named GENMAC. ACQU1 then controlled all data acquisition until expiration of the specified counting time. Control then returned to GEN, which afforded the operator a number of options as to what data would be in the output and in what form.

In this appendix, the flow charts are given first and then the listings for the above programs.

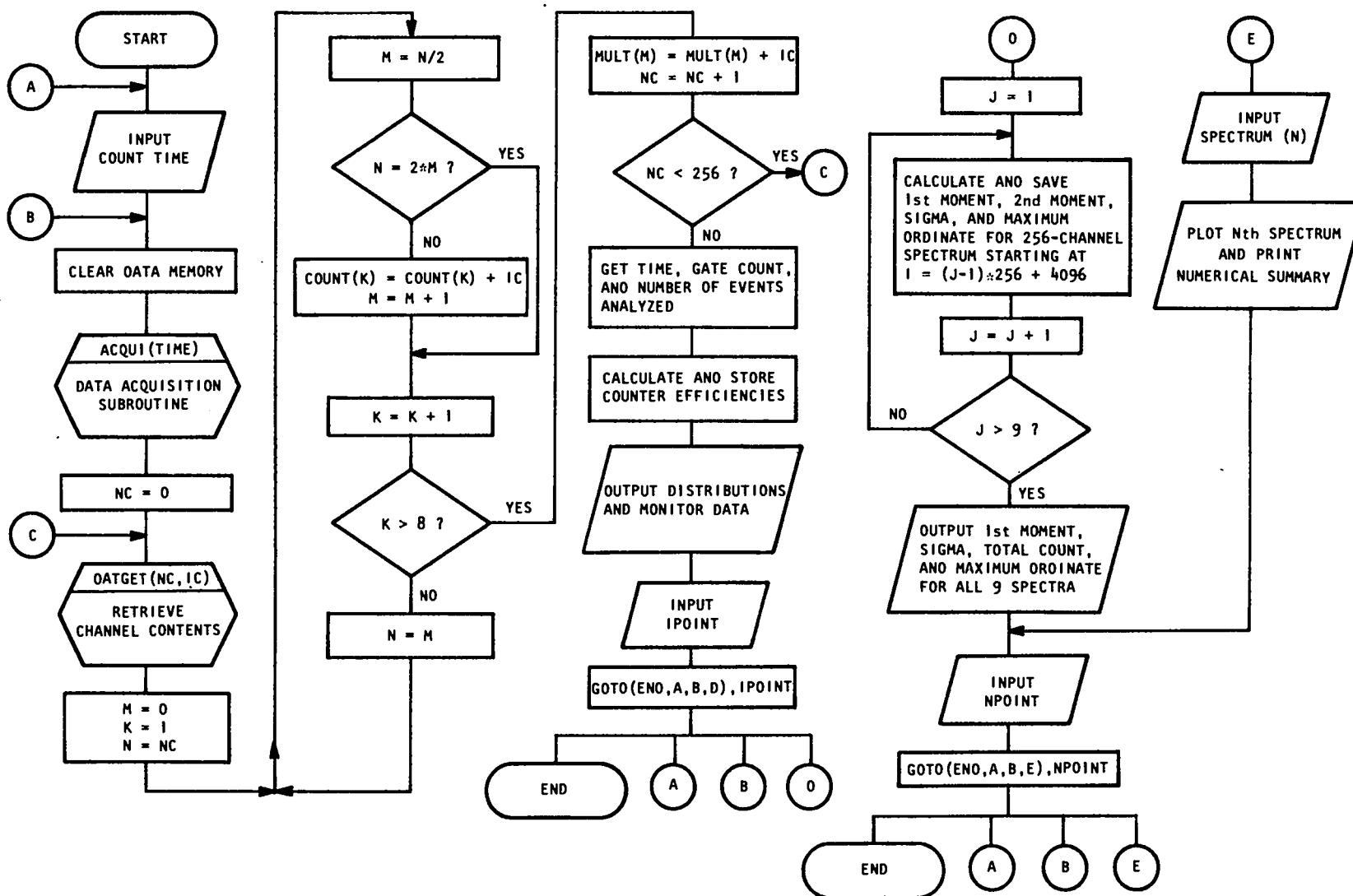


Fig. 16. Flow chart for the main data acquisition program.

