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REMARKS ON THE FISSION-CAPTURE RATIO OF 25

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The capture-fission ratio  $\propto = C_{capture}/C$  fission of 25 can be messured at thermal emergies but this measurement cannot be extended to omorgies occuring in the fission spectrum. Therefore theoretical conriderations to estimate the energy dependence of & are relevant.

The ratio & is given by

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$$o_{t} = \frac{\frac{T_{r}}{r}}{T_{o}}$$

where  $\Gamma_r$  and  $\Gamma_r$  are the partial widths in respect to radiation and to fission respectively. We are interested in the ratio K of the values of at low energies to the values around 1 MeV:

 $\mathbf{x} = \frac{(\mathbf{r}_{\mathbf{r}})_{1 \text{ MeV}}}{(\mathbf{r}_{\mathbf{e}})_{1 \text{ eV}}} \qquad \frac{(\mathbf{r}_{\mathbf{r}})_{1 \text{ eV}}}{(\mathbf{r}_{\mathbf{r}})_{1 \text{ eV}}}$ (1)The expression for G at 1 MeV is given by:  $\sigma_{\rm f} = \xi \sigma_{\rm o} \frac{\Gamma_{\rm f}}{\Gamma_{\rm p} + \Gamma_{\rm n} + \Gamma_{\rm n}}$ (2)  $\sigma$  is taken as  $\pi R^2$  (R nuclear radius), which defines the sticking probability  $\xi$ . Since for 1 MeV: R/X = 2.3,  $\xi$  won't be UNCLASSIFIED **VERIFIED UNCLASSIFIED** ntains information affecting the na 11-55-16 ing o transmus 1995 manner to an unamiorized meon sprohibited by lay - TARPBOYED, FRE PURLIC RELEASE

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larger than unity.  $\Gamma_n$  is the neutron width, which consists of several terms:  $\Gamma_n = \Gamma_{n0} + \Gamma_{n1} + \Gamma_{n2} + \dots$ , corresponding to the elastic recmission  $(\Gamma_{n0})$  and to the inelastic emissions  $\Gamma_{n1}$ . If R > X, the following relation holds:<sup>1)</sup>

$$\Gamma_{no} = \begin{cases} \sum_{l=0}^{L} (2l+1) \frac{D_{l}}{T} \end{cases}$$

L is the maximum  $\ell$  which can get into the nucleus  $(L = R/\lambda)$  and  $D_{\ell}$ is the average distance between the states which can be formed by neutrons with the angular momentum  $\ell$ , by collision with a nucleus in a state with definite quantum numbers. Let us call  $\Delta$  the distance between degenerate levels of given angular momentum, J, and assume  $\Delta$  independent of J, we get  $D_{\ell} = \frac{\Delta}{2(2\ell+1)}$ , the factor 2 coming from the spin. One then gets

$$\Gamma_{\rm no} = \begin{cases} \frac{L+1}{\pi} & \frac{\Delta}{2} \end{cases}$$

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We then write:  $\Gamma_n = N \Gamma_{no}$  where  $N \ge 1$  is related to the number of levels which can be reached via inelastic scattering of 1 MeV neutrons. It is somewhat smaller than this number because  $\Gamma_{nj}$  is

We may neglect  $\Gamma_r$  in (2) for 1 MeV and get for  $\Gamma_f$ :

 $\Gamma_{f} = \frac{\Gamma_{n}}{\left\{\frac{\sigma_{o}}{\sigma} - 1\right\}}$ 

1) See Bohr and Wheeler Phys. Rev. <u>56</u>, 426 (1939) and also LA-24 (33 and 34) page 8, formula 17.

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smaller than  $\Gamma_{no}$ . We then get for  $\Gamma_{f}$ :

$$\frac{\mathbf{i}_{\mathbf{f}}}{\mathbf{f}} = \left[ \frac{\mathbf{f}}{\mathbf{f}} \right] \left( \frac{\sigma_{0}}{\sigma_{\mathbf{f}}} - 1 \right) \frac{\mathbf{L} + 1}{\pi} \frac{\Delta}{2}$$

After inserting  $L = R/\lambda = 2.2$  and  $\sigma_0/\sigma_f = 2$  on the basis of the experimental value  $\sigma_f = 1.6 \text{ tarns}^2$  and  $\sigma_0 = \pi R^2 = 3 \text{ tarns}$ , one obtains:

$$\Gamma_{f} = \frac{H_{1}}{2i-1} \frac{\Delta}{2}$$

 $\Delta$  can be estimated in the following way: At thermal energies  $\Delta/2$  is just equal to the level distance  $D_0$  between the levels observed by McDaniel.<sup>2)</sup> (The factor two comes from the fact that the neutrons excite levels with two values of J.) The decrease of the level distance from thermal energies to 1 MoV can be estimated by using a dependence  $e^{-\sqrt{aE}}$ for the level distance (E is the excitation of the compound nucleus) and by adjusting the constant a so that the level distance decreases from 100 - 300 kilovolts at E = 0, to 2 eV at E = 6 MeV. We then obtain a = 22 (MeV)<sup>-1</sup> and a decrease of  $\Delta$  by a factor of 2.5 if E is raised from 6 to 7 MeV.

In order to get a lower limit on  $T_{f}$ , we put  $\begin{cases} = 1, N = 2, \\ \Delta/2 = 1/3 D_{0}, \text{ and } D_{0} = 1 \text{ eV}. \end{cases}$ This gives:

I) A. O. Hanson, Cf. G18 LA Report in preparation APPROVED FOR BUBLIC. RELEASE



The most plausible value for  $T_f$  may be obtained by putting N = 2,  $\xi = 0.75$ ,  $D_0 = 2$ :

$$(\mathbf{r}_{\mathbf{f}})_{1 \text{ MeV}} \sim 2 \text{ eV}.$$

We used  $\Gamma_{f}$  at low energies in order to estimate K which is defined in (1).  $(\Gamma_{f})_{low}$  is smaller or equal to the total width T of the levels observed by McDaniel. According to his curves one may put:

One obtains for the fission width at low energies:

$$\Gamma_{p} = \frac{\Gamma}{1+\alpha}$$

which gives 0.2 eV with a value & of 0.25. Thus the ratio K obeys the relation:

$$K \ge 2.6 (1+\alpha)$$

"ith the most probable value of  $\Gamma_{f}$ , we obtain

x~8.



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