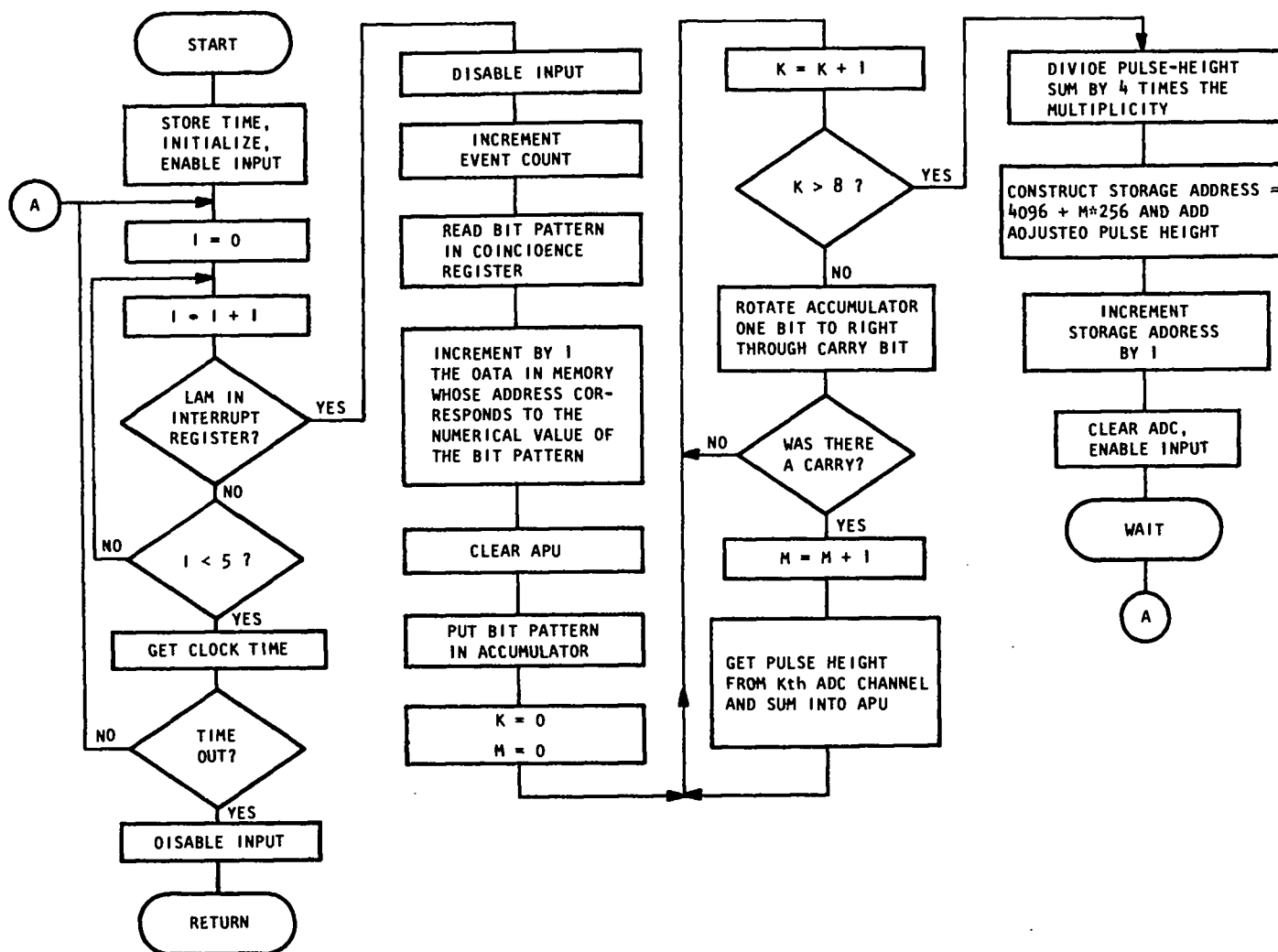


Fig. 17. Flow chart for the assembly language subroutine called by the main program to acquire data.

```

1:      PROGRAM GEN
2:      C
3:      C      THIS IS A FORTRAN PROGRAM TO DRIVE A MACHINE LANGUAGE
4:      C      SUBROUTINE TO ACCUMULATE SIMULTANEOUS DATA ON
5:      C      MULTIPLICITY AND ENERGY.  THE ENERGIES REGISTERED
6:      C      IN THE TRIGGERED CHANNELS ARE SUMMED AND DIVIDED
7:      C      BY 4 TIMES THE MULTIPLICITY TO OBTAIN AN AVERAGE
8:      C      256 CHANNEL SPECTRUM FOR EACH MULTIPLICITY.
9:      C
10:     EXTERNAL ACQU,DATPUT,DATGET,PLINIT,PLOT,SCRDMP,
11:     1  AXIS,WRTSYN,NUMFMT
12:     C
13:     DIMENSION CHAN(10),CHANP(10),FOLD(10),YMAX(10),
14:     1  SGP(10),SENS(10),ESUM(10),CSUM(10),EBAR(10),SIG(10)
15:     C
16:     111  FORMAT(1H )
17:     112  FORMAT(1H0)
18:     113  FORMAT(1H1)
19:     10   FORMAT(' HOW MANY TIME UNITS? UNIT = 65.54 SEC  >>> ')
20:     11   WRITE(1,10)
21:     12   FORMAT(I5)
22:     READ(1,12)NTS
23:     NT=NTS
24:     C
25:     7    DO 20 I=1,300
26:     CALL DATPUT(I,0)
27:     20   CONTINUE
28:     C
29:     CALL DATPUT(8000,0)
30:     CALL DATPUT(0,0)
31:     C
32:     DO 58 K=3000,3100
33:     CALL DATPUT(K,0)
34:     58   CONTINUE
35:     C
36:     DO 59 K=4096,6400
37:     CALL DATPUT(K,0)
38:     59   CONTINUE
39:     C
40:     CALL ACQU(NT)
41:     C
42:     DO 69 M=1,9
43:     CHAN(M)=0
44:     FOLD(M)=0
45:     YMAX(M)=0
46:     ESUM(M)=0
47:     CSUM(M)=0
48:     69   CONTINUE
49:     C
50:     DO 100 J=1,256
51:     JEST=J-1
52:     JADR=JEST

```

```

53:      JM=1
54:      CALL DATGET(JADR,X)
55:      DAT=X
56: C
57:      DO 75 JT=1,8
58:      JQ=JEST/2
59:      JSTAR=2*JQ
60:      IF(JSTAR,EQ,JEST)GO TO 74
61:      JM=JM+1
62:      CHAN(JT)=CHAN(JT)+DAT
63:      74 JEST=JQ
64:      IF(JEST,EQ,0)GO TO 99
65:      75 CONTINUE
66: C
67:      99 FOLD(JM)=FOLD(JM)+DAT
68: C
69:      100 CONTINUE
70: C
71: C      GET INTERNAL COUNT OF ANALYSIS CYCLES
72:      CALL DATGET(8000,X)
73:      TOT1=X
74:      CALL DATPUT(3250,TOT1)
75: C
76: C      GET COUNT OF STROBE GATES
77:      CALL DATGET(8001,X)
78:      TOT2=X
79:      CALL DATPUT(3251,TOT2)
80: C
81: C      GET TIME
82:      CALL DATGET(8002,X)
83:      TIME=X/1000.
84:      CALL DATPUT(3252,X)
85: C
86:      DO 711 JT=1,9
87:      J=JT+260
88:      CALL DATPUT(J,FOLD(JT))
89:      711 CONTINUE
90: C
91:      DO 712 KT=1,8
92:      K=KT+270
93:      CALL DATPUT(K,CHAN(KT))
94:      712 CONTINUE
95: C
96:      PRQ=1
97:      DO 700 J=1,8
98:      CHANP(J)=CHAN(J)/TOT1
99:      Q=1-CHANP(J)
100:      IF(Q,LE,0)GO TO 700
101:      PRQ=PRQ*Q
102:      SENS(J)=ALOG(Q)/(-7.95)
103:      TEMP=SENS(J)*1000000.
104:      STEMP=AIN(TEMP)

```

```

105:      NC=3240+J
106:      CALL DATPUT(NC,STEMP)
107:      700  CONTINUE
108:  C
109:      14  FORMAT(' TOTAL CYCLES = ',F10.0,' TIME = ',
110:      1  F10.1,' SECONDS')
111:      WRITE(2,14)TOT1,TIME
112:      15  FORMAT(' TOTAL INPUT GATES = ',F10.0)
113:      WRITE(2,15)TOT2
114:      WRITE(2,111)
115:      427  FORMAT(' CHANNEL DISTR. OF GAMMAS
116:      1  AND GAMMA SENSITIVITY')
117:      WRITE(2,427)
118:  C
119:      33  FORMAT(1H0,10F9.6)
120:      37  FORMAT(1H0,10F9.0)
121:      WRITE(2,37)(CHAN(I),I=1,8)
122:      WRITE(2,33)(CHANP(I),I=1,8)
123:      WRITE(2,33)(SENS(I),I=1,8)
124:      WRITE(2,111)
125:      S1=0
126:      S2=0
127:      S3=0
128:  C
129:      DO 702 J=1,9
130:      SGP(J)=FOLD(J)/TOT1
131:      S1=S1+CHAN(J)
132:      S2=S2+(J-1)*FOLD(J)
133:      S3=S3+FOLD(J)
134:      702  CONTINUE
135:      612  FORMAT(5X,' GAMMAS:(PRQ =',F9.6,')')
136:      WRITE(2,111)
137:      WRITE(2,111)
138:      WRITE(2,612)PRQ
139:      WRITE(2,37)(FOLD(I),I=1,9)
140:      WRITE(2,33)(SGP(I),I=1,9)
141:      WRITE(2,111)
142:      WRITE(2,111)
143:      271  FORMAT(' TOTAL GAMMAS FROM DETECTORS = ',F10.0)
144:      WRITE(2,271)S1
145:  C
146:      272  FORMAT(' TOTAL GAMMAS FROM DAT.MATR. = ',F10.0)
147:      WRITE(2,272)S2
148:  C
149:      273  FORMAT(' TOTAL CYCLES FROM DAT.MATR. = ',F10.0)
150:      WRITE(2,273)S3
151:  C
152:  C
153:      DO 703 M=1,9
154:      MP=M+3219
155:      CALL DATPUT(MP,FOLD(M))
156:      703  CONTINUE

```

```

157: C
158:      WRITE(1,112)
159: C
160:      22  FORMAT(' WHAT NEXT? 1=QUIT,2=CHGE TIME,3=REPEAT,
161:      1  RTN=LOOK AT ENERGIES' )
162:
163:      WRITE(1,22)
164:      READ(1,12)NEXT
165:      IF(NEXT.EQ.1)GO TO 7000
166:      IF(NEXT.EQ.2)GO TO 11
167:      IF(NEXT.EQ.3)GO TO 7
168: C
169: C      FOLLOWING IS THE ROUTINE FOR ENERGY
170: C      ANALYSIS AND PLOTTING
171: C
172:      300  DO 302 J=1,9
173:           L1=4096+(J-1)*256
174:           L2=L1+255
175:           SUM1=0
176:           SUM2=0
177:           SUM3=0
178:           INDEX=0
179: C
180:           DO 301 K=L1,L2
181:           CALL DATGET(K,X)
182:           SUM1=SUM1+X
183:           FIN=FLOAT(INDEX)
184:           SUM2=SUM2+X*FIN
185:           SUM3=SUM3+X*(FIN**2)
186:           YMAX(J)=AMAX1(X,YMAX(J))
187:           INDEX=INDEX+1
188:      301  CONTINUE
189: C
190:           CSUM(J)=SUM1
191:           ESUM(J)=SUM2
192:           IF(SUM1.LE.0)GO TO 901
193:           EBAR(J)=SUM2/SUM1
194:           TEMP=SUM3/SUM1
195:           TEMP1=(TEMP-(EBAR(J)**2))
196:           IF(TEMP1.LE.0)GO TO 911
197:           SIG(J)=SQRT(TEMP1)
198:      302  CONTINUE
199: C
200:           GO TO 399
201: C
202:      900  FORMAT(' CSUM(' ,I2,' ) = 0' )
203:      901  WRITE(2,900)J
204:      902  FORMAT(' SUM1=',F10.1,' SUM2=',F10.1,' SUM3=',F10.1)
205:           WRITE(2,902)SUM1,SUM2,SUM3
206:           CSUM(J)=9999.
207:           GO TO 302
208: C

```

```

209: 910  FORMAT(' SQRT ARG =',F10.4)
210: 911  WRITE(2,910)TEMP1
211: 912  FORMAT(' J= ',I2,' SUM1=',F10.1,' SUM3=',F10.1,
212: 1    ' EBAR =',F8.4)
213:      WRITE(2,912)J,SUM1,SUM3,EBAR(J)
214:      SIG(J)=88.88
215:      GO TO 302
216: 380  FORMAT(2X,'BLOCK',8X,'COUNTS',10X,'CHAX',
217: 1    10X,'EBAR',8X,'SIGMA')
218: 399  WRITE(2,380)
219:      WRITE(2,111)
220: 381  FORMAT(2X,I3,7X,F10.1,6X,F8.1,5X,F9.5,5X,F9.5)
221:      DO 303 I=1,9
222:          K=I-1
223:          WRITE(2,381)K,CSUM(I),YMAX(I),EBAR(I),SIG(I)
224: 303  CONTINUE
225:      WRITE(2,111)
226: C
227: 382  FORMAT(' CHOICES: 1=QUIT, 2=CHANGE TIME, 3=REPEAT
228: 1    RTN=LOOK AT SPECTRA')
229:      WRITE(1,382)
230: 383  FORMAT(I2)
231:      READ(1,383)KPOINT
232:      IF(KPOINT.EQ.1)GO TO 7000
233:      IF(KPOINT.EQ.2)GO TO 11
234:      IF(KPOINT.EQ.3)GO TO 7
235: C
236: 384  FORMAT(' WHICH SPECTRUM??')
237: 304  WRITE(1,384)
238:      READ(1,383)MPOINT
239:      JPPT=MPOINT+1
240:      L1=4096+MPOINT*256
241:      L2=L1+255
242:      XQ=FLOAT(MPOINT)
243: C
244:      IFLAG=0
245: 339  CALL PLINIT
246: C
247:      DO 328 I=1,225
248:          IX=2*I+61
249:          LP=L1+I-1
250:          CALL DATGET(LP,X)
251:          TEMP=200*X/YMAX(JPPT)
252:          IY=INT(TEMP)+55
253:          CALL PLOT(IX,IY,3)
254: 328  CONTINUE
255:      XL=450./82.
256:      YL=200./51.
257:      XM=0.
258:      YM=0.
259:      XDL=82.
260:      YDL=YMAX(JPPT)/YL

```

```

261:      CALL AXIS(61,55,'CHANNEL NUMBER',14,XL,0.,XM,XDL)
262:      CALL AXIS(61,55,'COUNTS',-6,YL,90.,YM,YDL)
263:      CALL PLOT(511,55,3)
264:      CALL PLOT(511,255,2)
265:      CALL PLOT(61,255,2)
266:      CALL WRTSYM(300,150,'BLOCK',5,0.)
267:      CALL NUMFMT(360.,150,XQ,'(F4.0)',0.)
268:  C
269:      IF(IFLAG.EQ.1)GO TO 387
270:      CALL WRTSYM(300,180,'PLOT? 1=YES,RTN=NO',18,0.)
271:      READ(1,383)JPT
272:      IFLAG=1
273:      IF(JPT.EQ.1)GO TO 339
274:      GO TO 386
275: 387  CALL SCRDMF
276: 401  FORMAT(' COUNT TIME =',F11.4)
277: 402  FORMAT(' MAX ORDINATE =',F12.4)
278: 403  FORMAT(' EBAR   =',6X,F9.4)
279: 404  FORMAT(' SIG    =',6X,F9.4)
280: 405  FORMAT(' TOTAL COUNT =',F12.1)
281:      WRITE(2,401)TIME
282:      WRITE(2,402)YMAX(JPPT)
283:      WRITE(2,403)EBAR(JPPT)
284:      WRITE(2,404)SIG(JPPT)
285:      WRITE(2,405)CSUM(JPPT)
286:      WRITE(2,113)
287:  C
288: 386  FORMAT(' NEXT? 1=QUIT,2=CHANGE TIME,
289: 1    3=REPEAT, RTN=NEW BLOCK' )
290:      WRITE(1,386)
291:      READ(1,383)KPT
292:      IF(KPT.EQ.1)GO TO 7000
293:      IF(KPT.EQ.2)GO TO 11
294:      IF(KPT.EQ.3)GO TO 7
295:      GO TO 304
296:  C
297: 7000  END

```

```

1: ;          PROGRAM GENMAC.MAC
2: ;
3: ;THIS DATA ACQUISITION ROUTINE IS BASED ON SIMMAC.MAC.
4: ;IT USES AN ADDITIONAL 8K OF DATA MEMORY TO STORE
5: ;SPECTRAL INFORMATION CORRELATED WITH THE MULTIPLICITY
6: ;OF THE DETECTION.
7: ;
8: ;          SLOT ASSIGNMENTS
9: ;
10: ;SLOT 0          2341A COINCIDENCE REGISTER
11: ;SLOT 1          MODEL 3610 (TIMING SCALER)
12: ;                A=0 >> GATE COUNTER
13: ;                A=1 >> TIMER(POSITIVE INPUT)
14: ;SLOT 2          MODEL 41 OUTPUT REGISTER
15: ;SLOT 3          IRO26 INPUT REGISTER
16: ;SLOT 5          2249A ADC(1ST 8 CHANNELS)
17: ;SLOT 7          CLOCK(NOT ADDRESSED)
18: ;
19: ;THE PROGRAM EXTENDS SIMMAC.MAC TO SUM AND AVERAGE THE
20: ;GAMMA ENERGIES FROM A GIVEN EVENT. THESE ARE REDUCED
21: ;TO 256 CHANNEL SPECTRA AND STORED IN DIFFERENT MEMORY
22: ;SEGMENTS ACCORDING TO THE INDEX OF THE MULTIPLICITY.
23: ;UNLIKE SIMMAC.MAC,THIS PROGRAM REQUIRES THAT EACH
24: ;EVENT BE ANALYZED IMMEDIATELY.
25: ;
26: ;IT USES THE APU ON THE DISPLAY CONTROL BOARD
27: ;TO SUM THE ENERGIES AND PERFORM THE DIVISION.
28: ;
29: ;          ;CAMAC I/O ADDRESSES
30: ;          LO      EQU      0D0H      ;LO BYTE(IN/OUT)
31: ;          MID     EQU      0D1H      ;MID BYTE(IN/OUT)
32: ;          HI      EQU      0D2H      ;HI BYTE(IN/OUT)
33: ;          XQL     EQU      0D4H      ;XQL REGISTER(IN)
34: ;          FA      EQU      0D4H      ;FA REGISTER(OUT)
35: ;          NC      EQU      0D5H      ;NCREGISTER(OUT)
36: ;          CYCLE   EQU      0D7H      ;CYCLE COMMAND(OUT)
37: ;
38: ;          ;DATA MEMORY ADDRESSES
39: ;          ADDRLO  EQU      0          ;LOW ADDRESS BYTE
40: ;          ADDRHI  EQU      1          ;HI ADDRESS BYTE
41: ;          WRLO    EQU      2          ;LO DATA BYTE(OUT)
42: ;          WRMID   EQU      3          ;MID DATA BYTE(OUT)
43: ;          WRHI    EQU      4          ;HI DATA BYTE (OUT)
44: ;          RDLO    EQU      5          ;LO DATA BYTE(IN)
45: ;          RDMID   EQU      6          ;MID DATA BYTE(IN)
46: ;          RDHI    EQU      7          ;HI DATA BYTE(IN)
47: ;
48: ;
49: ;          ;DEFINE ACCESS TO APU
50: ;
51: ;          BASE    EQU      0F0H
52: ;          CMDBR   EQU      0F1H

```

```

53:          SDIV    EQU    06FH
54:          SADD    EQU    06CH    ;ADD COMMAND
55:          STATR   EQU    0F1H    ;APU STATUS
56:          ;
57:          ;DEFINE STORAGE
58:          ;
59:  TIME:    DS      1
60:  BYTE:    DS      1
61:  ;
62:          ;
63:  LOOK     MACRO
64:          MVI      A,6            ;LOOK FOR LAM
65:          OUT      NC            ;IN IRO
66:          MVI      A,8H          ;N=3,A=0,F=8
67:          OUT      FA
68:          OUT      CYCLE
69:          IN       XQL
70:          ANI      40H           ;IS THERE A LAM?
71:          JNZ      READ          ;YES
72:          ENDM
73:          ;
74:  FILE     MACRO                ;ROUTINE TO TAKE DATA FROM A
75:          IN       LO            ;SPECIFIED CAMAC ADDRESS AND
76:          OUT      WRLO          ;WRITE TO A SPECIFIED DATA
77:          IN       MID           ;ADDRESS.(IN 1ST MEMORY)
78:          OUT      WRMID
79:          IN       HI
80:          OUT      WRHI
81:          ENDM
82:          ;
83:  INK      MACRO                ;ROUTINE TO INCREMENT A SPECIFIED
84:          IN       RDLO          ;CHANNEL IN 1ST DATA MEMORY.
85:          ADI      1
86:          OUT      WRLO
87:          IN       RDMID
88:          ACI      0
89:          OUT      WRMID
90:          IN       RDHI
91:          ACI      0
92:          OUT      WRHI
93:          ENDM
94:          ;
95:          ;
96:          ENTRY   ACQU
97:          ;
98:  ACQU:    MOV     A,M
99:          STA     TIME
100:         ;
101:         ;ROUTINE TO TURN EXPERIMENT ON
102:         MVI     A,0FH          ;WRITE 1'S INTO LO HALF
103:         OUT     LO             ;OF OUTPUT REGISTER
104:         MVI     A,4            ;IN SLOT #2

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```

105:      OUT      NC          ;N=2,A=0,F=16
106:      MVI      A,10H
107:      OUT      FA
108:      OUT      CYCLE
109:      ;
110:      LXI      B,0A09H      ;CLEAR ADC
111:      CALL     CYC          ;N=5,A=0,F=9
112:      ;
113:      LXI      B,0A18H      ;DISABLE ADC LAM
114:      CALL     CYC          ;N=5,A=0,F=24
115:      ;
116:      LXI      B,0209H      ;CLEAR GATE COUNTER
117:      CALL     CYC          ;N=1,A=0,F=9
118:      ;
119:      LXI      B,0229H      ;CLEAR TIME
120:      CALL     CYC          ;N=1,A=1,F=9
121:      ;
122:      LXI      B,0218H      ;DISABLE GATE LAM
123:      CALL     CYC          ;N=1,A=0,F=24
124:      ;
125:      ;
126:      MVI      A,0F0H      ;WRITE 1'S INTO HI HALF
127:      OUT      LO          ;OF OUTPUT REGISTER
128:      MVI      A,4         ;IN SLOT #2
129:      OUT      NC          ;N=2,A=0,F=16
130:      MVI      A,10H
131:      OUT      FA
132:      OUT      CYCLE
133:      MVI      A,6         ;CLEAR IRO & IRO LAM
134:      OUT      NC          ;N=3,A=0,F=10
135:      MVI      A,0AH
136:      OUT      FA
137:      OUT      CYCLE
138:      MVI      A,0         ;CLEAR COINC REG
139:      OUT      NC          ;N=0,A=0,F=9
140:      MVI      A,9
141:      OUT      FA
142:      OUT      CYCLE
143:      MVI      A,6         ;ENABLE IRO LAM
144:      OUT      NC          ;N=3,A=0,F=26
145:      MVI      A,01AH
146:      OUT      FA
147:      OUT      CYCLE
148:      JMP      WAIT
149:
150:      ;
151:      ;ROUTINE TO TURN EXPERIMENT OFF
152:  OFF:  MVI      A,0FH      ;WRITE 1'S INTO LO HALF
153:      OUT      LO          ;OF OUTPUT REGISTER
154:      MVI      A,4         ;IN SLOT #2
155:      OUT      NC          ;N=2,A=0,F=16
156:      MVI      A,10H

```

```

157:      OUT      FA
158:      OUT      CYCLE
159:      LXI      B,0202H      ;READ GATE COUNTER
160:      CALL     CYC          ;N=1,A=0,F=2
161:      ;
162:      MVI      A,1FH
163:      OUT      ADDRHI
164:      MVI      A,41H      ;PUTTING GATE COUNT
165:      OUT      ADDRLO      ;IN ABS CHAN 8001
166:      FILE
167:      LXI      B,0222H      ;READ TIME
168:      CALL     CYC          ;N=1,A=1,F=2
169:      ;
170:      MVI      A,42H      ;PUTTING TIME IN
171:      OUT      ADDRLO      ;ABS CHAN 8002
172:      FILE
173:      RET
174:      ;
175:      ;ROUTINE TO READ EXPERIMENTAL EVENT
176: READ:  MVI      A,0FH      ;WRITE 1'S INTO LO
177:      OUT      LO          ;OF OUTPUT REGISTER
178:      MVI      A,4         ;IN SLOT #2
179:      OUT      NC          ;N=2,A=0,F=16
180:      MVI      A,10H
181:      OUT      FA
182:      OUT      CYCLE
183:      MVI      A,1FH
184:      OUT      ADDRHI
185:      MVI      A,40H      ;TALLY THE NUMBER
186:      OUT      ADDRLO      ;OF CYCLES IN 8000
187:      INK
188:      MVI      A,0         ;READ & CLEAR COINC REG
189:      OUT      NC          ;N=0,A=0,F=2
190:      MVI      A,2
191:      OUT      FA
192:      OUT      CYCLE
193:      IN      MID          ;GET COINC MID BYTE
194:      STA      BYTE        ;SAVE PATTERN FOR SHUFFLING
195:      OUT      ADDRLO
196:      MVI      A,0         ;ZERO HIGH ADDRESS
197:      OUT      ADDRHI      ;FOR PATTERN STORAGE
198:      INK
199:      ;
200:      ;ENERGY ANALYSIS
201:      ;
202:      LDA      BYTE        ;GET BYTE PATTERN
203:      ADI      0           ;IS THERE DATA
204:      JZ       CLEAR       ;NO,GO BACK AND WAIT
205:      ;
206:      ;           YES, ANALYZE LIKE CRAZY
207:      MVI      B,0         ;CLEAR MULT. COUNTER
208:      MVI      C,0         ;SET SUBADDRESS INDICATOR

```

209:		MVI	A,0	
210:		OUT	BASE	
211:		OUT	BASE	;CLEAR APU TO ADD ENERGIES
212:		MVI	D,8	;SET LOOP COUNTER
213:				
214:	LOOP:	LDA	BYTE	;GET BYTE PATTERN
215:		RAR		
216:		STA	BYTE	
217:		JNC	AGAIN	
218:				
219:		INR	B	;YES, TALLY ONE GAMMA
220:		MOV	A,C	;GET SUBADDRESS OF GAMMA
221:		RRC		;PUT SUBADDRESS INBITS 5-7
222:		RRC		
223:		RRC		
224:		ANI	0E0H	
225:		OUT	FA	;F=0, SO OUT TO FA
226:		MVI	A,0AH	;N=5
227:		OUT	NC	
228:		OUT	CYCLE	;READ INDICATED MCA CHANNEL
229:		IN	LO	
230:		OUT	BASE	;PUT ADC DATA IN APU
231:		IN	MID	;FOR ADDING
232:		OUT	BASE	
233:		MVI	A,SADD	;GET ADD COMMAND
234:		OUT	CMDR	;AND SEND TO APU
235:				;AND LOOK FOR NEXT GAMMA
236:	AGAIN:	INR	C	;ADVANCE SUBADDRESS
237:		DCR	D	;CHECK OFF ONE ADC
238:		JZ	CALC	;IF DONE,CALC AVERAGE
239:		JMP	LOOP	;NOT DONE,LOOK FOR NEXT GAMMA
240:				
241:	CALC:	MOV	A,B	;GET NUMBER OF GAMMAS
242:		RLC		
243:		RLC		
244:		ANI	03CH	;MULTIPLY BY 4
245:		OUT	BASE	;LOW BYTE OF DIVISOR
246:		MVI	A,0	
247:		OUT	BASE	;HI BYTE OF DIVISOR
248:		MVI	A,SDIV	;GET DIVIDE COMMAND
249:		OUT	CMDR	;AND SEND TO APU
250:				;THIS GIVES AVERAGE ENERGY/4
251:				;GIVING A 256 CHANNEL SPECTRUM
252:				;NOW SET UP ADDRESS FOR STORAGE
253:	NTDN:	IN	STATR	;WAIT FOR APU
254:		ADI	0	
255:		JM	NTDN	
256:		IN	BASE	
257:		IN	BASE	;AND USE LOW BYTE OF QUOTIENT
258:				;FOR LOW ADDRESS
259:		OUT	ADDRLO	
260:		MVI	A,010H	

```

261:      OUT      ADDRHI      ;PUT ALL EVENTS IN ZEROth BASKET
262:      INK
263:      MVI      A,010H
264:      ADD      B            ;ADD MULTIPLICITY IN HI BYTE
265:      OUT      ADDRHI      ;AND ALSO STORE IN B-TH BASKET
266:      INK                ;@4096 + B * 256.
267:      ;
268:      ;NOW CLEAR 2249A
269:      ;
270:  CLEAR:  MVI      A,0AH      ;N=5,A=0,F=9
271:      OUT      NC
272:      MVI      A,9
273:      OUT      FA
274:      OUT      CYCLE
275:      ;NOW THROUGH WITH ENERGY
276:      ;AND BACK TO CONTINUE
277:      ;WITH CLEAN UP FOR
278:      ;NEXT EVENT
279:      ;
280:      MVI      A,6            ;CLEAR IRO & IRO LAM
281:      OUT      NC            ;N=3,A=0,F=10
282:      MVI      A,0AH
283:      OUT      FA
284:      OUT      CYCLE
285:      MVI      A,0F0H        ;WRITE 1'S INTO HI HALF
286:      OUT      LO            ;OF OUTPUT REGISTER
287:      MVI      A,4            ;IN SLOT #2
288:      OUT      NC            ;N=2,A=0,F=16
289:      MVI      A,10H
290:      OUT      FA
291:      OUT      CYCLE
292:      MVI      A,6            ;ENABLE IRO LAM
293:      OUT      NC            ;N=3,A=0,F=26
294:      MVI      A,01AH
295:      OUT      FA
296:      OUT      CYCLE
297:      JMP      WAIT
298:      ;
299:  CYC:    MOV      A,B        ;DOES CAMAC CYCLE FOR
300:      OUT      NC            ;CONTROL WORD IN BC
301:      MOV      A,C
302:      OUT      FA
303:      OUT      CYCLE
304:      RET
305:      ;
306:  WAIT:   LDA      TIME      ;PUT COUNT TIME IN D
307:      MOV      D,A
308:      LOOK
309:      LOOK
310:      LOOK
311:      LOOK
312:      LOOK

```

313:	MVI	A,02H	;CHECK TIME
314:	OUT	NC	;READ CLOCK
315:	MVI	A,020H	;N=1,A=1,F=0
316:	OUT	FA	
317:	OUT	CYCLE	;AND COMPARE
318:	IN	HI	
319:	CMP	D	;WITH TIME IN D
320:	JC	WAIT	;NOT OUT YET
321:	JMP	OFF	;TIME IS OUT
322:	END		

ACKNOWLEDGMENTS

I wish to acknowledge with gratitude the hospitality, interest, and assistance of the many Los Alamos people who contributed to this work. I thank William G. Davey, former Q (Energy) Division Leader; and Carl Henry and E. J. Dowdy, former and present Group Leaders of Group Q-2 (Advanced Nuclear Technology) for their personal and official support. Special thanks go to Gordon Hansen for hours of counsel and encouragement, and for performing many supporting calculations. John Pratt contributed greatly in the procurement of equipment and in putting it in operation. Everyone in Q-2 contributed with a helping hand and a smile for a visitor.

I appreciate the courtesy of W. K. McFarlane and the Experimental Elementary Particle Physics Group of Temple University who loaned the high-performance photomultipliers and preamplifiers used in the main detector array.

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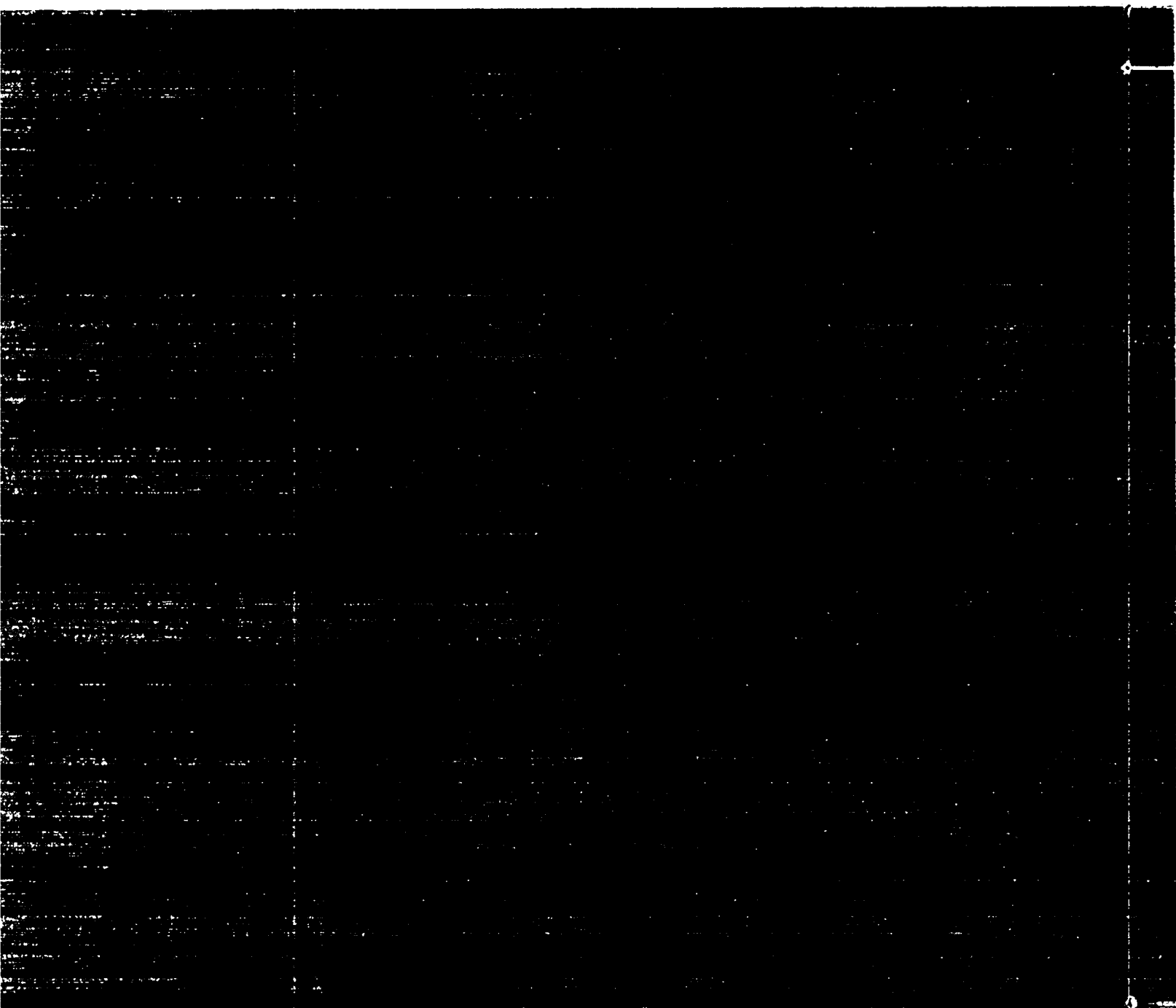
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